



# *Coherent control of light matter interaction*

Prof. Dr. Cleber Renato Mendonça

Photonic Group

University of São Paulo (USP), Institute of Physics of São Carlos



# ultrashort laser pulses

## Mode-locking

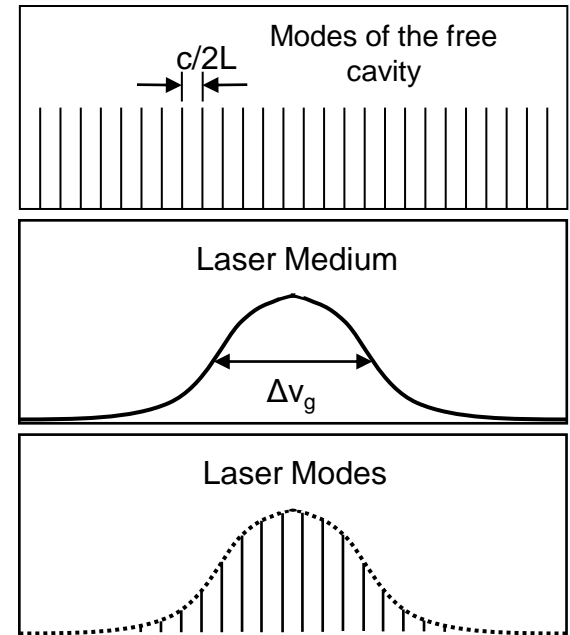
$$E_n(t) = E_n \cdot \exp[i(2\pi\nu_n t + \phi_n)] \quad n = 1, 2, 3, \dots$$

$$\nu_n = n \frac{c}{2L} \quad \nu_c \equiv \frac{c}{2L} \Rightarrow f_{\text{Laser System}}$$

$$N \approx \Delta\nu_g / \nu_c$$

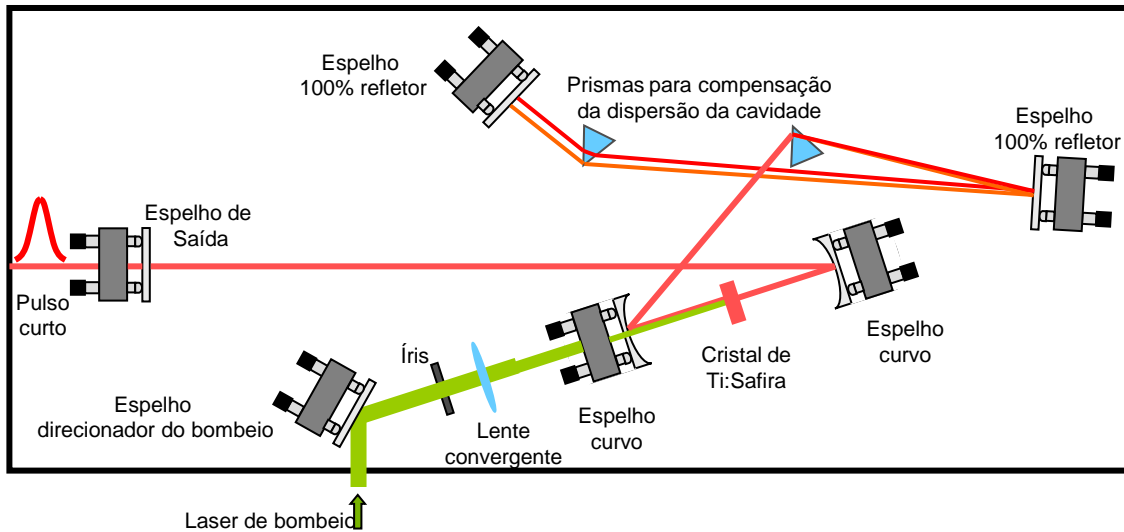
Femtosecond lasers :  $N \approx 10^5 - 10^6$

$$\tau_{\text{pulse}} \propto 1 / \Delta\nu_g$$

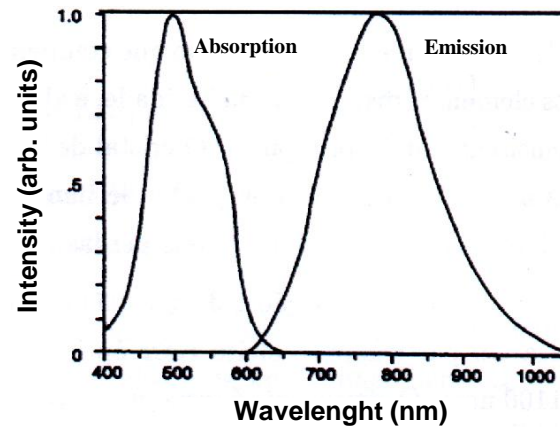


# ultrashort laser pulses

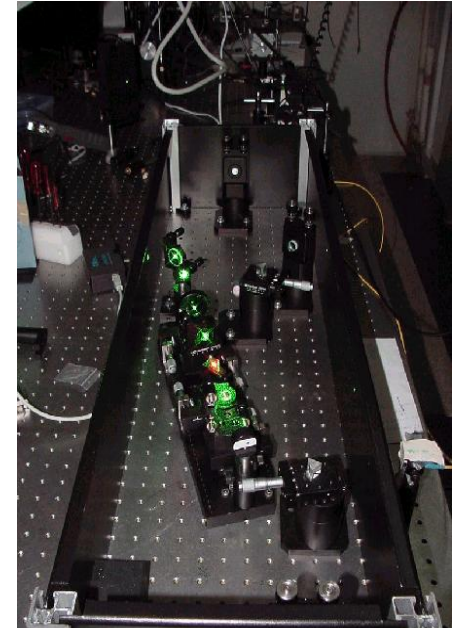
## Laser Cavity



## Ti:Sapphire Crystal



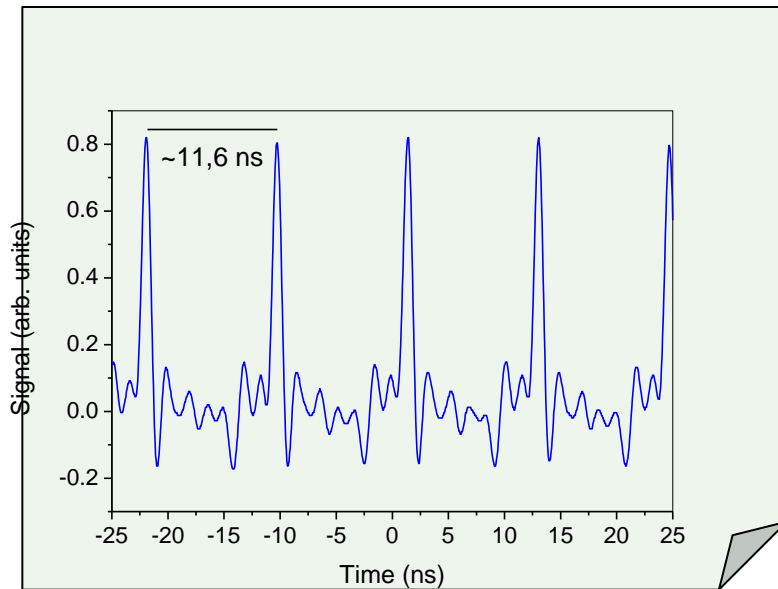
## Ti:Sapphire laser oscillator



$$\tau_{pulse} \propto 1/\Delta\nu_g$$

# *ultrashort pulses*

**Femtosecond pulses supplied for the  
Ti:Sapphire oscillator laser**



$$\Delta t_{peak} \sim 11,6 \text{ ns}$$

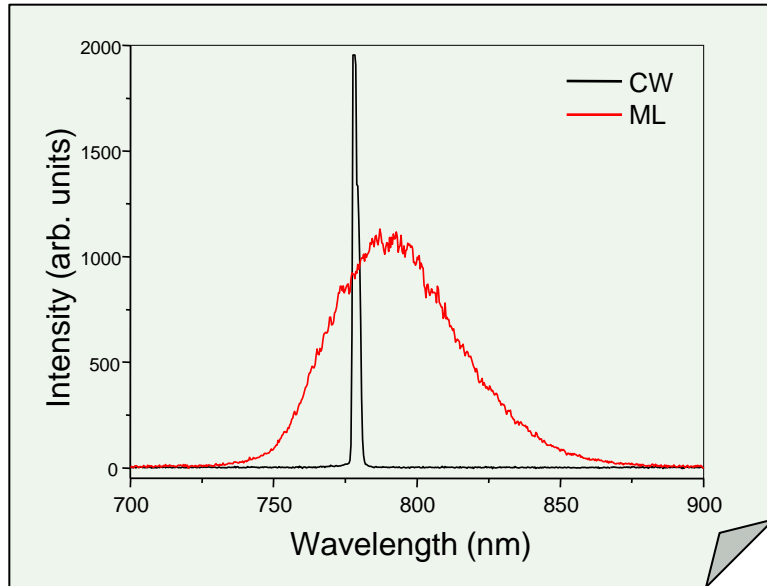
$$f \sim 86 \text{ MHz}$$

$$\overline{P}_{ML} \sim 450 \text{ mW}$$

$$\overline{E}_{pulse} \sim 5 \text{ nJ}$$

# *ultrashort pulses*

## Emission spectra of the Ti:Sapphire laser oscillator in CW and ML modes



$$\Delta\lambda_p^{FWHM} \sim 55 \text{ nm}$$

$$\lambda_0 \sim 785 \text{ nm}$$

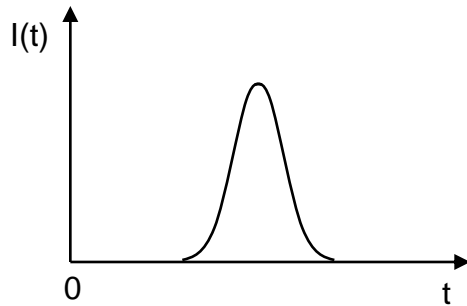
$$\overline{P}_{CW} \sim 250 - 300 \text{ mW}$$

$$\overline{P}_{ML} \sim 450 \text{ mW}$$

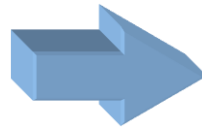
$$\tau_{TL} = 0,441 \cdot \frac{\lambda_0^2}{c\Delta\lambda_p} \Rightarrow \tau_{TL} \sim 16,5 \text{ fs}$$

Fourier Transform  
limited pulse

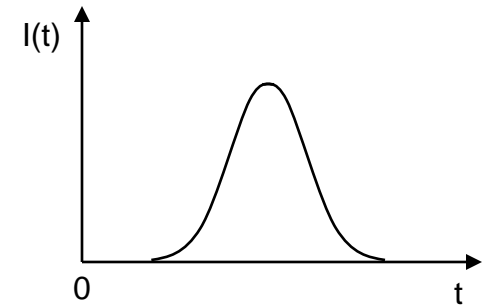
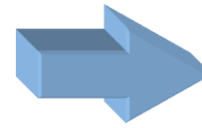
# *Dispersion of ultrashort pulse*



Initial pulse



**Dispersive  
medium**

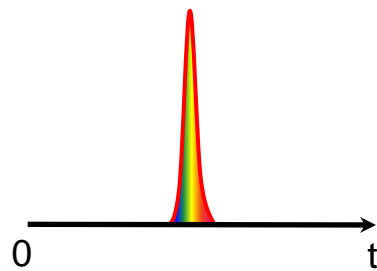
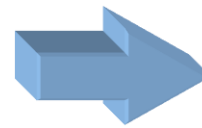


Final pulse

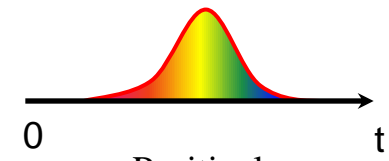
$$n(\lambda_{\text{Blue}}) > n(\lambda_{\text{redder}})$$



**Normal  
dispersive  
medium**



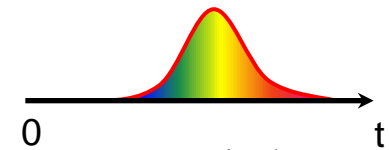
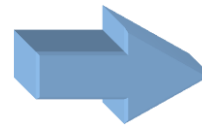
Fourier-transformed  
limit pulse



Positively  
chirped pulse



**Anomaly  
dispersive  
medium**



Negatively  
chirped pulse

$$n(\lambda_{\text{Redder}}) > n(\lambda_{\text{Blue}})$$

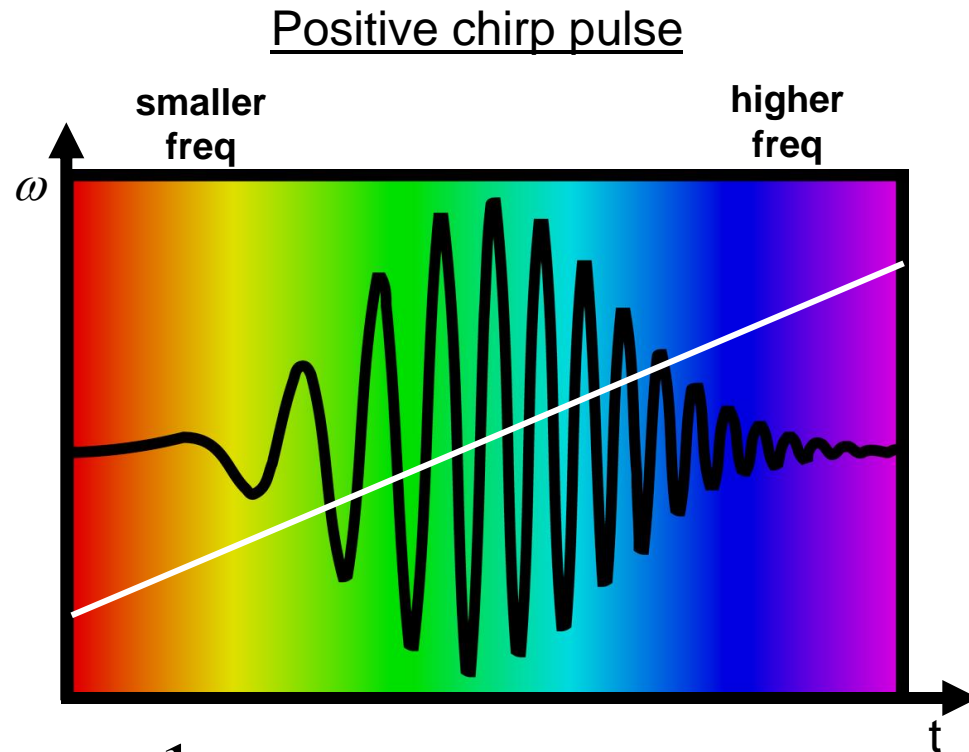
# *Dispersion of ultrashort pulse - Chirp*

- *chirp*  $E(t) = \text{Re} \left\{ \sqrt{I(t)} \exp \{ i [\omega_0 t - \phi(t)] \} \right\}$

$$\omega_{inst}(t) \equiv \omega_0 - \frac{d\phi}{dt}$$

$$\phi(t) = \phi_0 + \phi_1 t + \frac{1}{2} \phi_2 t^2 + \dots$$

$$\varphi(\omega) = \varphi_0 + \varphi_1 (\omega - \omega_0) + \frac{1}{2} \varphi_2 (\omega - \omega_0)^2 + \dots$$



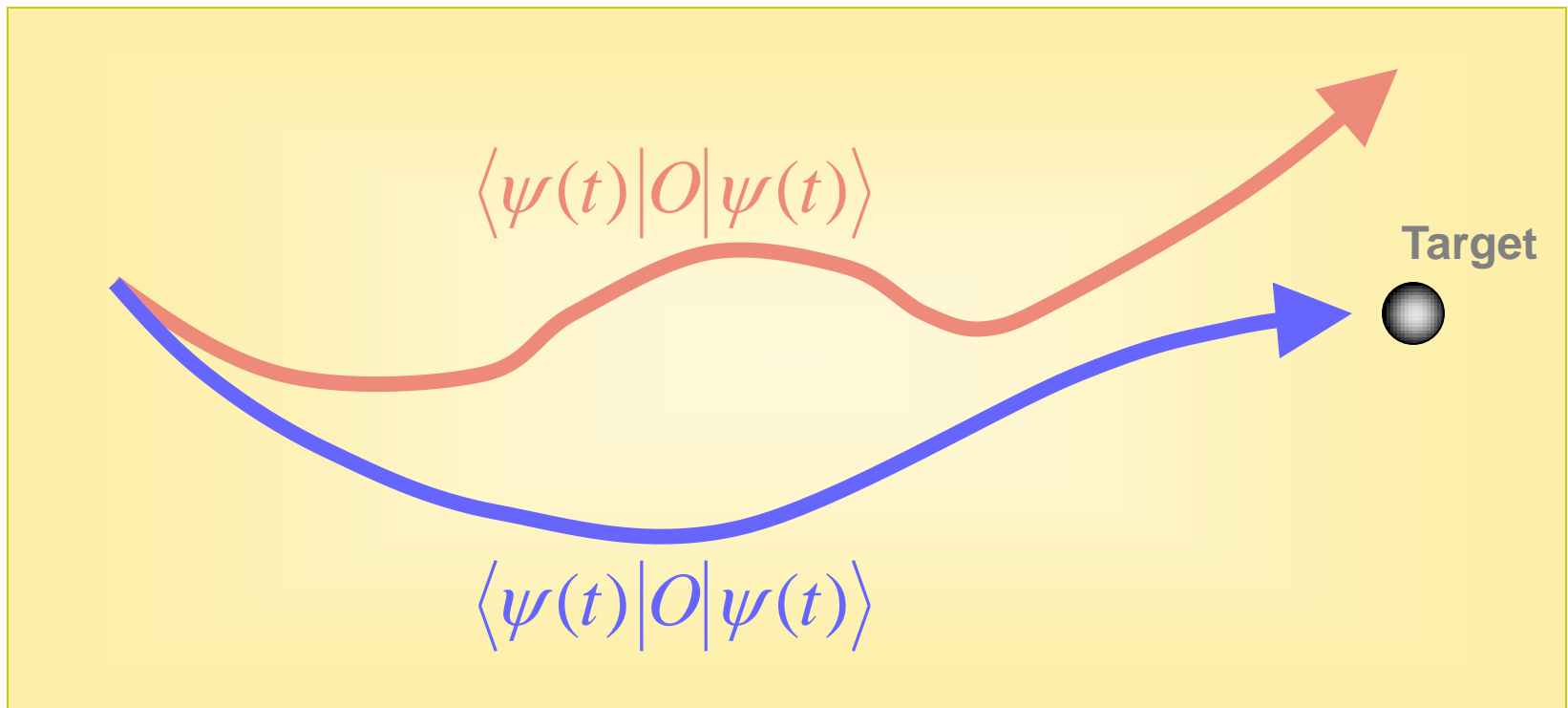
# *ultrashort pulses*

- Short temporal duration
  - High light intensity
  - Allows nonlinear effects
- wide spectral band
  - Control nonlinear process and photo-reactions

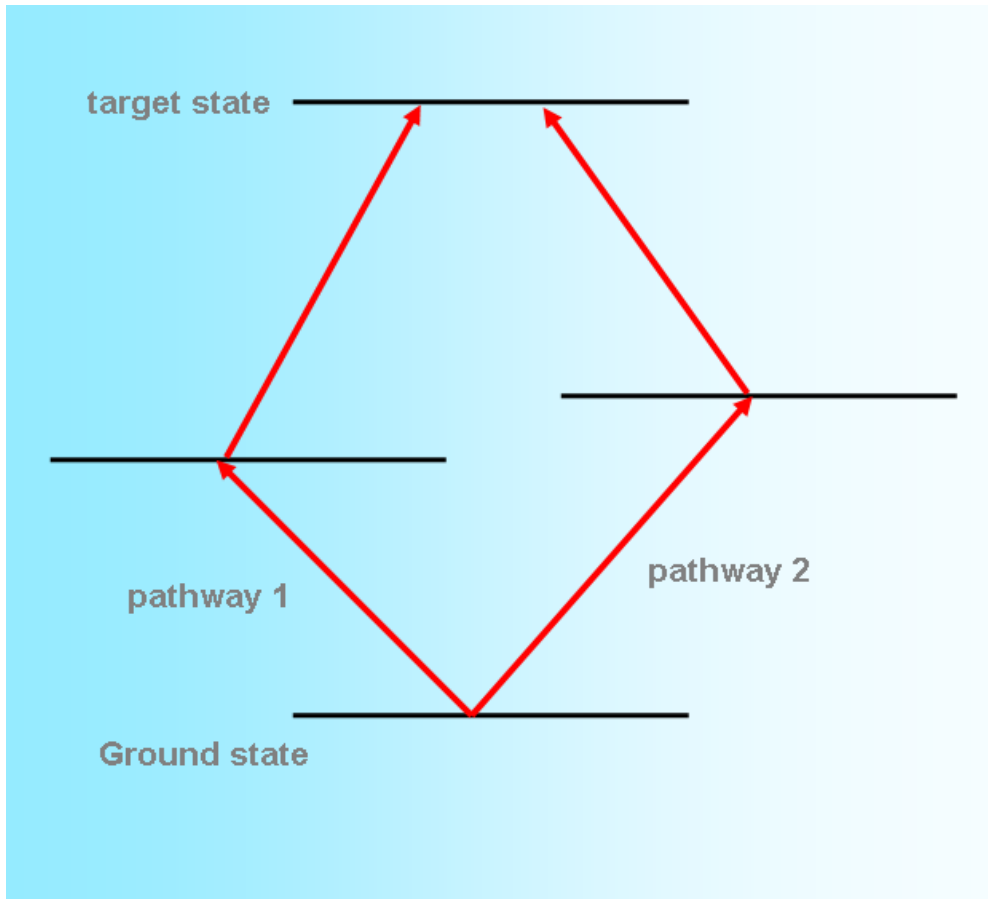


# *Coherent control*

**General idea:** use the broad spectral band of ultrashort pulses to direct a given optical process



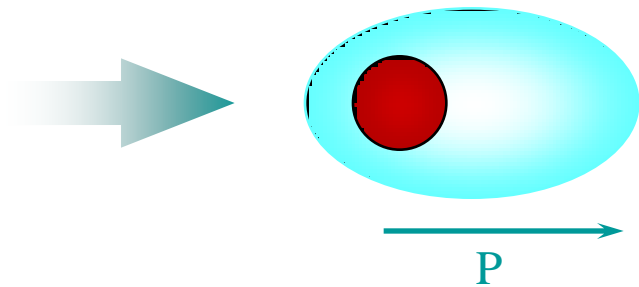
# *Coherent control*



Distinct combinations of photons of the same pulse can lead the system to a given final state through different pathways

How does that actually work ?

# Nonlinear optics



\* *High intensities*

$$\mathbf{E}_{\text{rad.}} \sim \mathbf{E}_{\text{inter.}}$$

Charges:  
Anharmonic Oscillator

$$\vec{P} = \chi^{(1)} \cdot \vec{E} + \chi^{(2)} : \vec{E}\vec{E} + \chi^{(3)} : \vec{E}\vec{E}\vec{E} + \dots$$

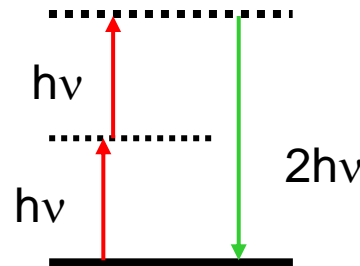
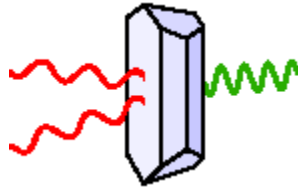
nonlinear effects are observed with high intensities

# Nonlinear optics

1961:

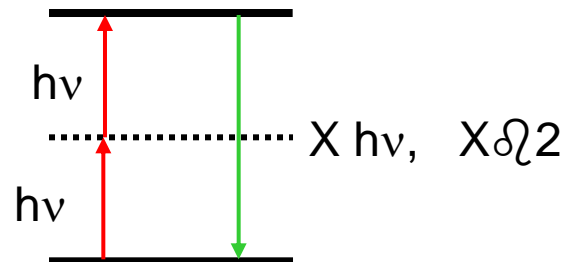
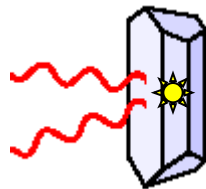
- Second Harmonic Generation (SHG), Franken *et al.*

Ruby laser beam (694.2nm) → Quartz Crystal ( $\text{SiO}_2$ ) → Laser light (347.1nm)



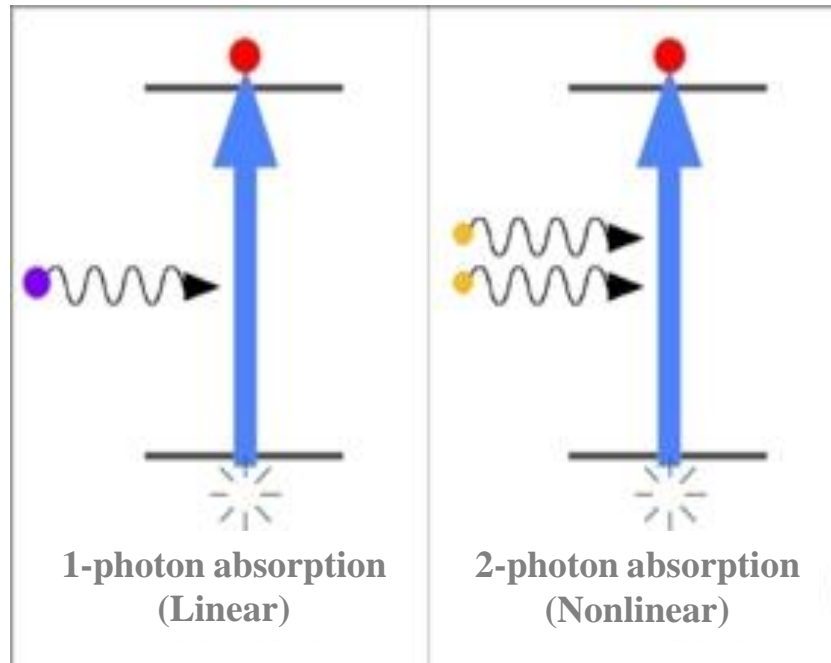
- Two-photon Absorption (2PA), Kaiser and Garret

Ruby laser beam (694.2nm) →  $\text{Eu}_2\text{:CaF}_2$  Crystal → Fluorescent light (425 nm)



# *Two-photon absorption (2PA) process*

Phenomenon does not described for the Classical Physics and **does not observed until the development of the Laser.**



**Maria Göppert-Mayer** was born June 28th 1906 in **Kattowitz**. In 1963 she received the Nobel Prize in Physics.

Theoretical model: Maria Göppert-Mayer, 1931

Two photons from an intense laser light beam are simultaneously absorbed in the same “quantum act”, leading the molecule to some excited state with energy equivalent to the absorbed two photons.

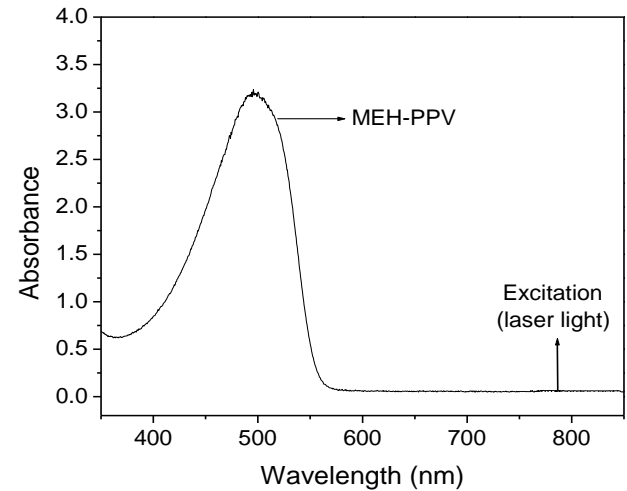
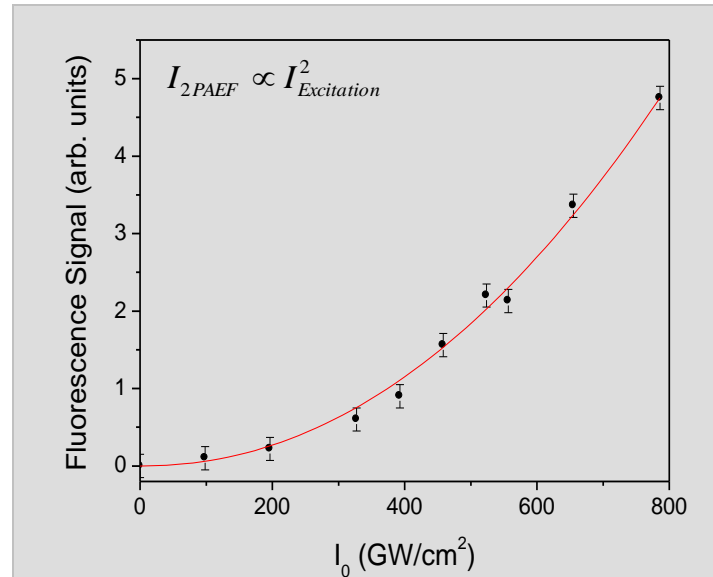
# Two-photon absorption

*Materials are excited by non-resonant light (outside of the absorption band)*

## Fluorescence (2PAEF)

$$\frac{dN_2}{dt} \approx N_1 \delta_{2PA} I^2$$

$$F \propto I^2$$



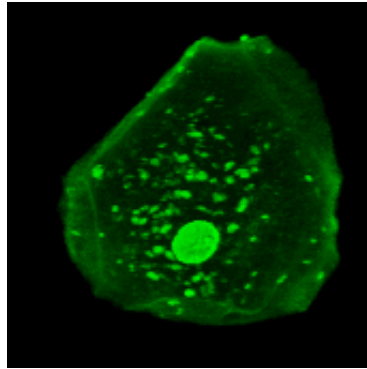
## Spatial localization of the excitation



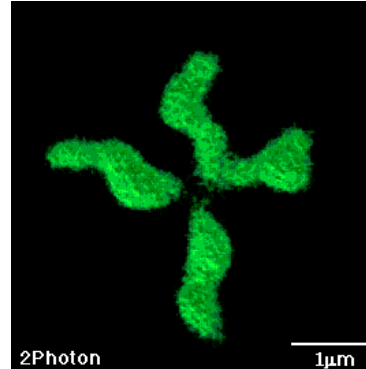
Amos, W.B. & White, J.G. (2003)

# *Applications of 2PA*

## ➤ **Microscopy by 2PAEF**



*Cell*

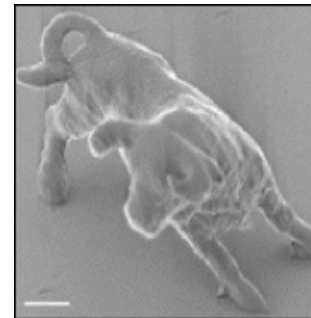


*Human chromosome*

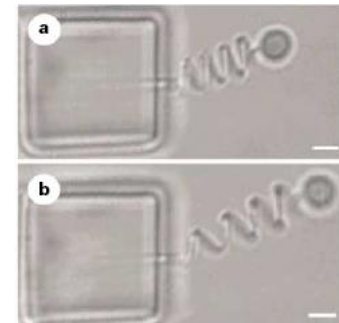
*3D Images obtained  
via 2PAEF*

## ➤ **3D Microfabrication by 2PA (Polimerization)**

*Nature 412, 697-698 (2001)*



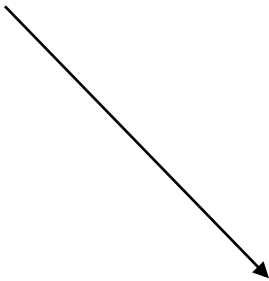
*Bull*



*Mass-spring system*

# *Applications of 2PA*

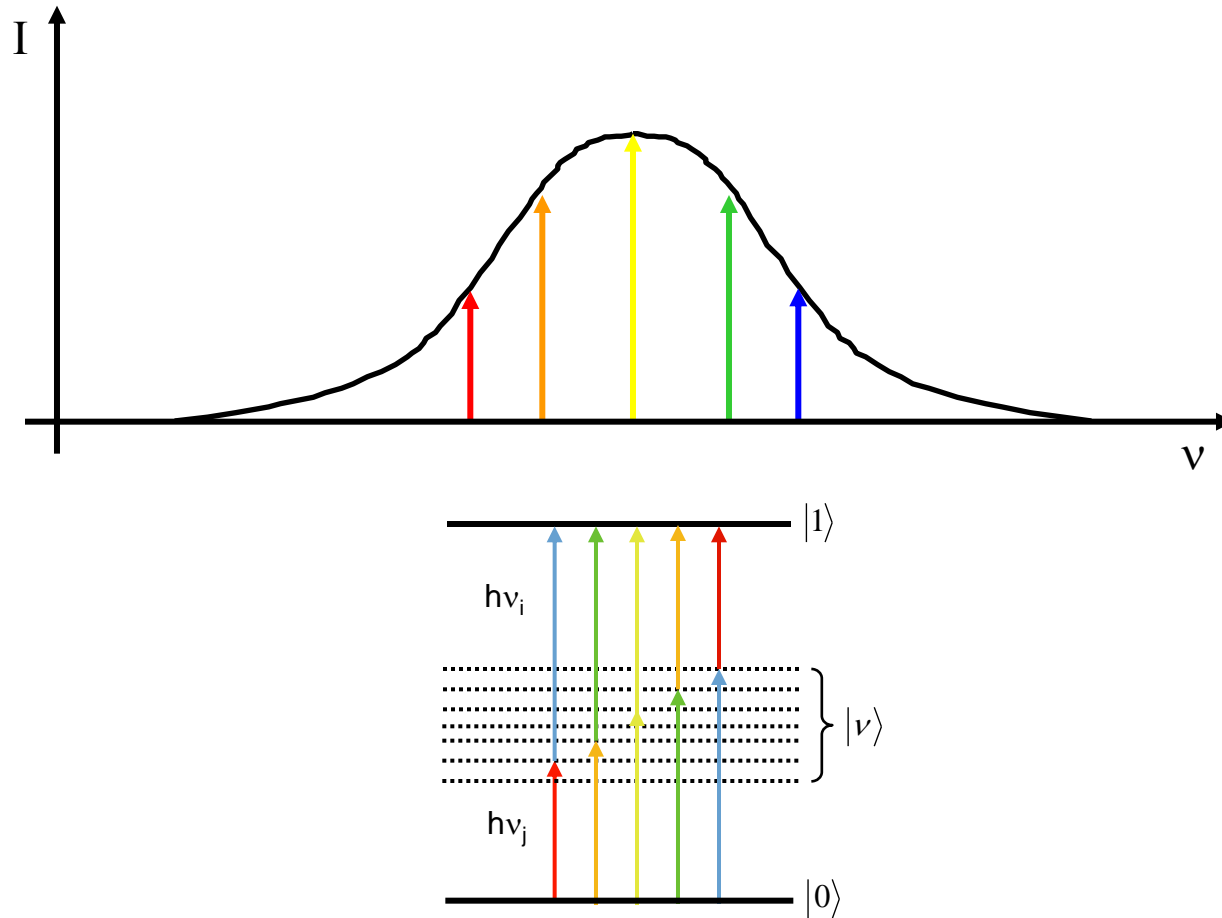
Given all the applications of 2PA, it seems to be interesting to be able to directly control it



Just playing with the fs-pulse “shape”

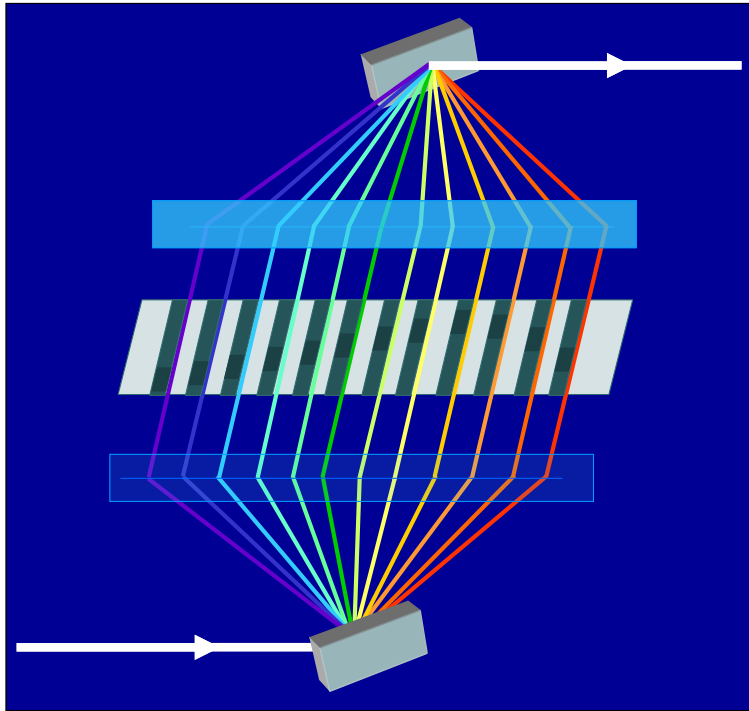


# Manipulating a two-photon process

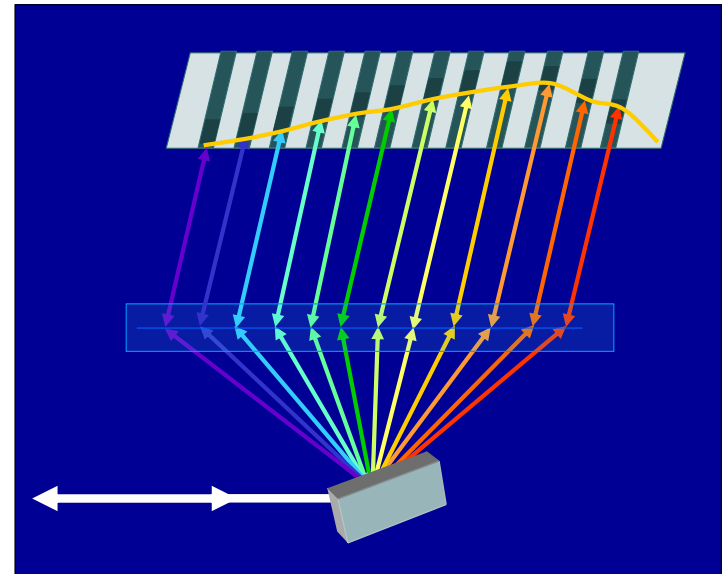


I'll have to "shape" the phase profile of the pulse

# Shaping the pulses



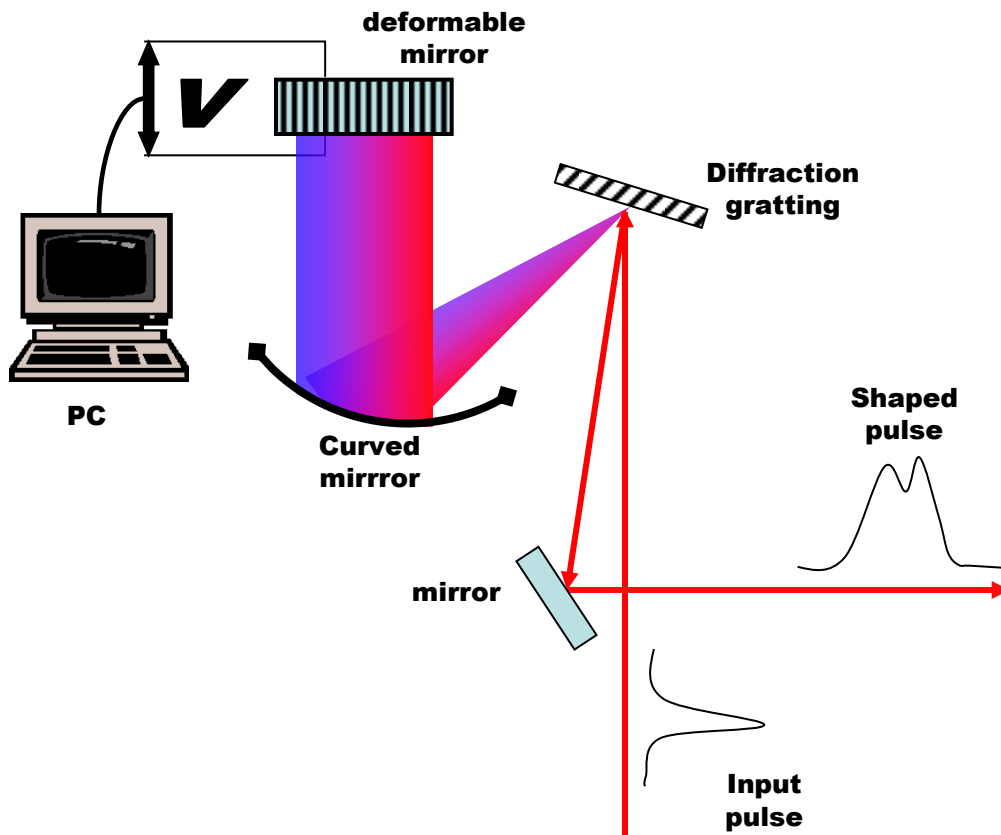
Transmission setup



Reflection setup

# Shapping the pulses

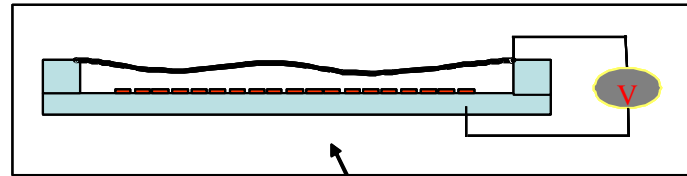
## Deformable Mirror



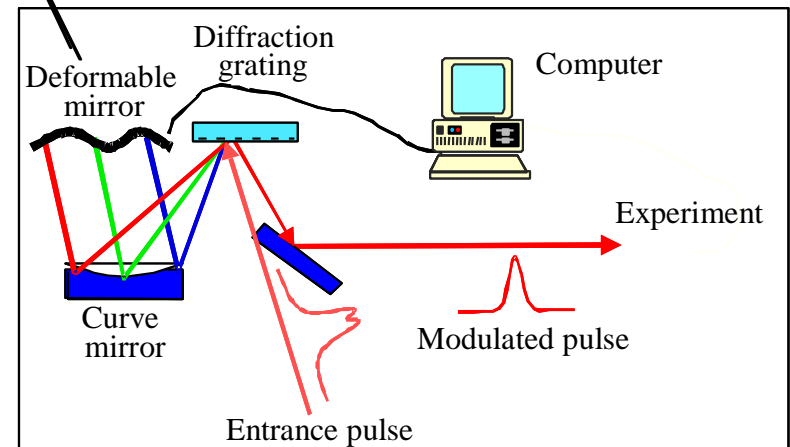
The deformable mirror acts in each spectral component of the pulse

# *Ultrashort pulse shaping technique*

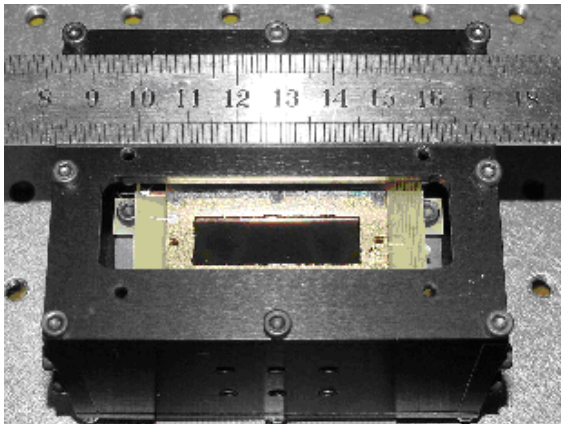
The MMDM is placed at the Fourier plane of a zero dispersion stretcher consisting of a 600 groove/mm ruled grating and a 25 cm focal-length mirror.



Experimental setup for ultrashort pulse shaping technique using deformable mirror



## **Micromachined deformable mirror (MMDM)**



Shaping the ultrashort pulse in the phase domain.

MMDM is a 600 nm gold-coated silicon nitride membrane (8 mm x 30 mm) suspended over an array of 19 actuator electrodes on a printed circuit board. Potential applied to the actuator creates an electrostatic attraction between the membrane and the electrode, deforming the mirror surface.

# Shapping the pulses

How to define which shape to use ?



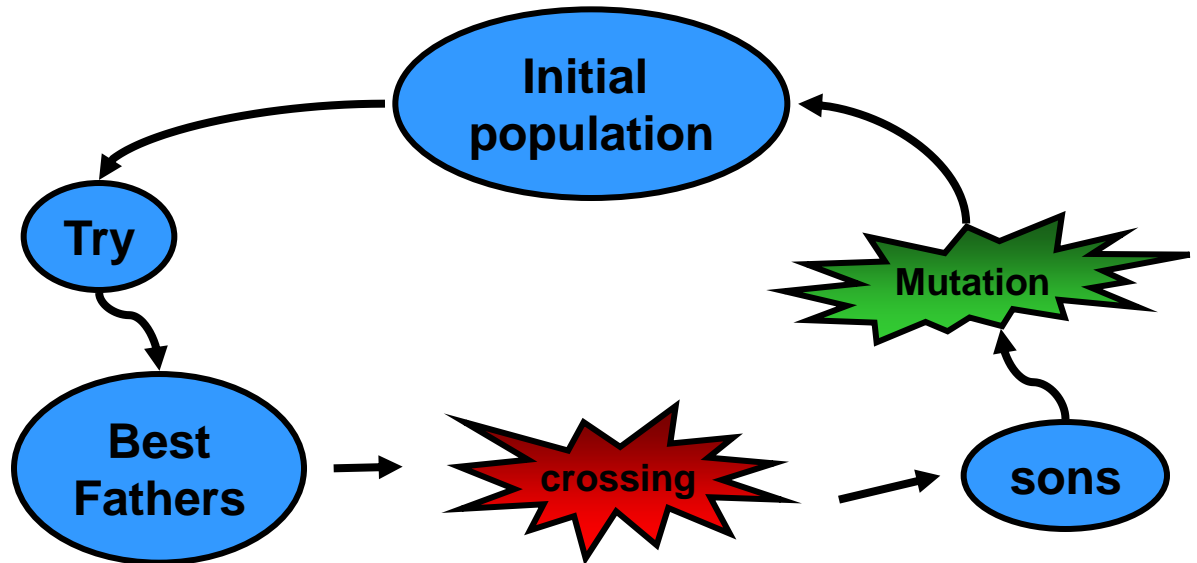
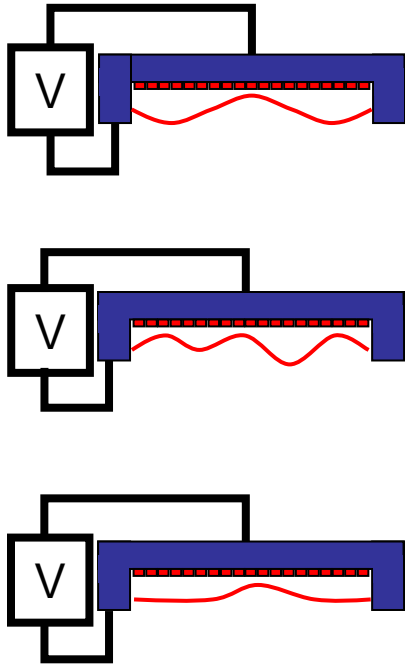
Genetic Algorithm



Phase Mask

# Genetic Algorithm

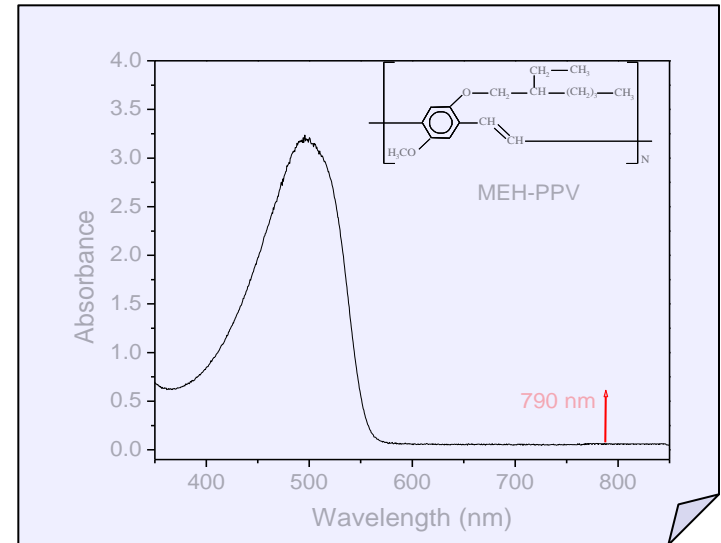
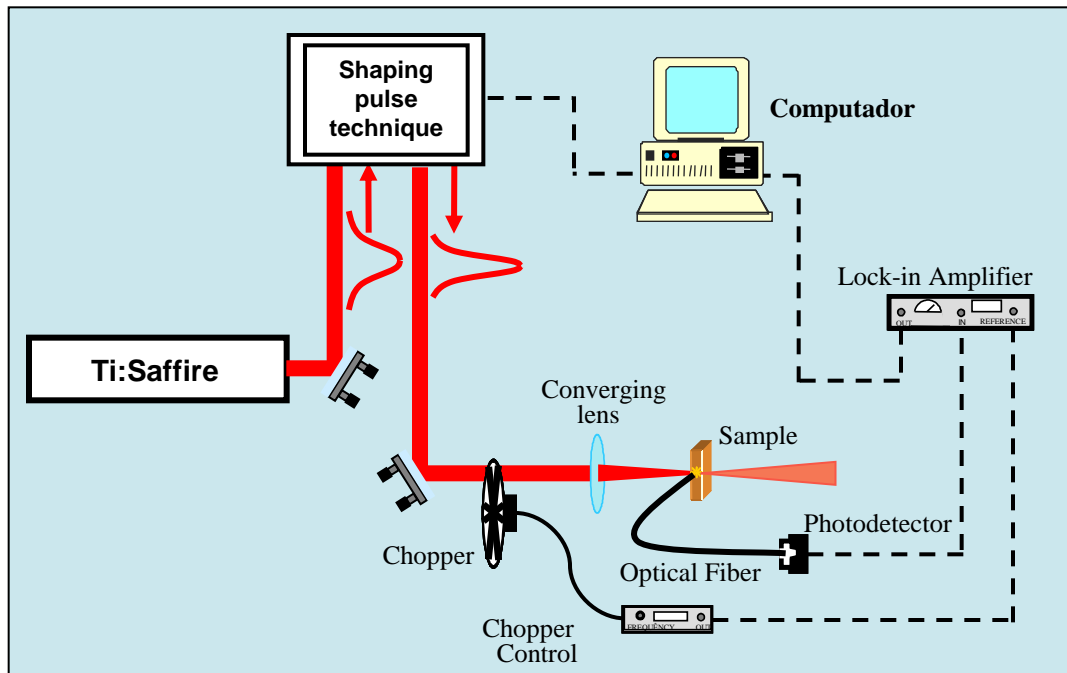
Shaping the pulse + computer algorithm



*Evolution*

In this case, the process is optimized but the mechanism is not well understood

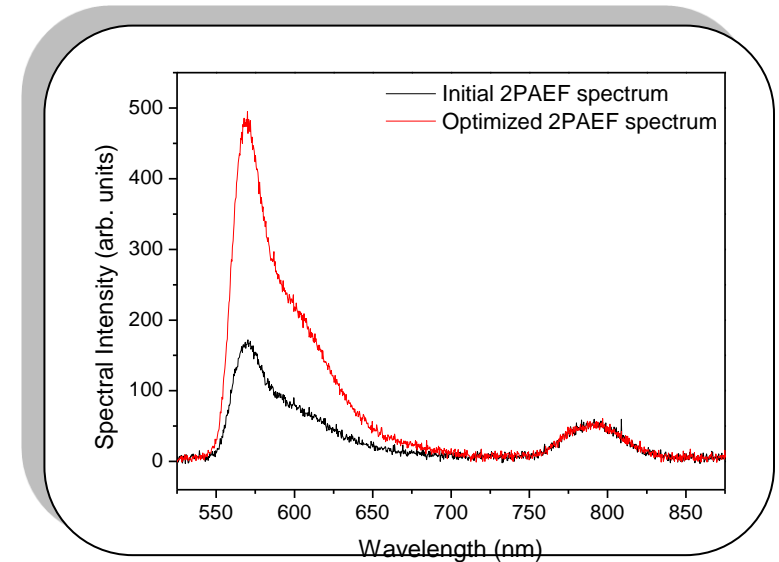
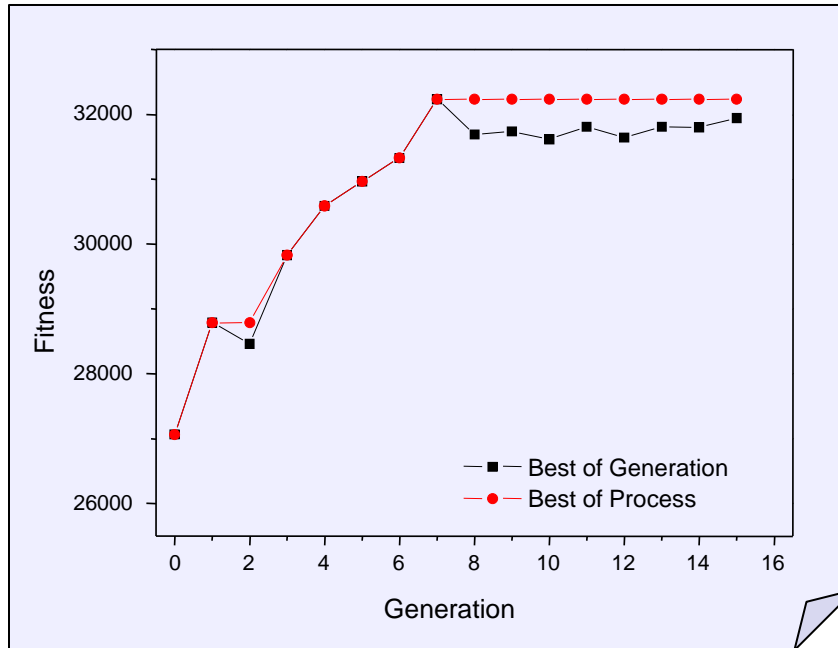
## Optimization of two-photon induced fluorescence



### UV-Vis spectra of MEH-PPV in chloroform

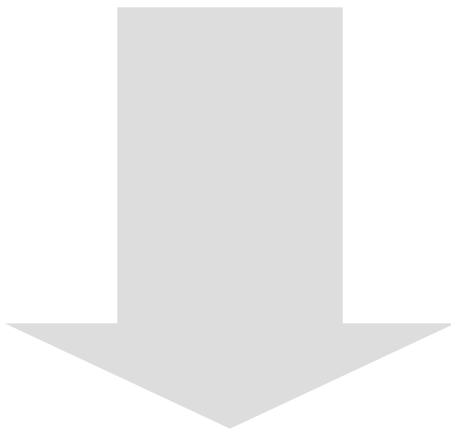
# *Optimization of two-photon induced fluorescence*

**Optimization  
process of the  
2PAEF  
spectrum**



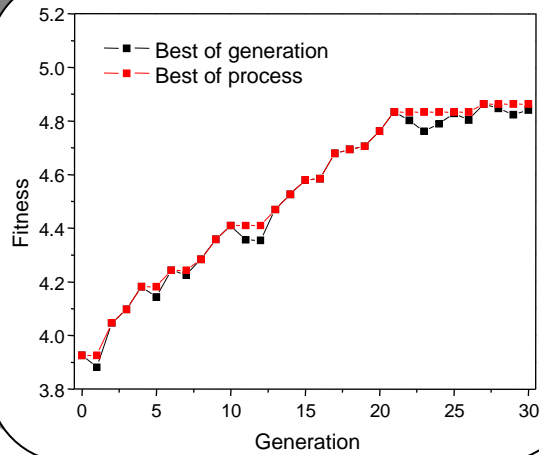


# *Optimization of two-photon induced fluorescence*

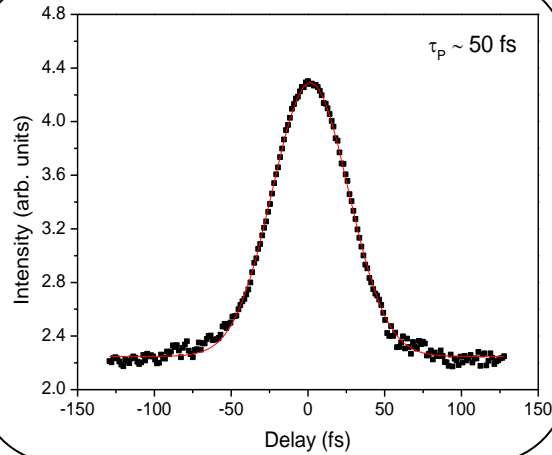


What is happening with the pulse during its evolution

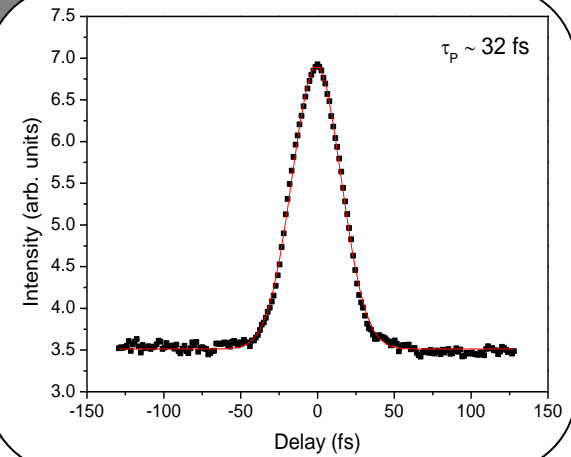
# *Optimization of the two-photon induced fluorescence*



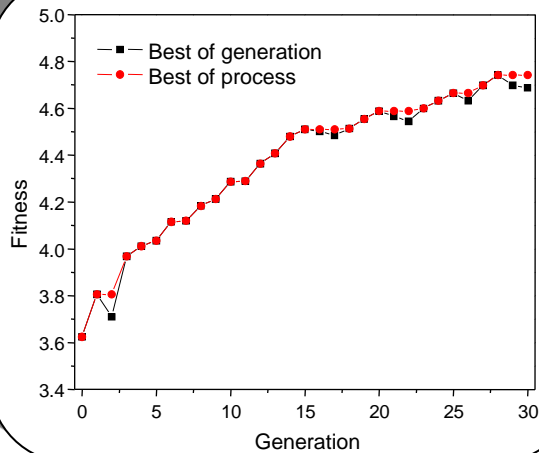
Evolution of the optimization process of the ultrashort pulse



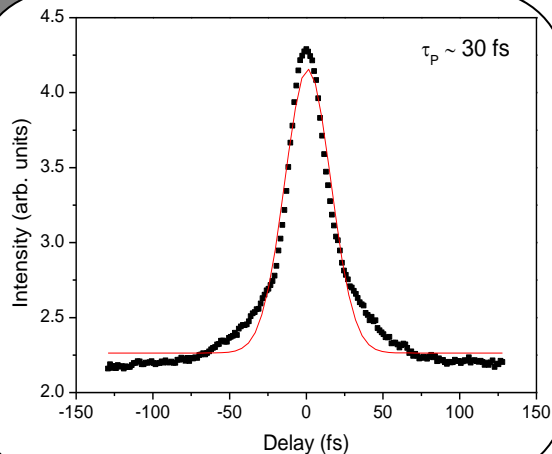
Autocorrelation obtained for non optimized pulse



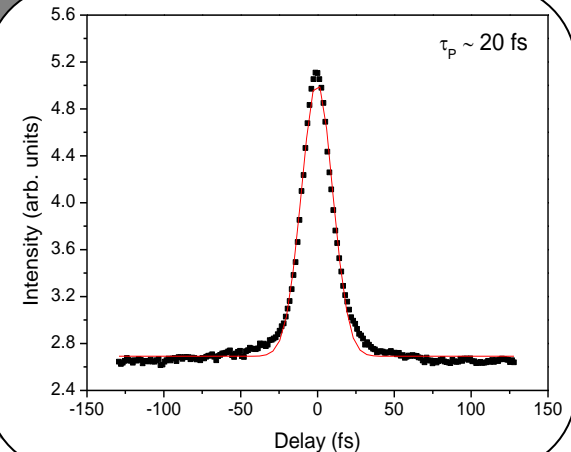
Autocorrelation obtained for optimized pulse



Evolution of the optimization process of the ultrashort pulse

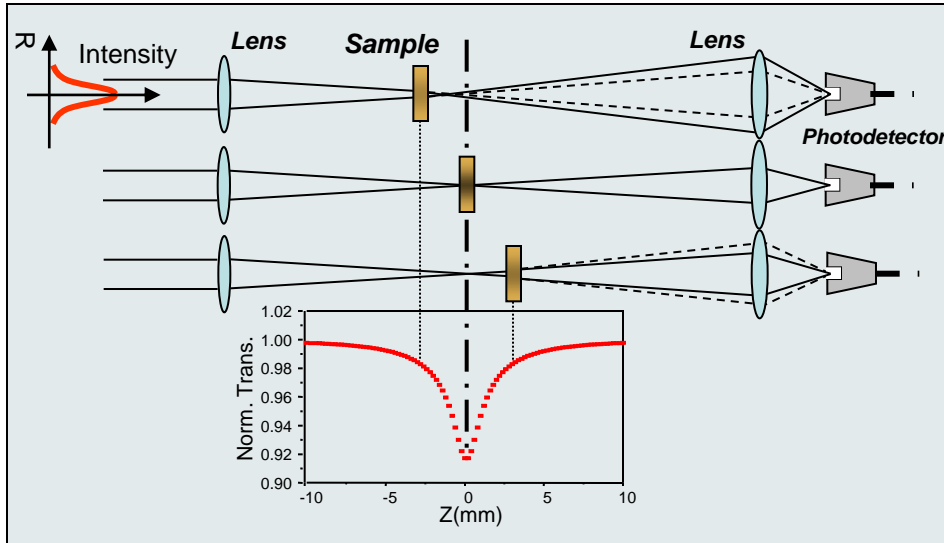


Autocorrelation obtained for non optimized pulse



Autocorrelation obtained for optimized pulse

# Optimization of two-photon induced thermal lens

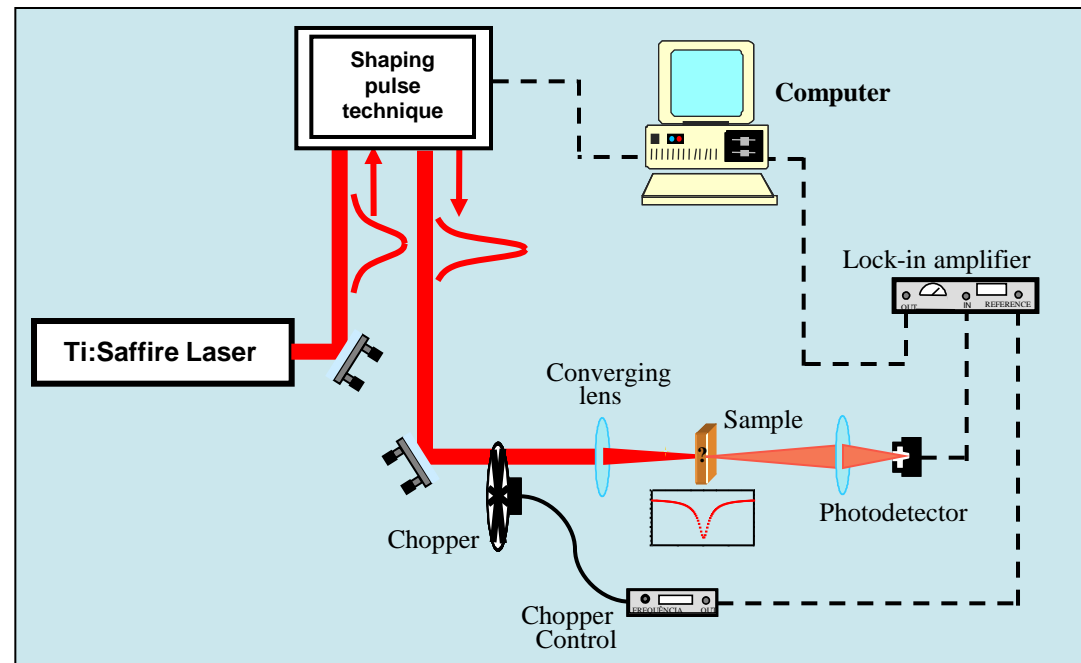


**Z-Scan technique**  
(Nonlinear absorption)

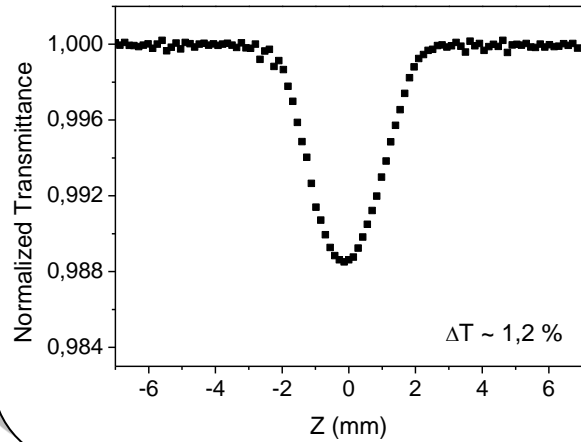
**Feedback Signal**  
**Nonlinear transmittance**

**Samples**

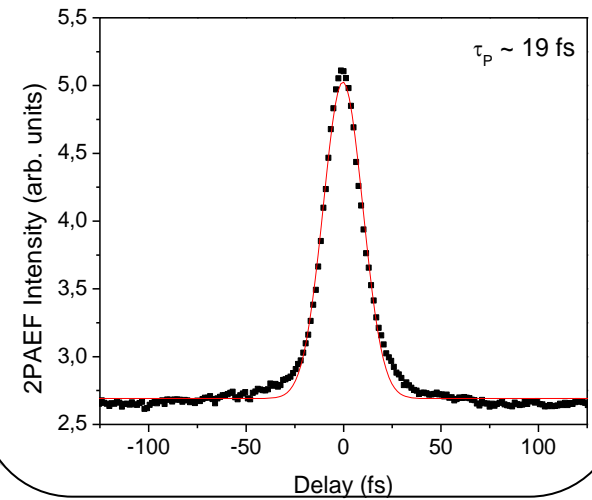
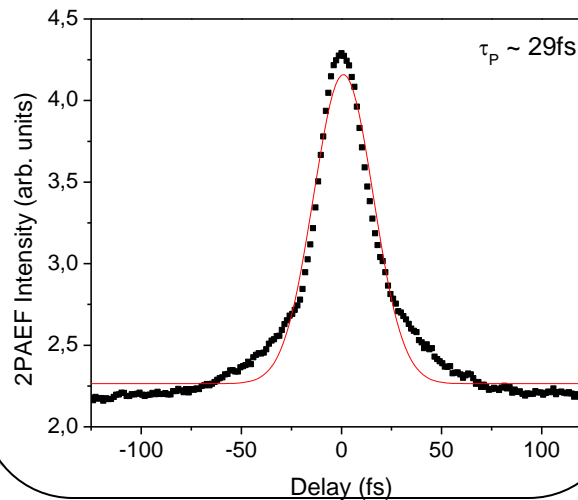
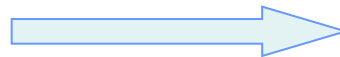
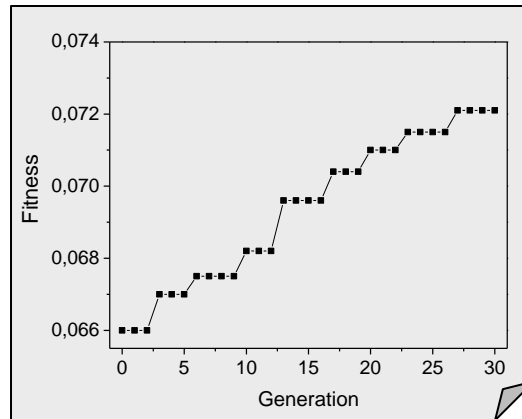
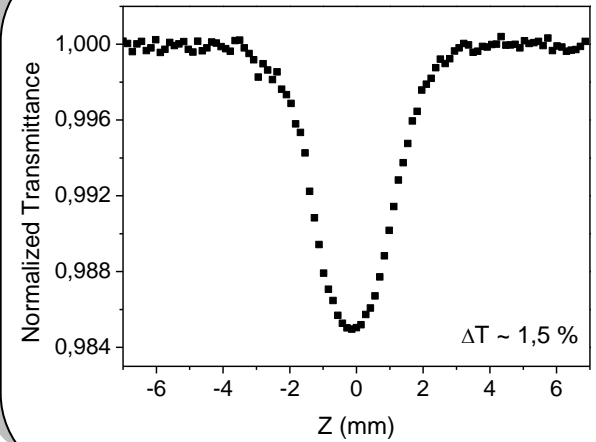
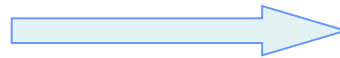
Fluorescents and  
non-fluorescents



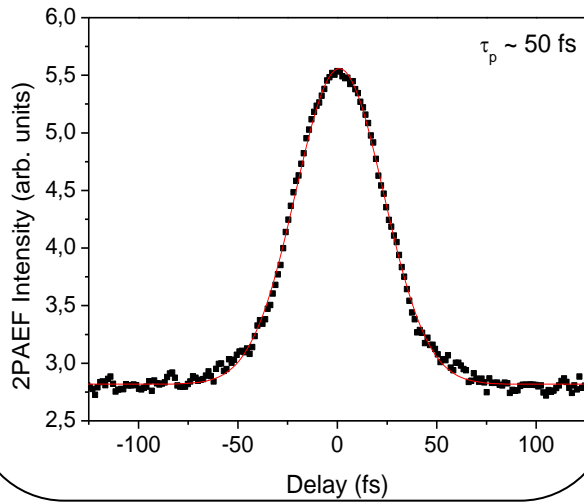
# Optimization of two-photon induced thermal lens



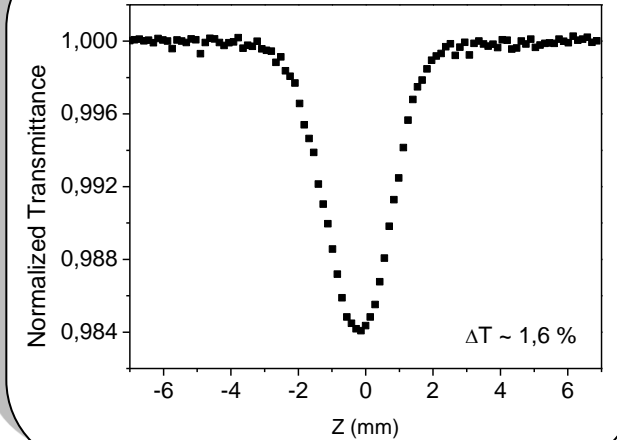
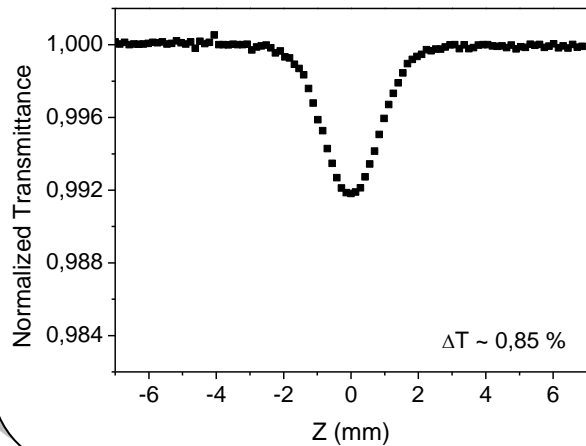
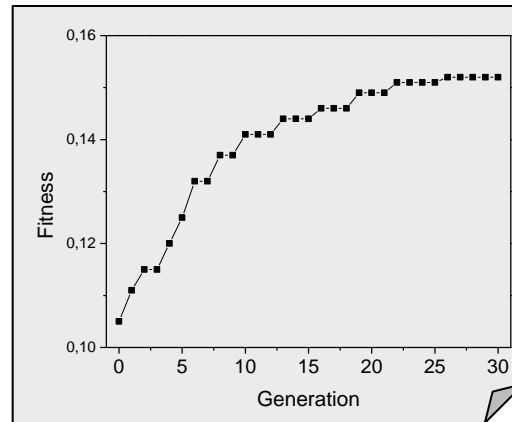
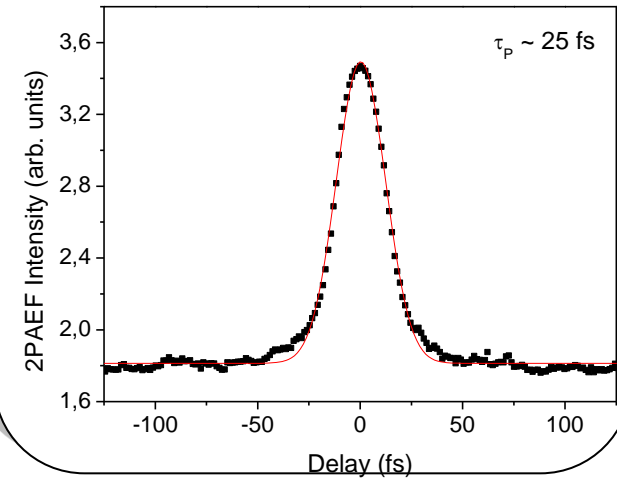
*Sample: MEH-PPV*



# Optimization of two-photon induced thermal lens



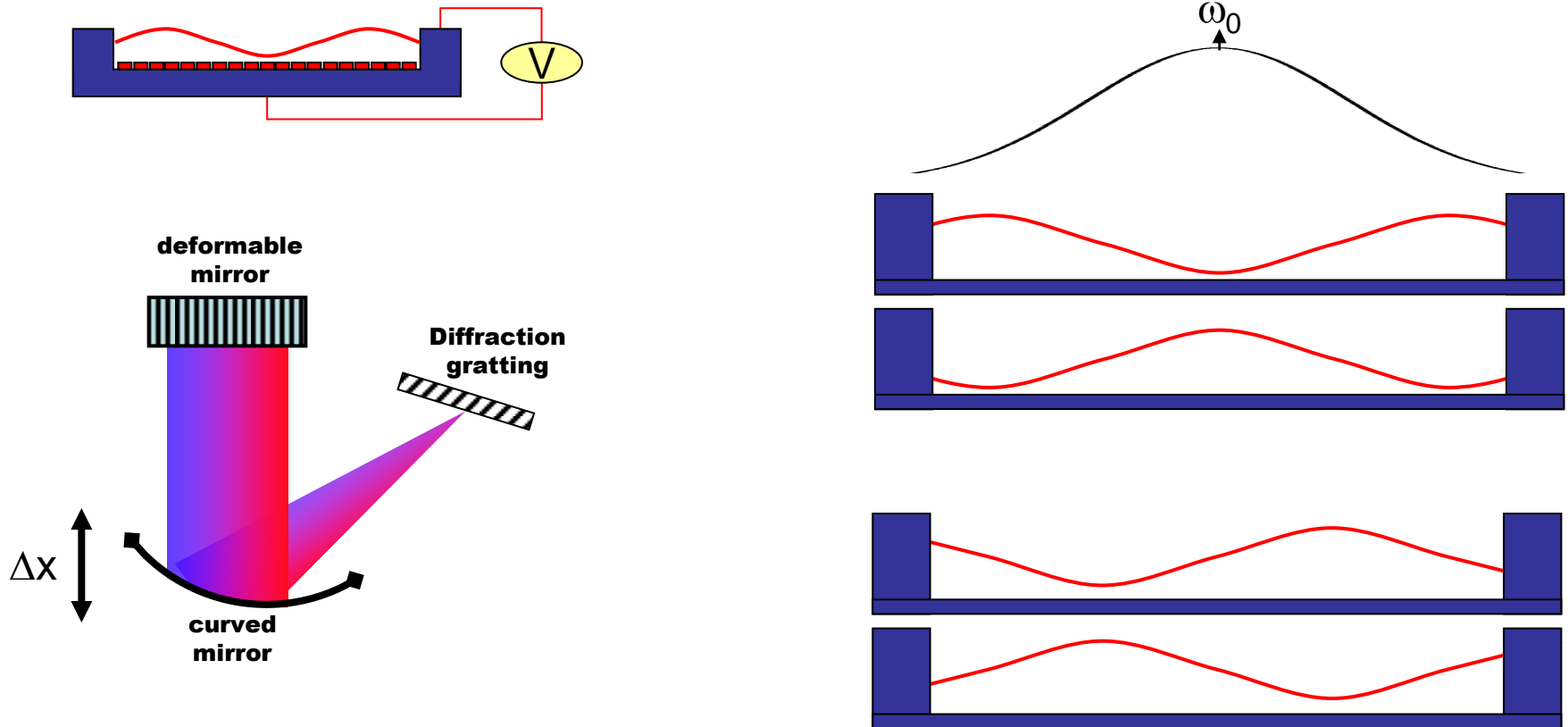
*Sample: MEH-PPV*



# Phase Mask

In this case, we impose a known phase mask function to the ultra-short pulse

$$\Theta(\Omega) = \alpha \cos(\gamma\Omega + \delta)$$



# *Phase Mask*

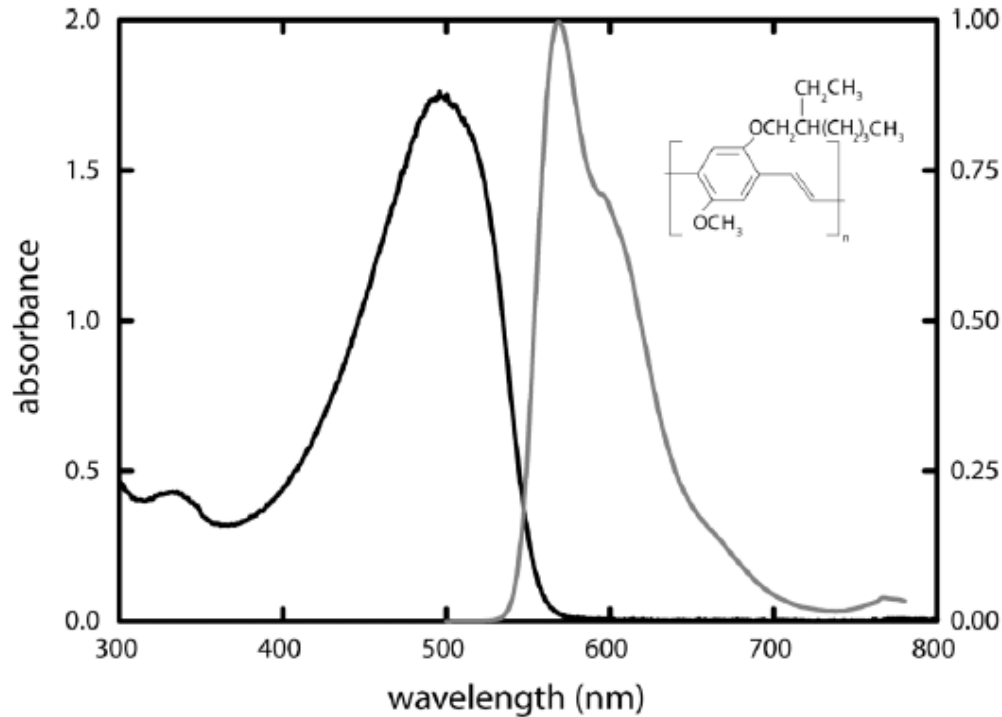
- Two-photon absorption transition is given by

$$S_2(\omega_0) \propto \left| \int_{-\infty}^{+\infty} A(\omega_0/2 + \Omega) A(\omega_0/2 - \Omega) \cdot \exp \{ i [ \varphi(\omega_0/2 + \Omega) + \varphi(\omega_0/2 - \Omega) ] \} d\Omega \right|^2$$

Anit-symmetric phase mask:  $S_2$  is independent of the phase

Symmetric phase mask:  $S_2$  is minimized

# *Controlling the MEH-PPV photo-degradation*

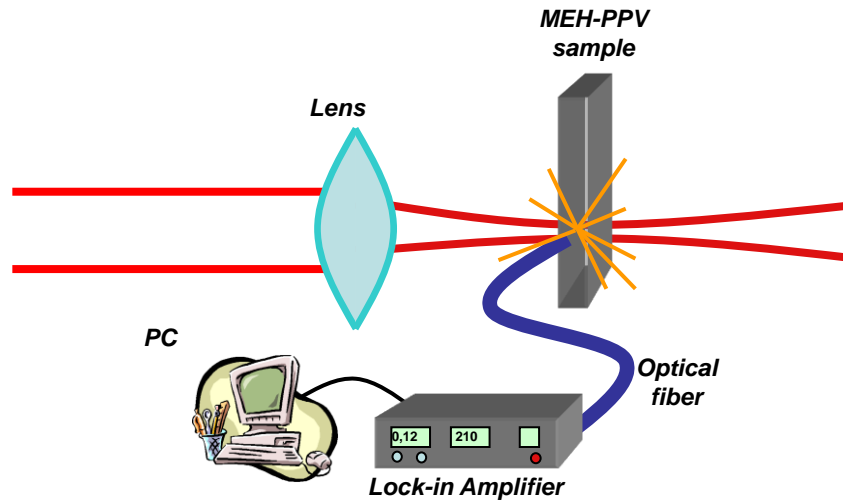


**MEH-PPV**: conductive and luminescent polymer with interesting properties for applications

However, MEH-PPV photo-bleaches due to a photooxidation reaction, causing a decrease in its emission



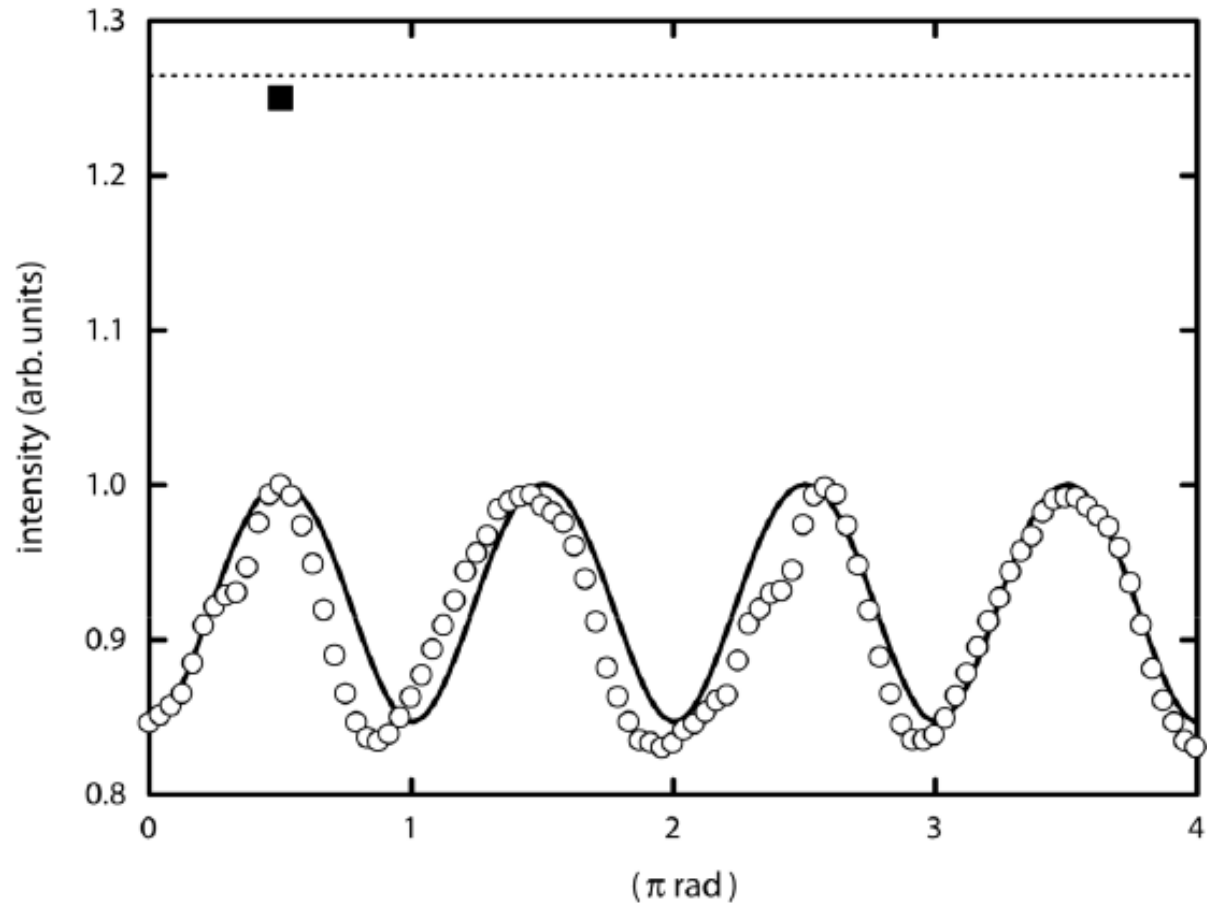
# *Controlling the MEH-PPV photo-degradation*



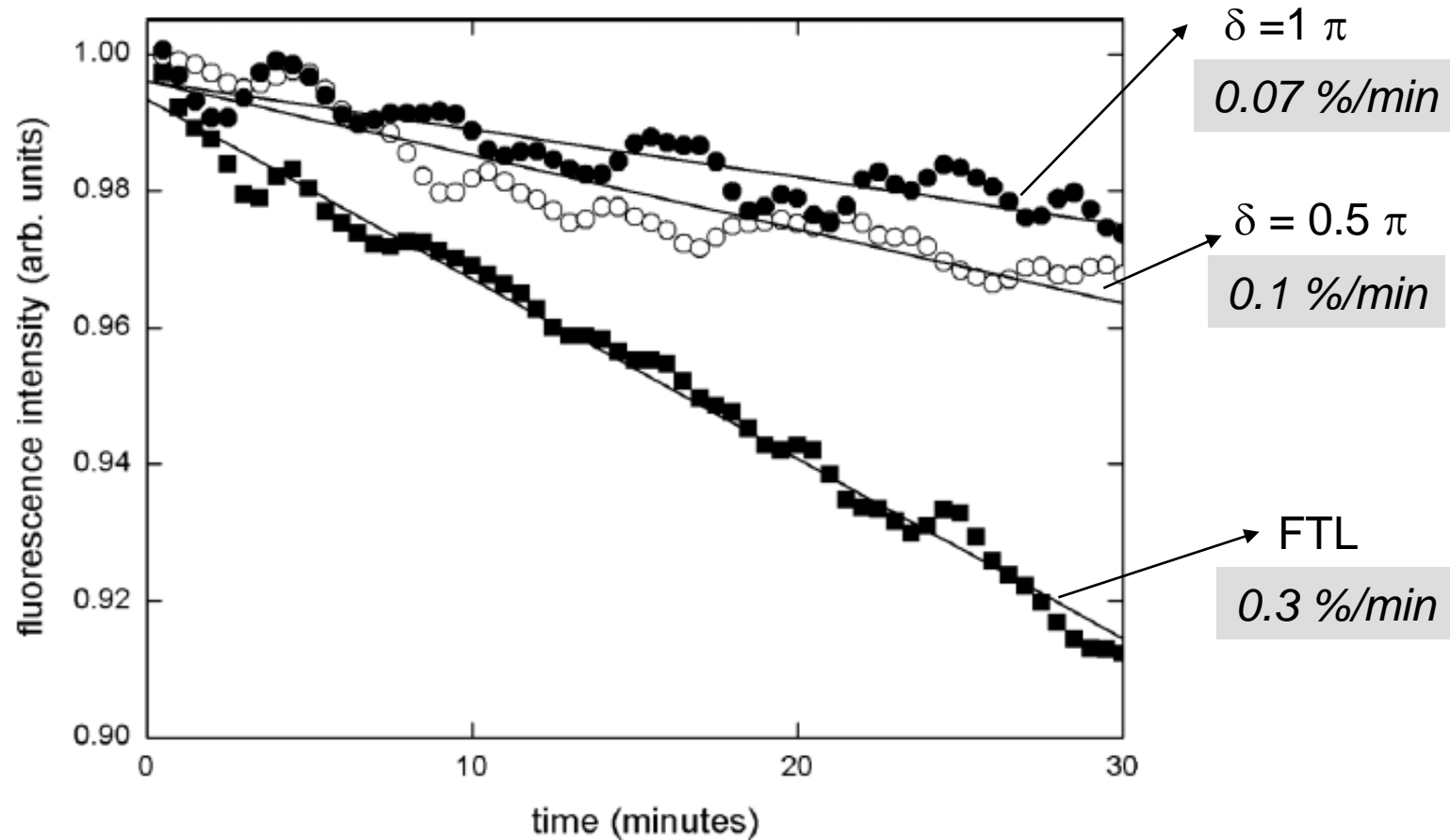
Observe the two-photon excited emission as a function of the phase-mask

Observe the photodegradation for distinct phase masks

# *Controlling the MEH-PPV photo-degradation*

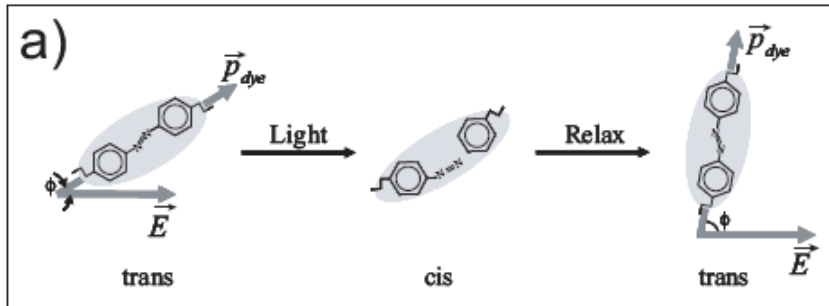


# Controlling the MEH-PPV photo-degradation

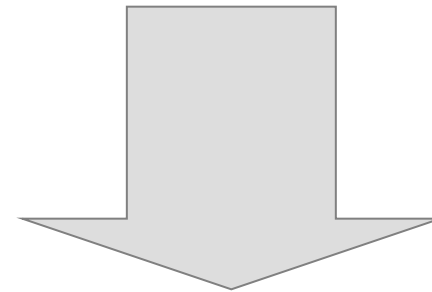
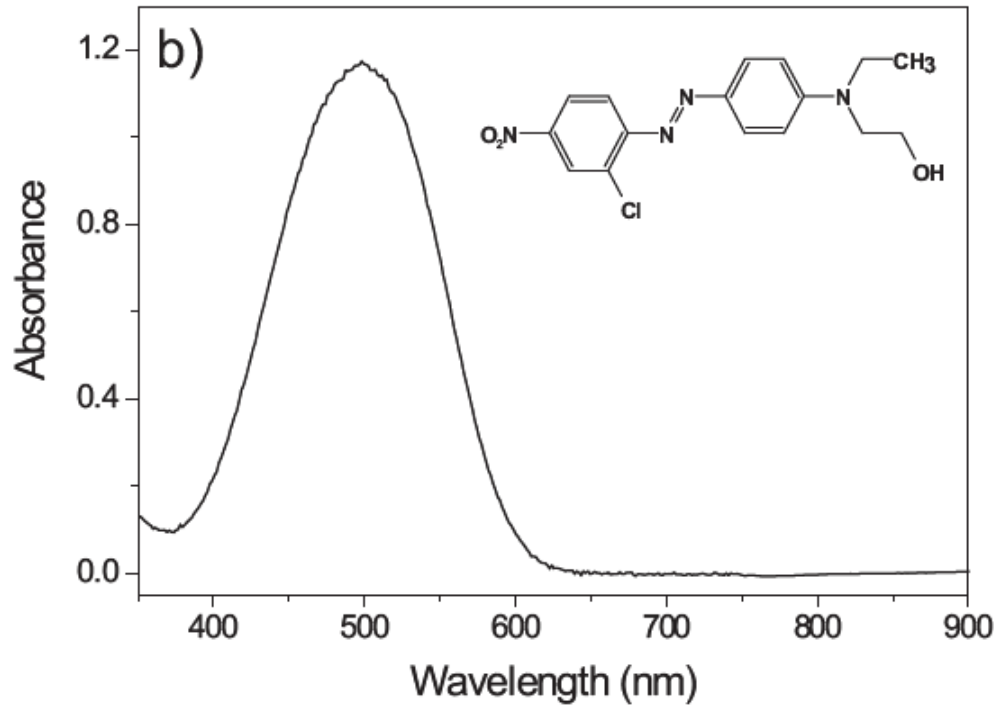


*Photobleaching rate is smaller for the phase-masked pulses*

# Coherent control of molecular orientation

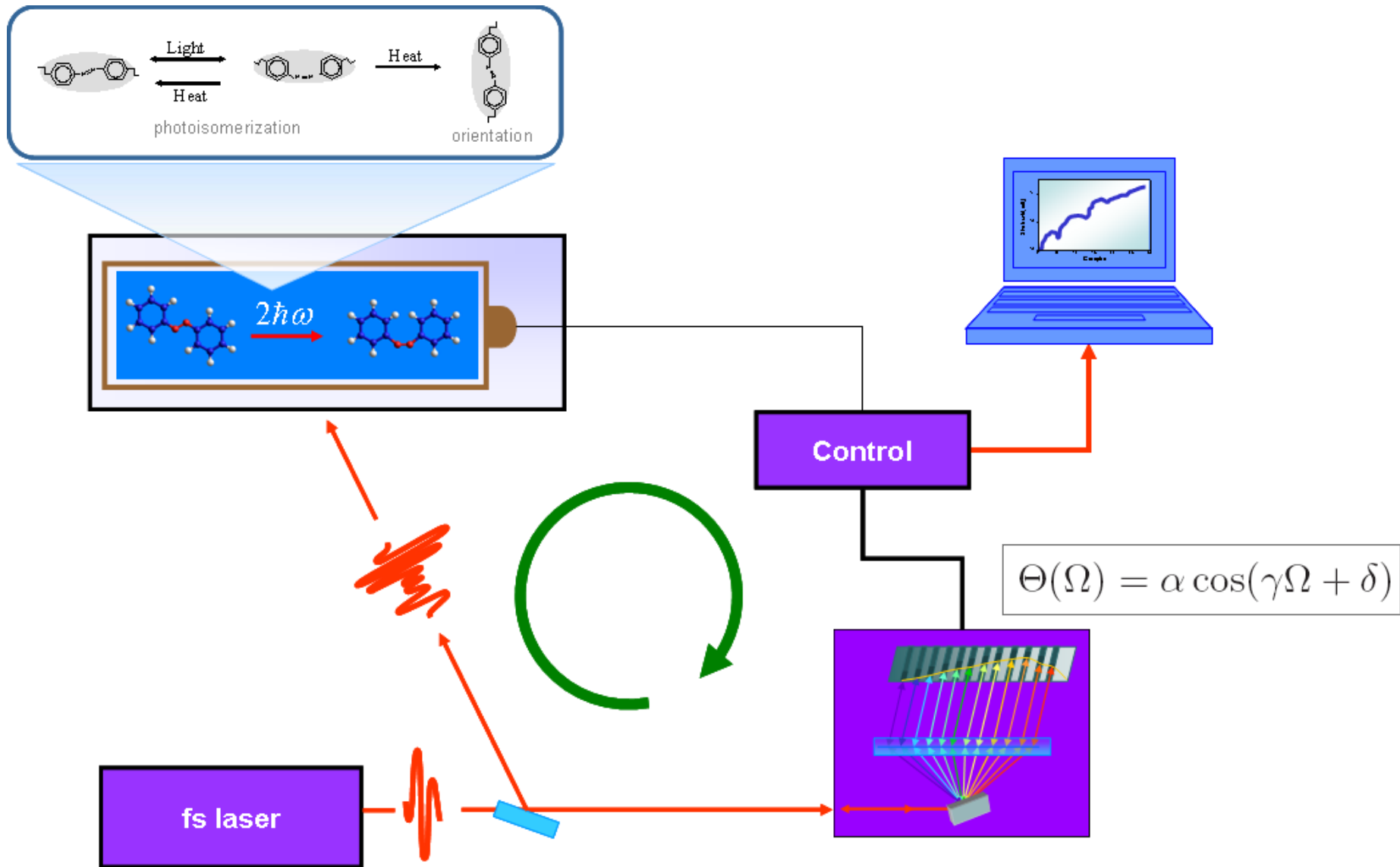


Use fs-laser (broad band) to induce two-photon absorption and, consequently, molecular orientation (optical storage)



*Coherently control the molecular orientation*

# Coherent control of molecular orientation



# Coherent control of molecular orientation

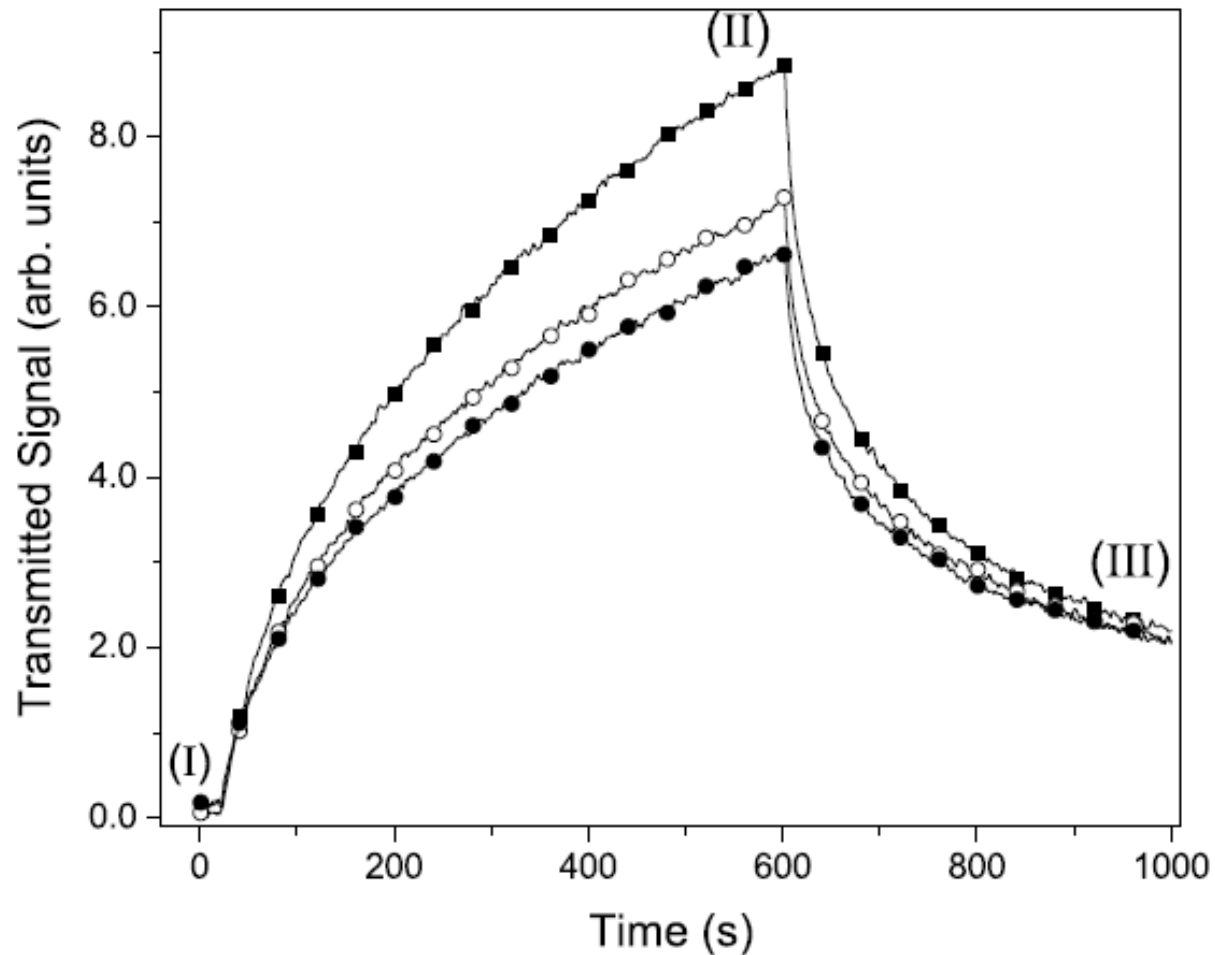
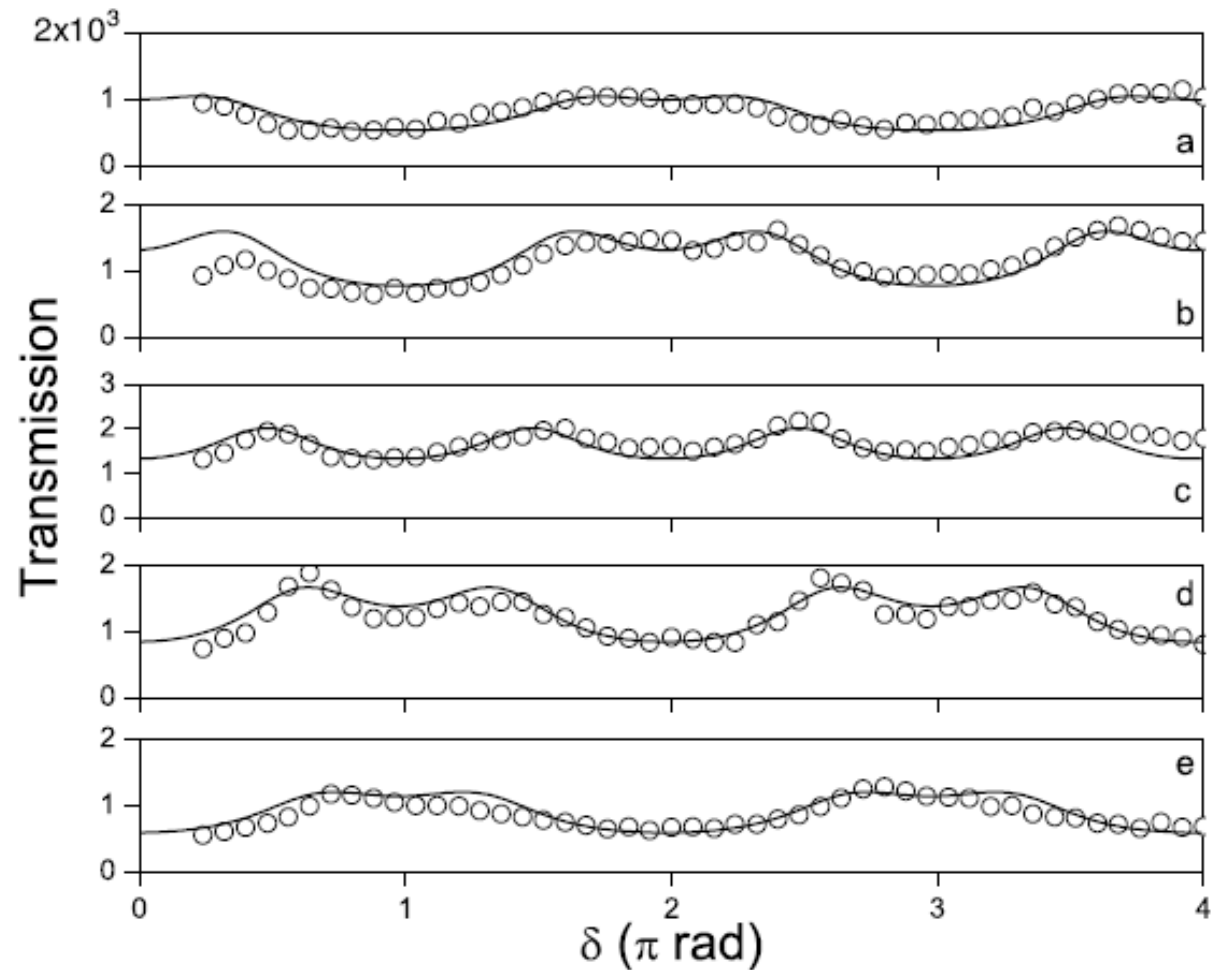


FIG. 2. Optical storage curves for the PMMA/DR13 film for three distinct  $\delta$  values:  $5\pi/10$ (■),  $7\pi/10$ (○), and  $\pi$ (●).

# Coherent control of molecular orientation

$$\phi(\Omega) = \alpha \cos(\gamma\Omega + \delta) + 1/2\beta\Omega^2$$



# *Conclusions*

- Pulse shaping methods + coherent control of the nonlinear interaction seems to be an interesting method to further control nonlinear optical processes.



Thanks for your attention!!!