

Optical nonlinearities in organic materials

Prof. Cleber R. Mendonca



<http://www.fotonica.ifsc.usp.br>

University of Sao Paulo - Brazil



students 77.000
52.000 undergrad.
25.000 grad.
employers 15.000
professors 6.000

- Sao Paulo
- Sao Carlos (9.000)
- Ribeirao Preto





University of Sao Paulo – in Sao Carlos





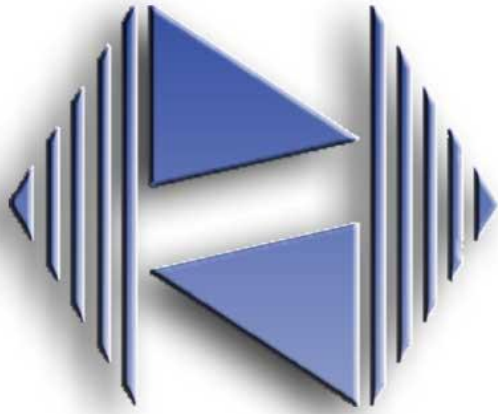
University of Sao Paulo – in Sao Carlos



University of Sao Paulo – in Sao Carlos Campus II (new Campus)



Instituto de Física de São Carlos



Professors: 68

Employers: 160
(technical and administration)

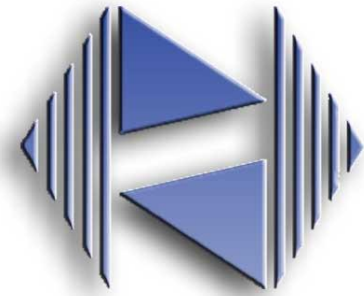
Students: 450 (undergrad)
100 (master)
140 (phD)

Several research areas in Physics
and Material Sciences





Grupo de Fotonica Photonics Groups

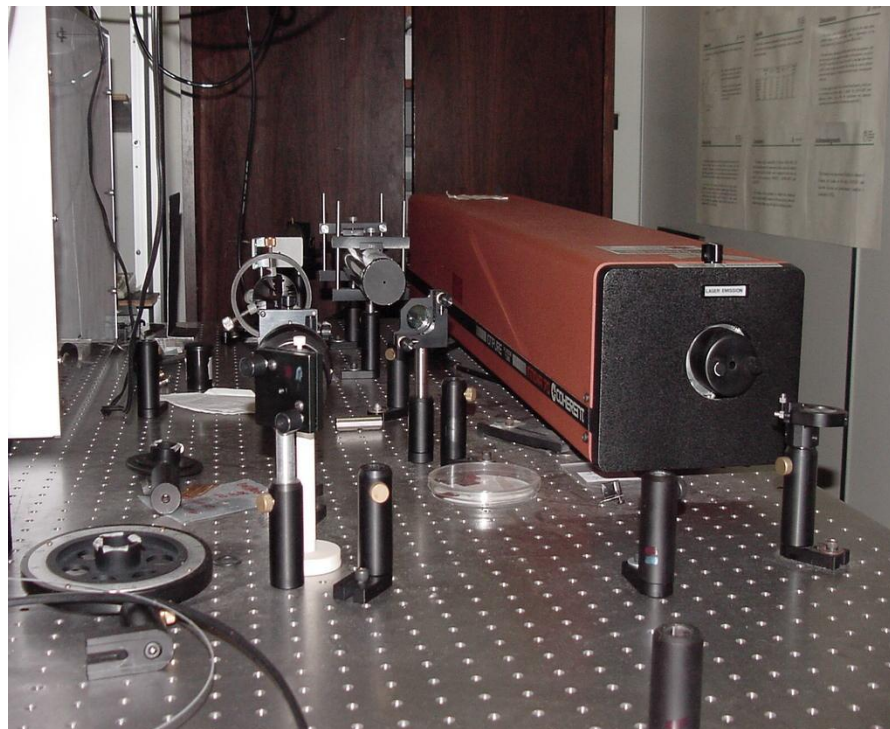
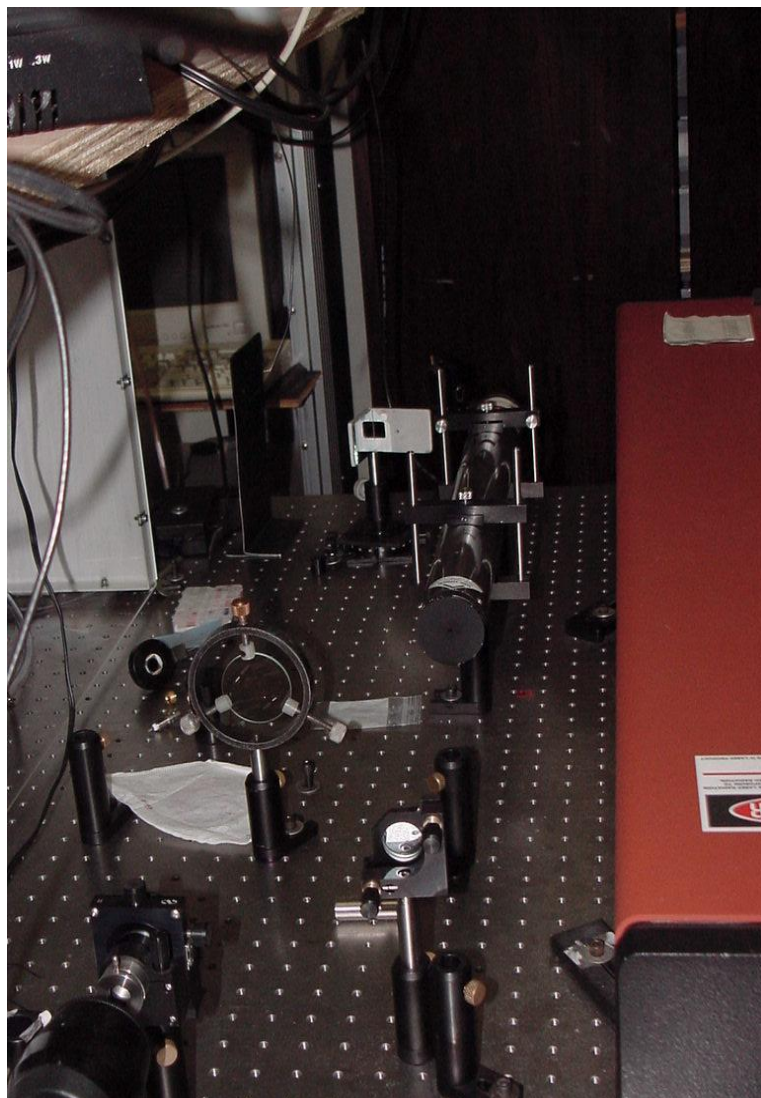


The purpose of the Photonics Group is to develop fundamental science and applied technology *in Optics and Photonics*

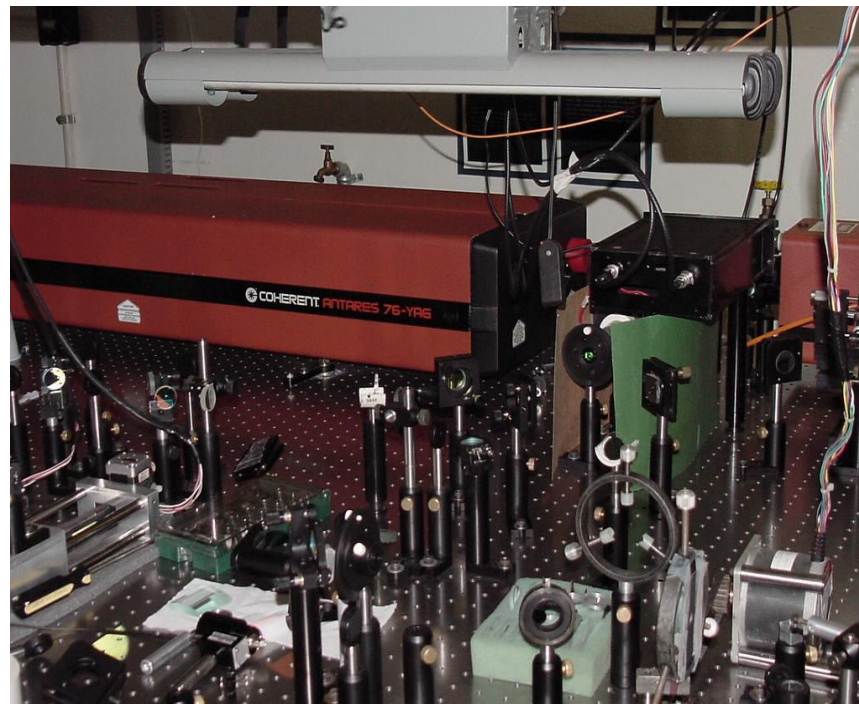
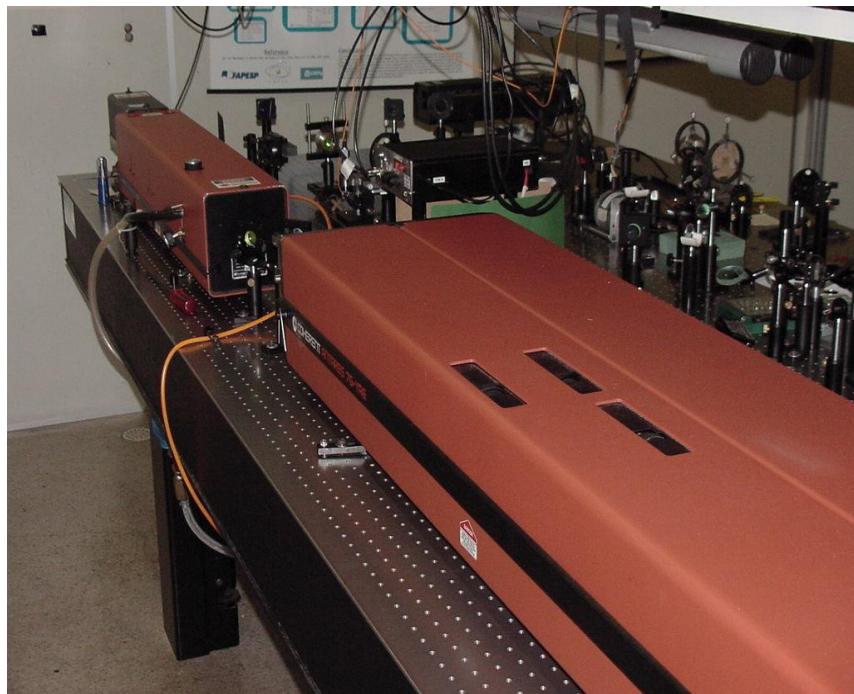
Some of the research areas

- Nonlinear optics
- Coherent control of light matter interaction
- fs-laser microfabrication and micromachining
- Optical spectroscopy
- Optical storage

cw laser



Picosecond laser



- ✓ *Laser Nd:YAG*
- Qswitched/modelocked*
 - ✓ *532nm and 1064 nm*
 - ✓ *100 ps*

150-fs amplified laser system



- ✓ Amplificador Ti:safira
 - ✓ 775 nm
 - ✓ 150 fs
 - ✓ 800 μJ

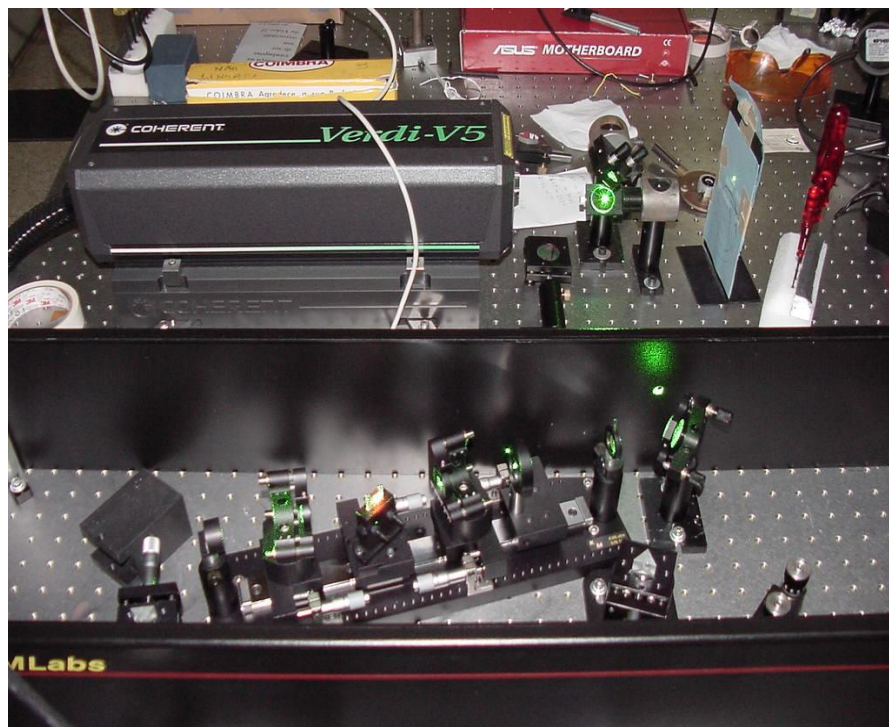
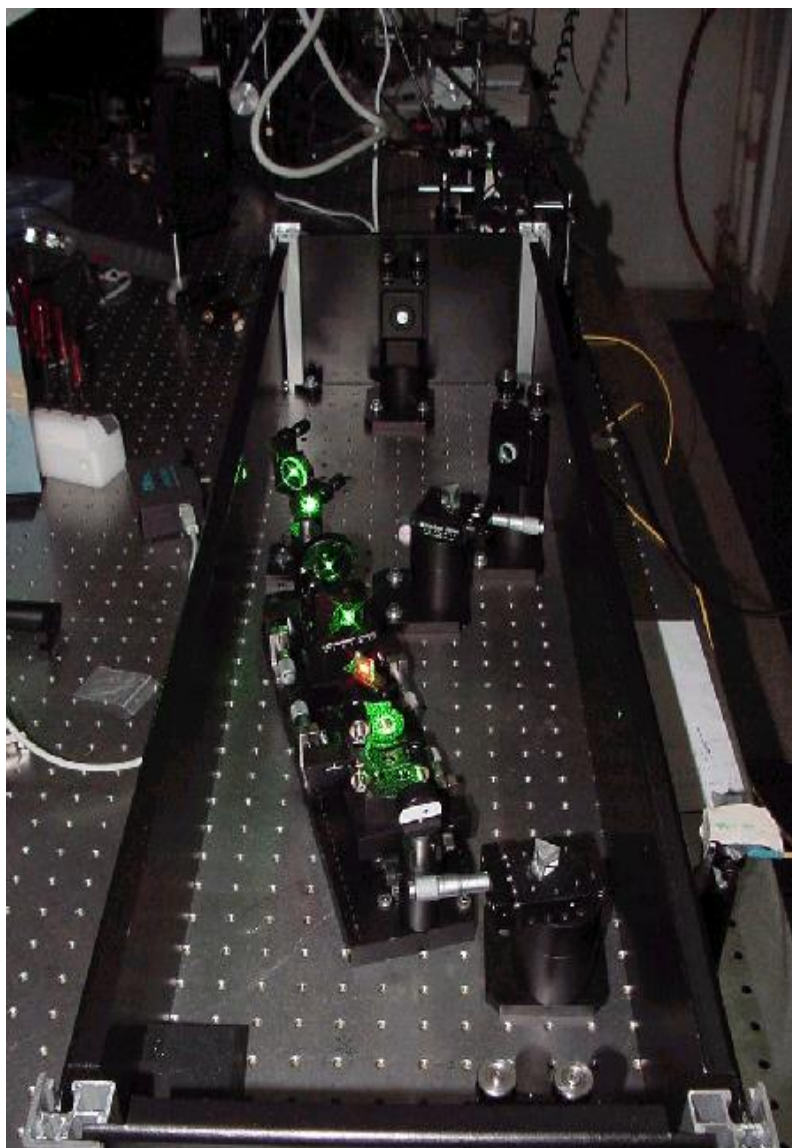
Optical Parametric Amplifier



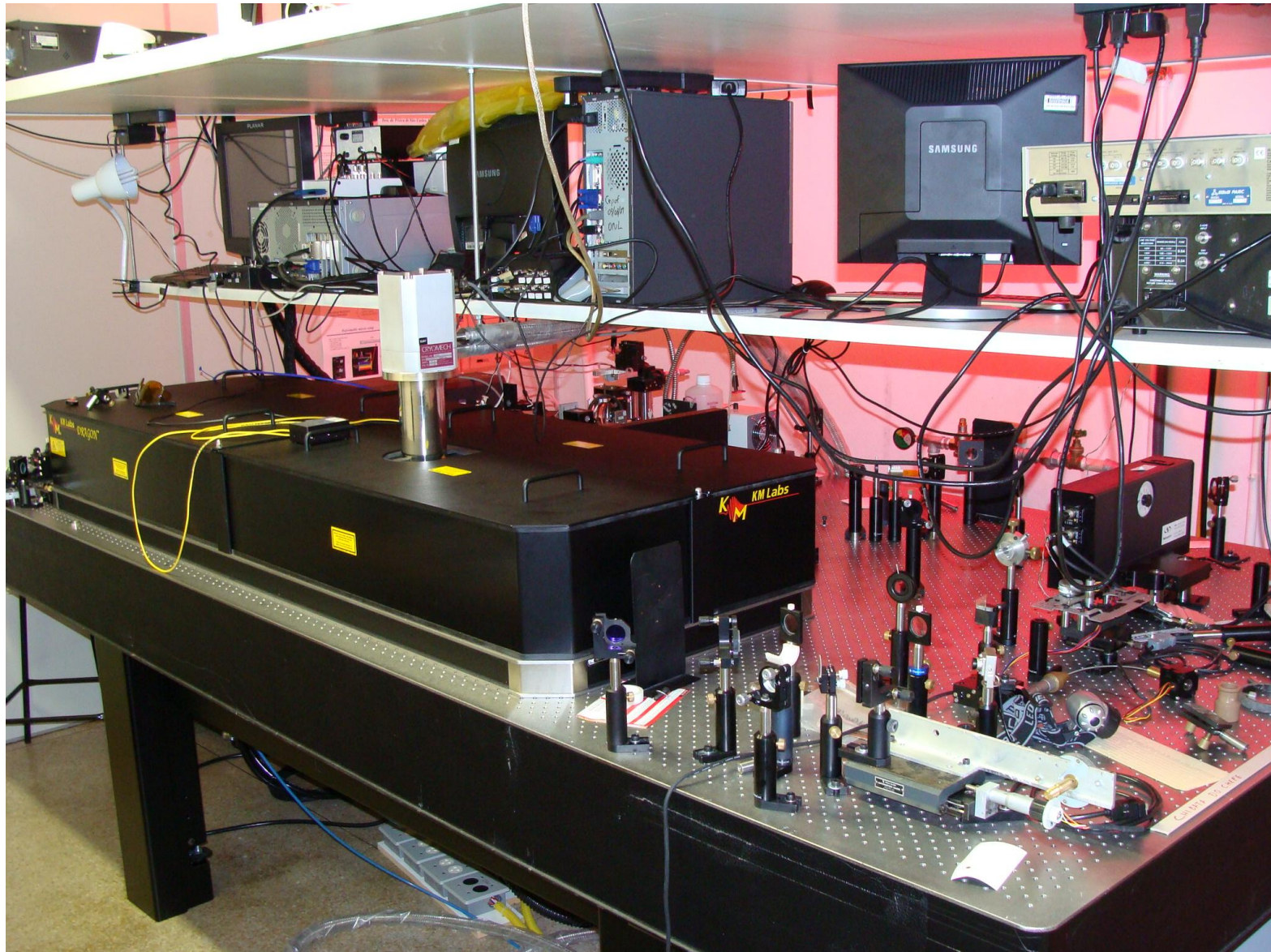
TOPAS

- ✓ pumped - Laser Clark
 - ✓ 460 - 2600 nm
 - ✓ ≈ 120 fs
 - ✓ 20-60 μ J

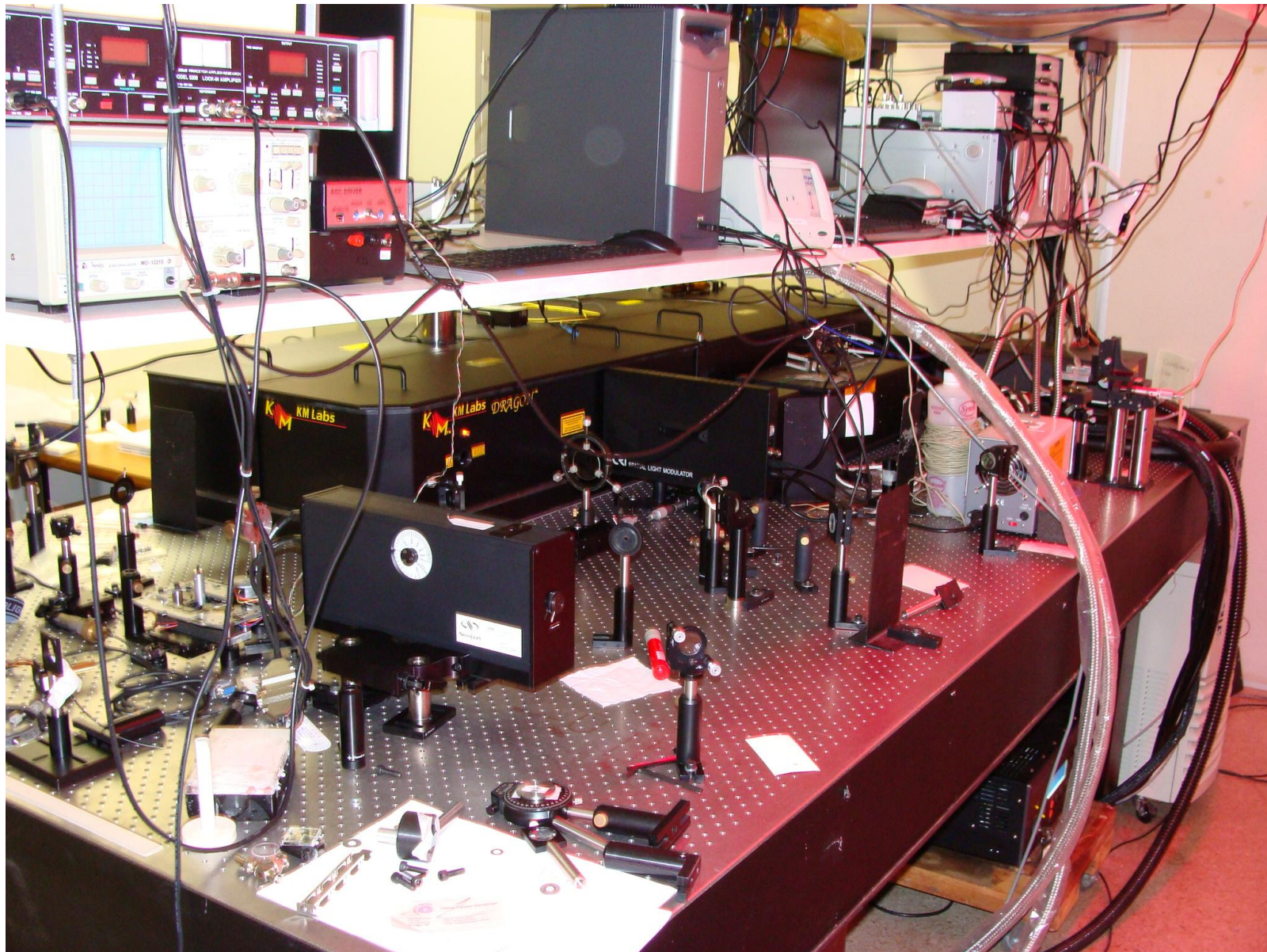
15-fs laser



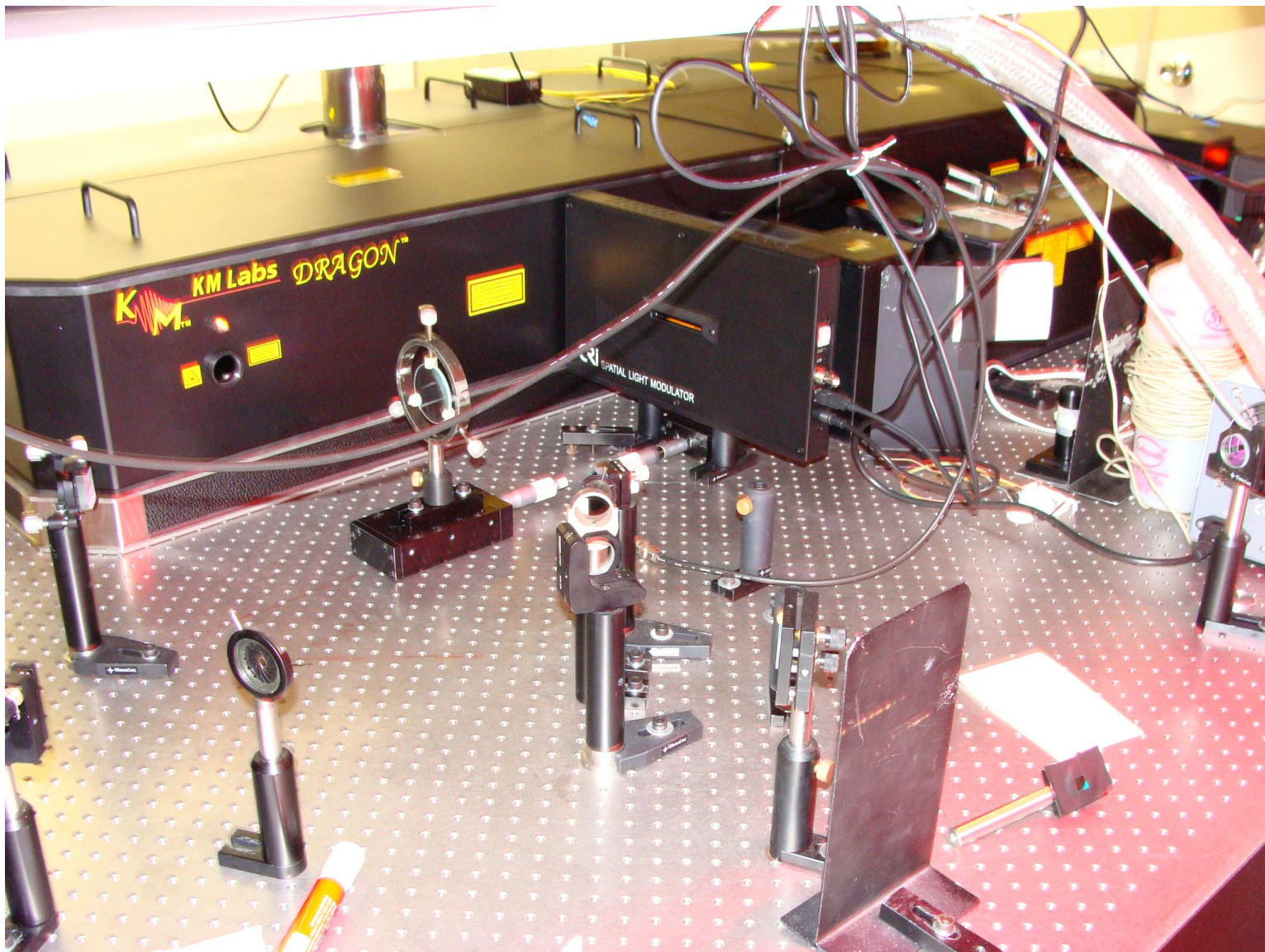
40-fs amplified laser system



40-fs amplified laser system



40-fs amplified laser system



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Outline

introduction to nonlinear optics
nonlinear optics in organic materials
experimental methods
examples of some results

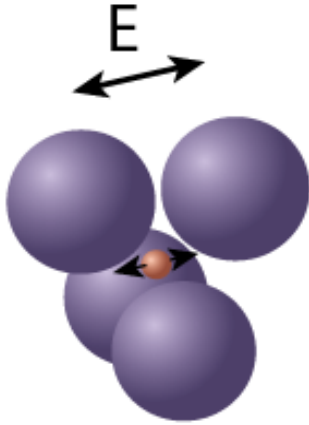
Linear optics vs Nonlinear optics

Optics is a branch of physics that describes the behavior and properties of light and the interaction of light with matter. Explains optical phenomena.

Nonlinear Optics

The branch of optics that describes optical phenomena that occur when very intense light is used

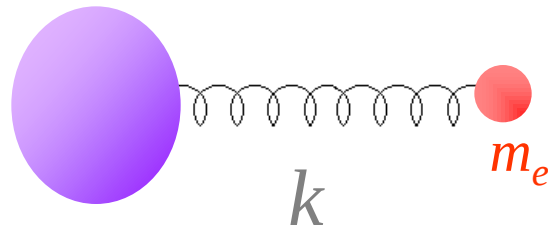
Linear optics



$$E_{\text{rad.}} \ll E_{\text{inter.}}$$

harmonic oscillator

electron on a spring



oscillation frequency

$$\omega_0 = \sqrt{\frac{k}{m_e}}$$

Linear optics

electron on a spring

$$F_{binding} = -m_e \omega_o^2 x$$

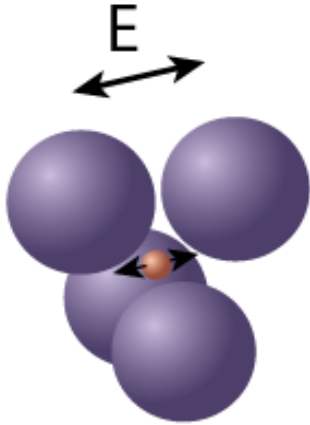
$$F_{damping} = -m_e \gamma \frac{dx}{dt}$$

$$F_{driving} = -eE = -eE_o e^{-i\omega t}$$

equation of motion

$$m \frac{d^2 x}{dt^2} = \sum F$$

Linear optics



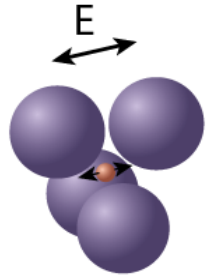
harmonic oscillator

$$m \frac{d^2x}{dt^2} + m\gamma \frac{dx}{dt} + m\omega_o^2 x = -eE$$

Steady state: electron oscillates at driving frequency

$$x(t) = x_o e^{-i\omega t} \quad x_o = -\frac{e}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o$$

Linear optics



Oscillating dipole

$$p(t) = -e x(t) = \frac{e^2}{m} \frac{1}{(\omega_o^2 - \omega^2) - i\gamma\omega} E_o e^{-i\omega t}$$

Polarization oscillator

$$P(t) = \left(\frac{Ne^2}{m} \right) \sum_j \frac{f_j}{(\omega_j^2 - \omega^2) - i\gamma_j\omega} E(t) \equiv \epsilon_o \chi_e E(t)$$

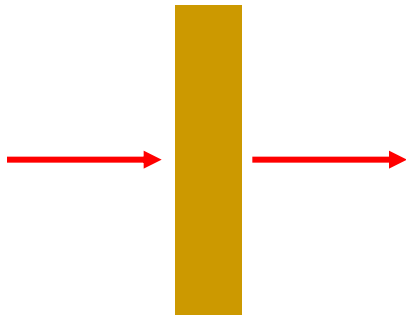
linear response

$$P = \chi E$$

Linear optical process

absorption

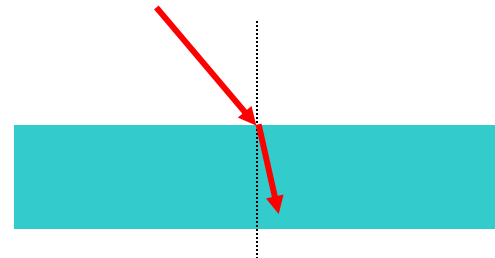
α_0 does not depend on light intensity



absorption of 10 %

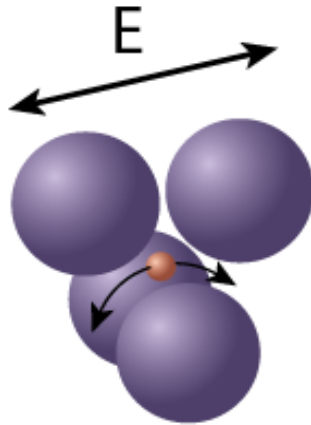
refraction

n_0 does not depend on light intensity



index of refraction 1.3

Nonlinear optics



anharmonic oscillator

high light intensity

$$E_{\text{rad.}} \sim E_{\text{inter.}}$$

nonlinear polarization response

$$P = \chi^{(1)}E + \chi^{(2)}E^2 + \chi^{(3)}E^3 + \dots$$

Nonlinear optics

nonlinear expansion of the polarization

$$\vec{P} = \chi^{(1)} \cdot \vec{E} + \chi^{(2)} : \vec{E}\vec{E} + \chi^{(3)} : \vec{E}\vec{E}\vec{E} + \dots$$



**linear
processes**



SHG

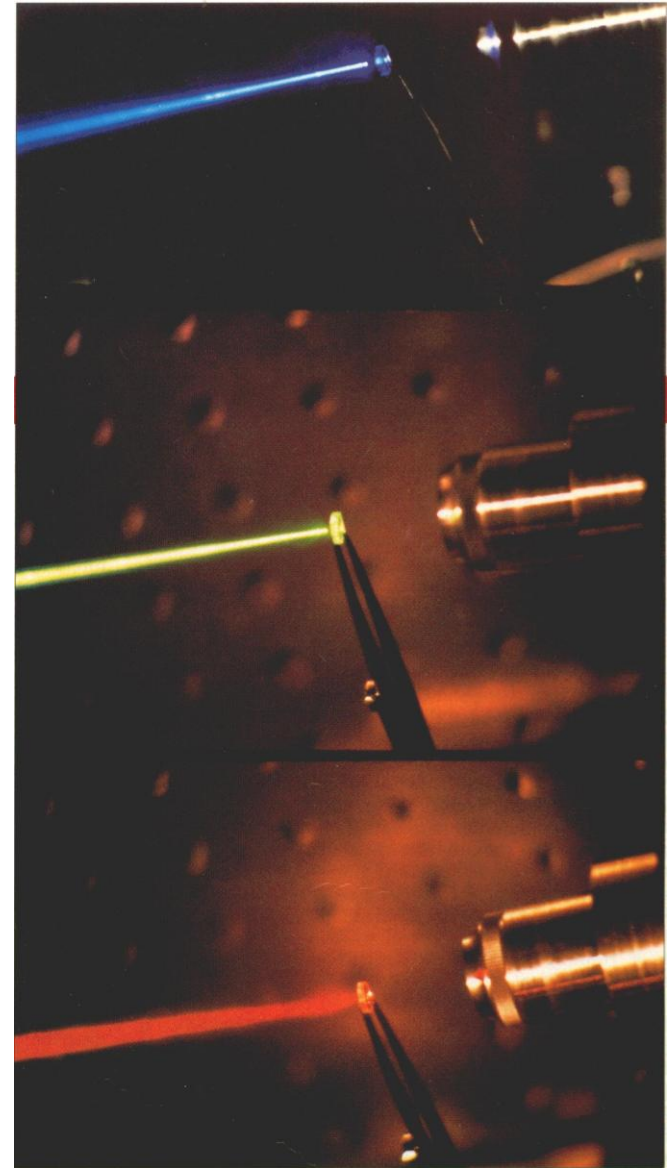
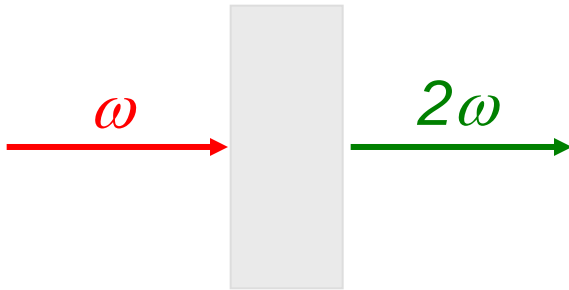


**THG
Kerr effect**

Nonlinear Optics

Second order processes $\chi^{(2)}$

Second Harmonic Generation



Nonlinear Optics

in medium with inversion symmetry

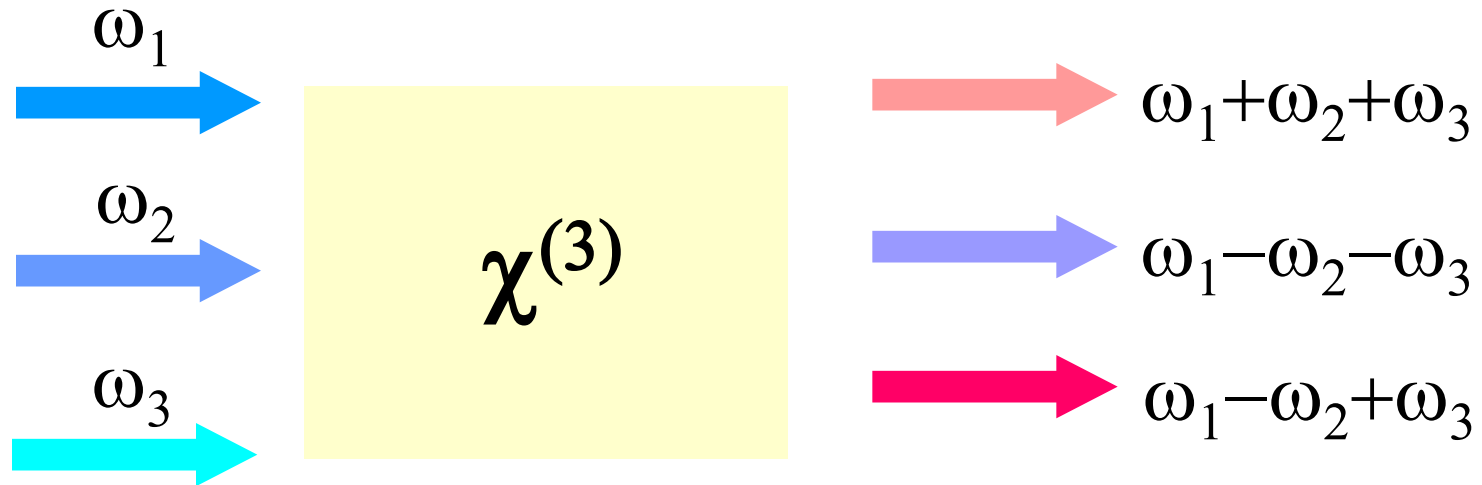
$$\vec{P}^{(2)} = \vec{\chi}^{(2)} : \vec{E} \vec{E} \quad \Rightarrow \quad -\vec{P}^{(2)} = \vec{\chi}^{(2)} : (-\vec{E})(-\vec{E})$$

and consequently

$$\chi^{(2)} = -\chi^{(2)} = 0$$

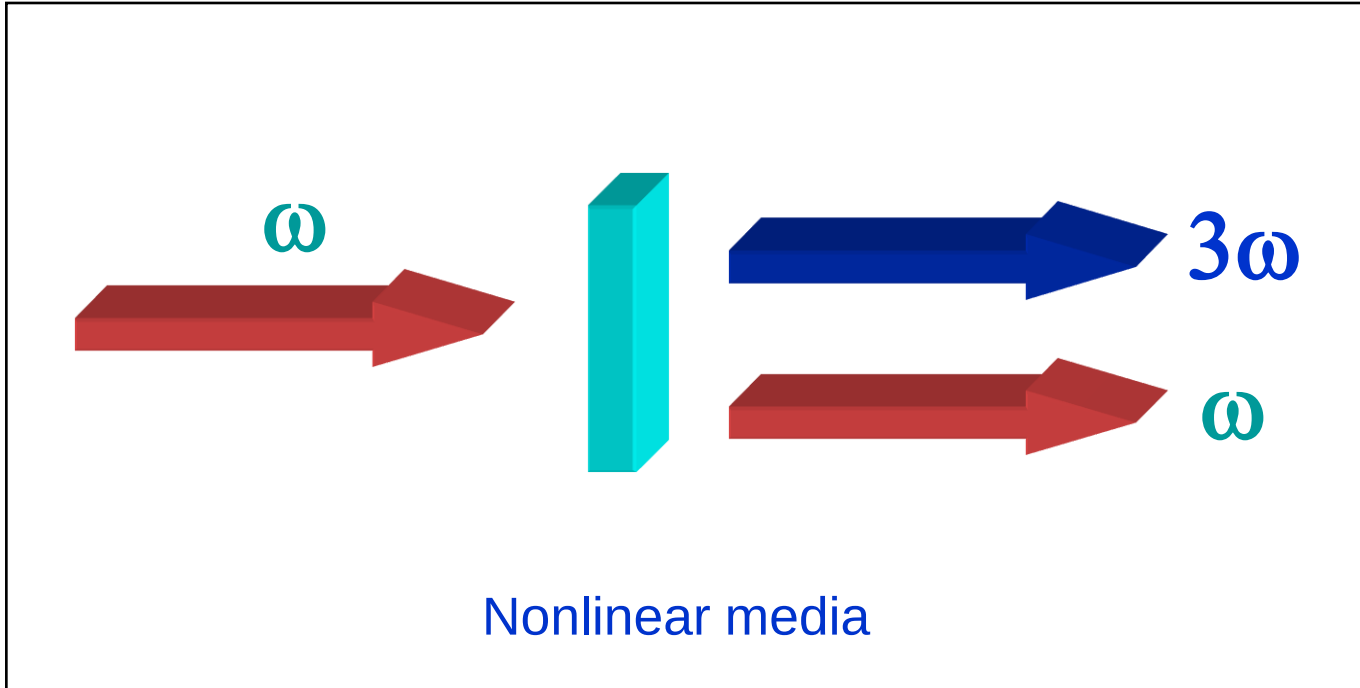
Nonlinear Optics

Third order processes $\chi^{(3)}$



- If $\omega_1 = \omega_2 = \omega_3$
 - Third harmonic generation: 3ω
 - Self phase modulation: ω

Third Harmonic Generation $\chi^{(3)}$



Nonlinear Optics

$$\chi^{(2)} = 0$$



$\chi^{(3)}$ Third order processes

Nonlinear polarization

$$P = \chi^{(1)}E + \chi^{(3)}E^3 + \dots$$

Third order polarization

$$P^{(3)}(t) = \chi^{(3)}E(t)E^*(t)E(t) = \chi^{(3)}I(t)E(t)$$

Nonlinear Optics

consequently

$$P = P^{(1)} + P^{(3)} = (\chi^{(1)} + \chi^{(3)}I)E \equiv \chi_{eff}E$$

and

$$n = \sqrt{\epsilon} = \sqrt{1 + \chi_{eff}} \approx \sqrt{1 + \chi^{(1)}} + \frac{1}{2} \frac{\chi^{(3)}I}{\sqrt{1 + \chi^{(1)}}} = n_o + n_2I$$

Kerr media

$$n = n_0 + n_2I$$

Nonlinear Optics

Third order processes $\chi^{(3)}$

$$n_2 \approx \chi^{(3)}$$

Kerr media:

$$n = n_0 + n_2 I$$

Index of refraction depends on the light intensity

Self phase modulation

Kerr media:

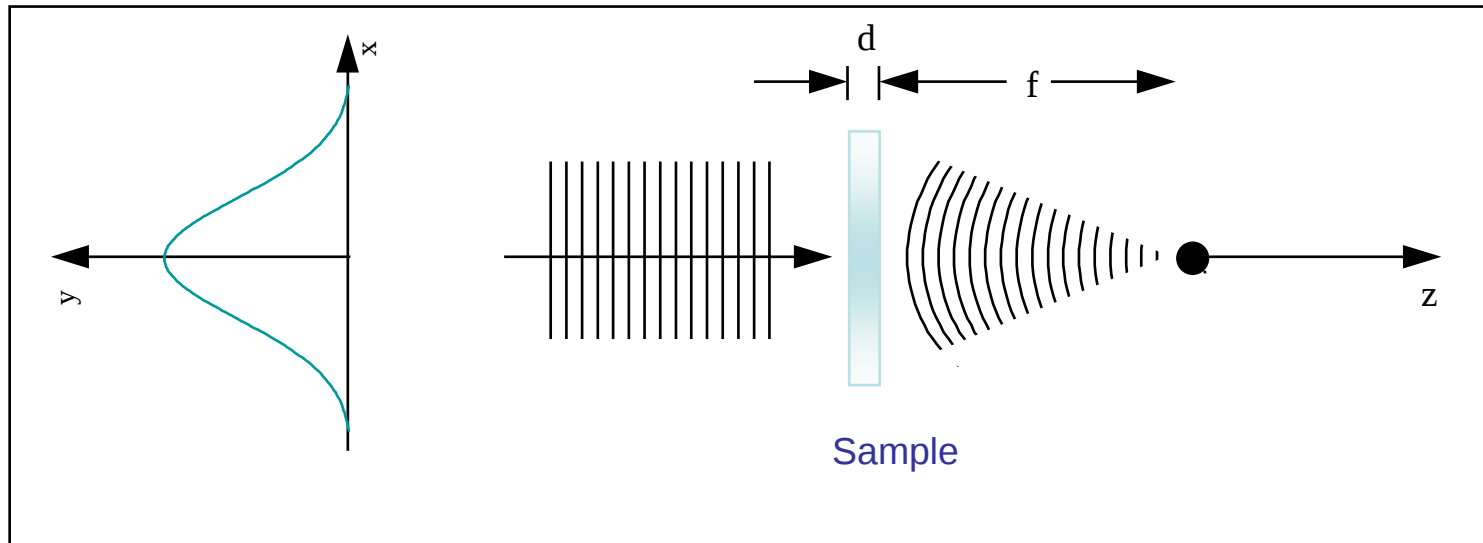
centre symmetric: $\chi^{(2)} = 0$

$$n = n_0 + n_2 I$$

$$P_{NL} = \chi^{(3)} E^3$$

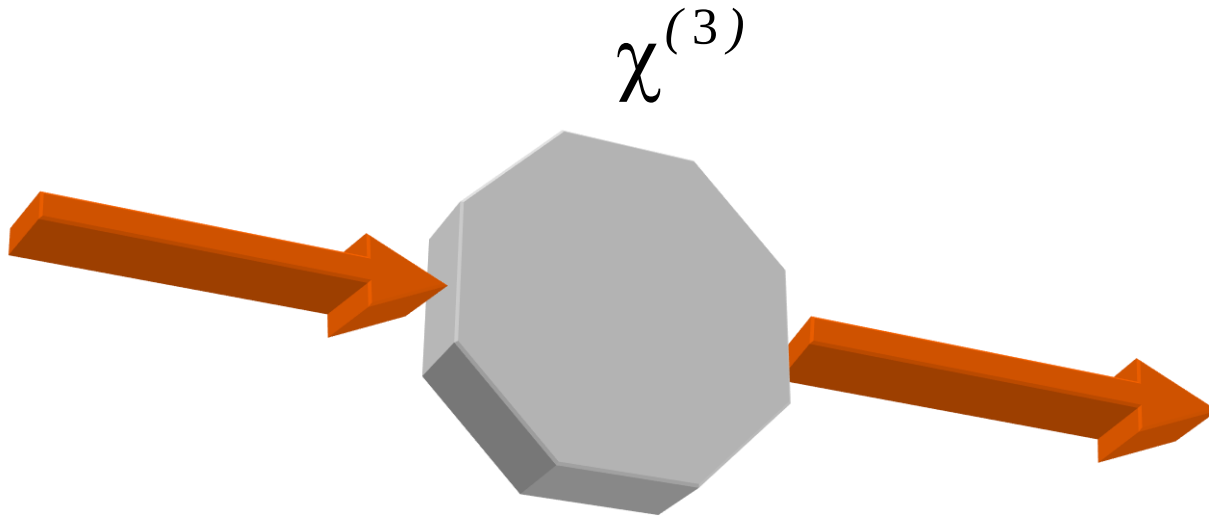
$$n_2 > 0$$

Material behaves as a convergent lens



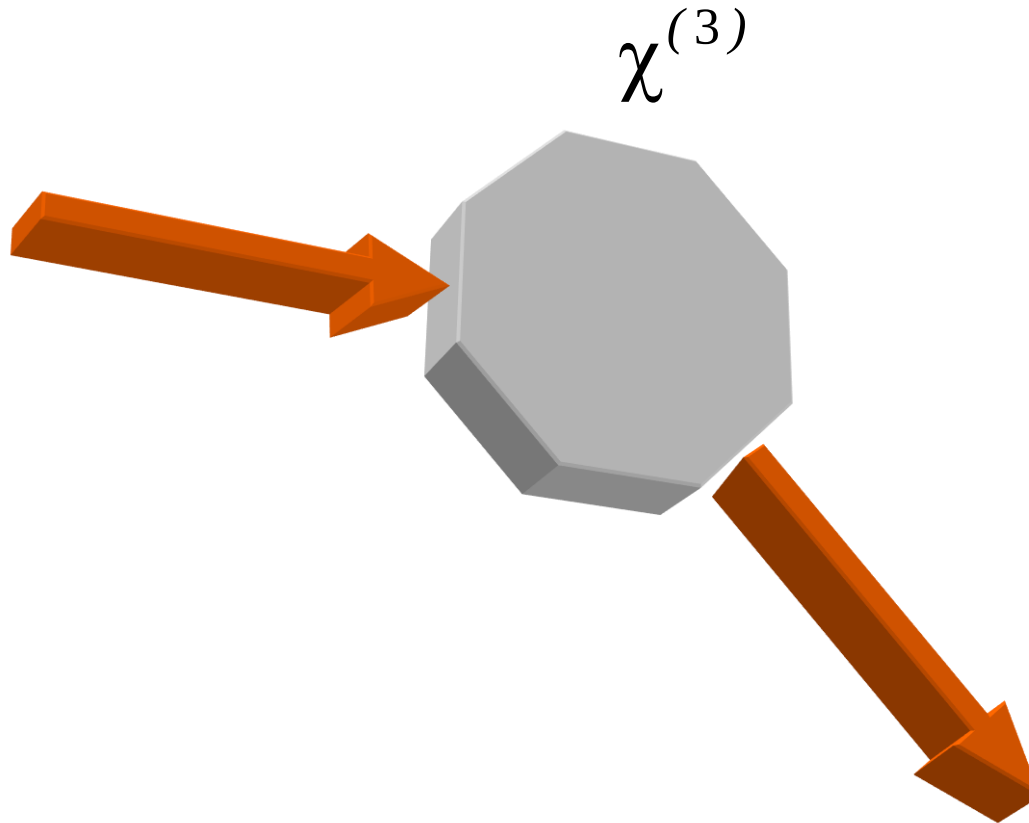
Optical switching

- low intensity



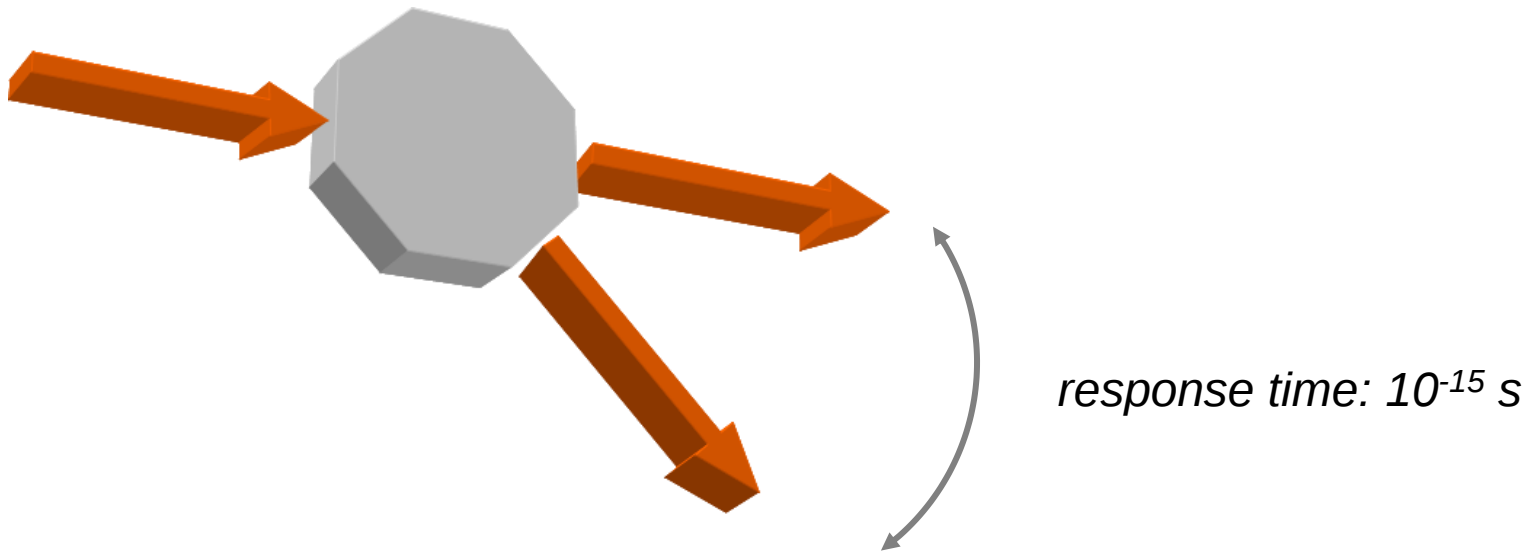
Optical switching

- high intensity



Self action process

Optical switching

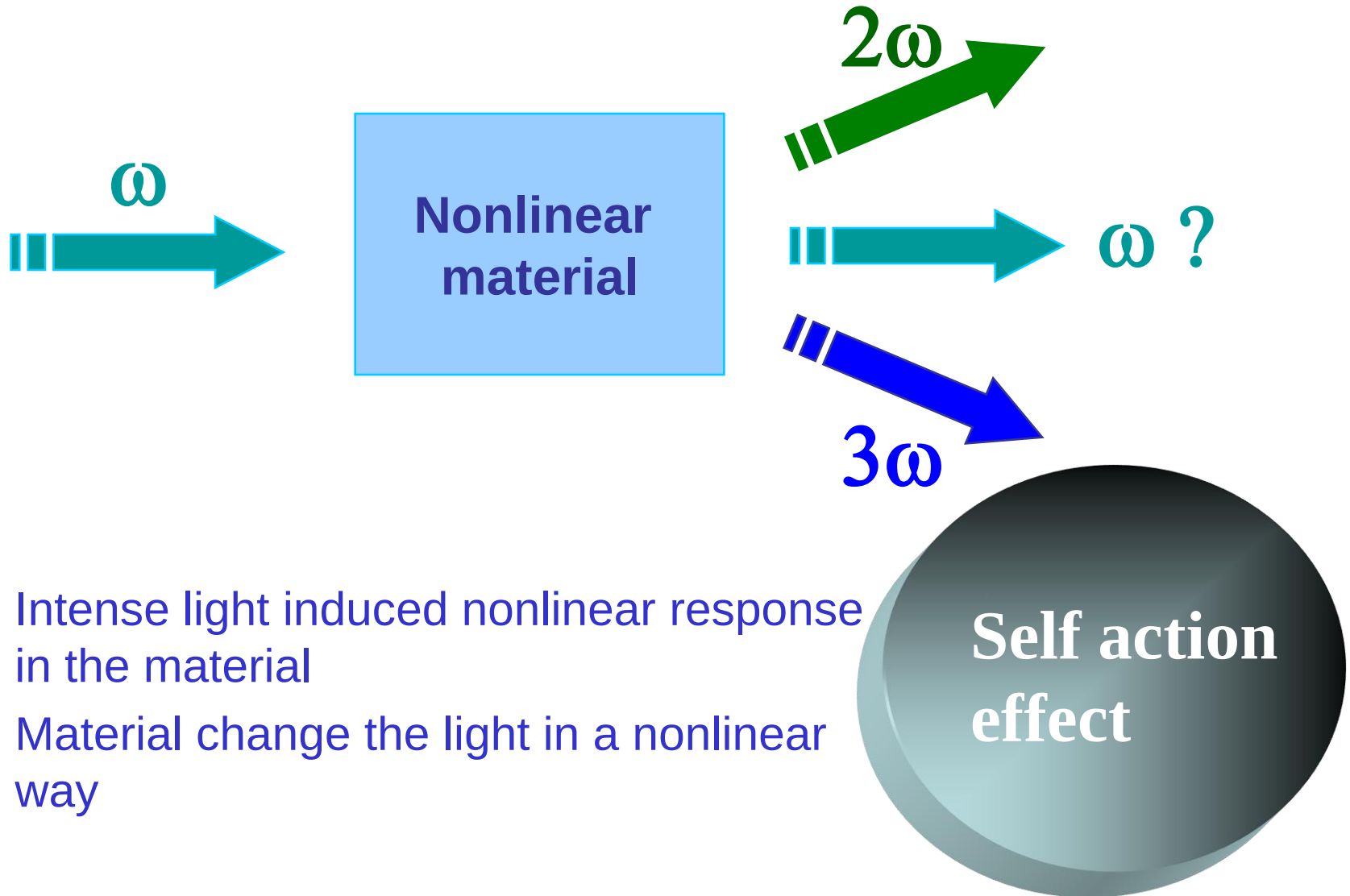


response time 1×10^{-9} 1×10^{-15}

1GHz → 1 THz

1 million times faster

Nonlinear optics



- Intense light induced nonlinear response in the material
- Material change the light in a nonlinear way

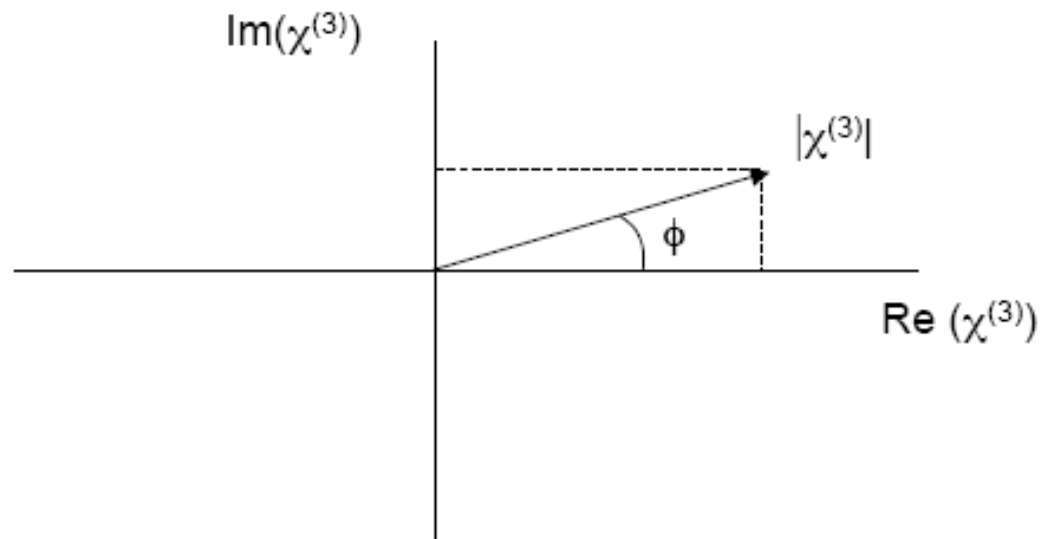
Nonlinear Optics

$\chi^{(3)}$ is a complex quantity

$$\chi^{(3)} = \text{Re}(\chi^{(3)}) + i \text{Im}(\chi^{(3)})$$

Related to intensity
dependent refractive index

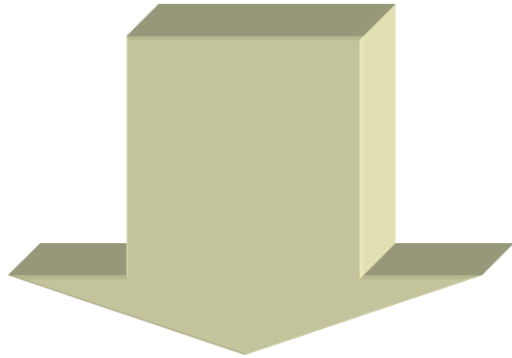
Related to two-photon
absorption



Third order processes: $\chi^{(3)}$

Refractive process:

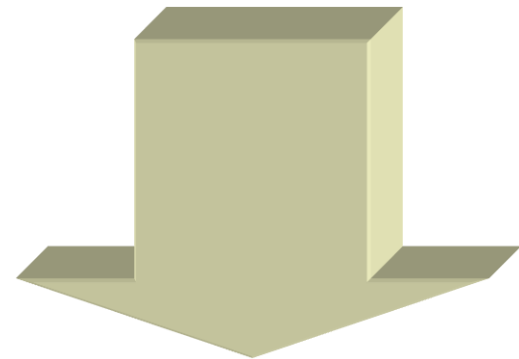
$$n = n_0 + n_2 I$$



- self-phase modulation
- lens-like effect

Absorptive process:

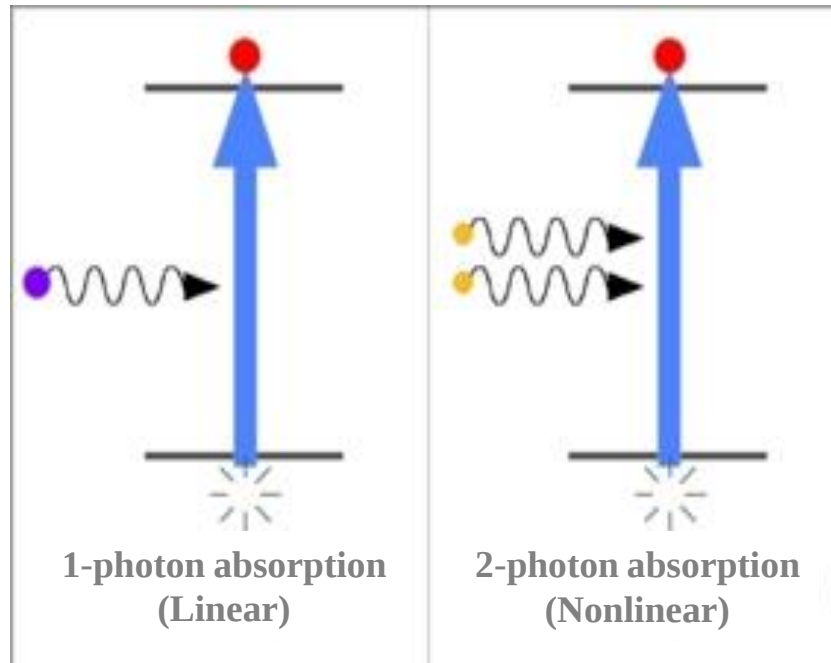
$$\alpha = \alpha_0 + \beta I$$



- nonlinear absorption
- two-photon absorption

Two-photon absorption (2PA) process

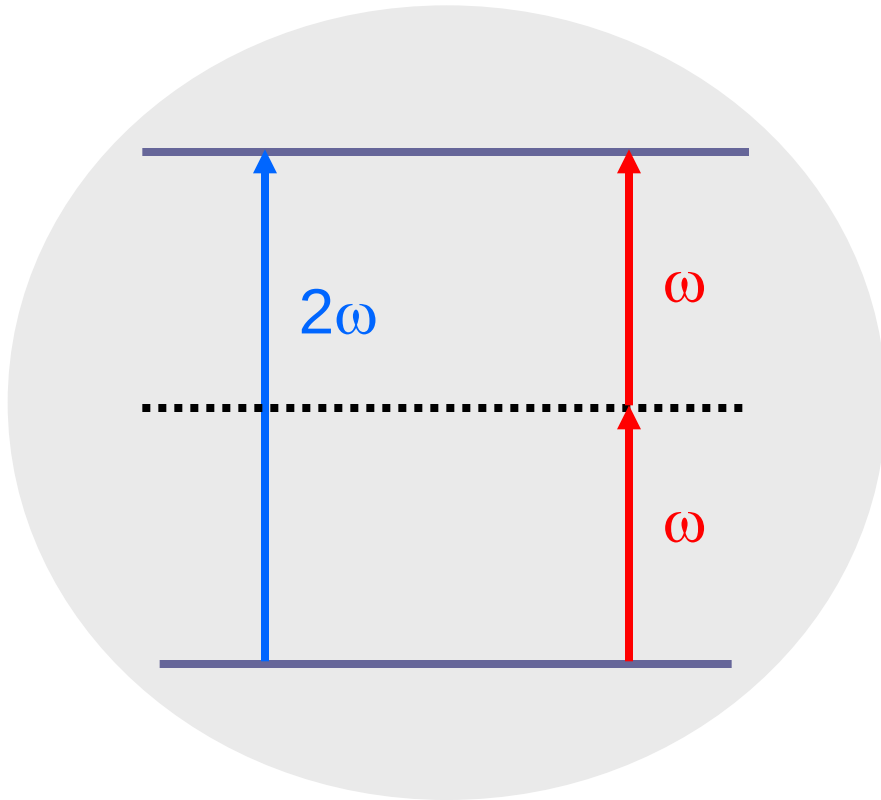
Phenomenon does not described for the Classical Physics and **does not observed until the development of the Laser.**



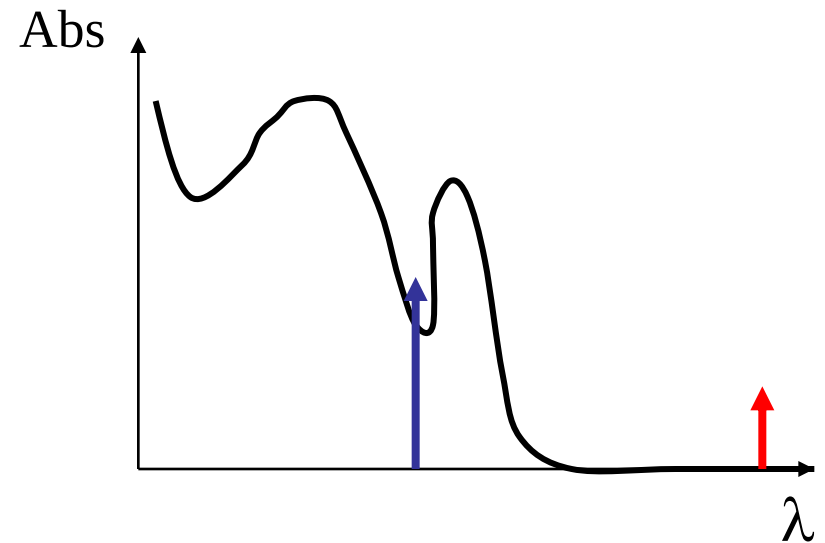
Theoretical model: Maria Göppert-Mayer, 1931

Two photons from an intense laser light beam are simultaneously absorbed in the same “quantum act”, leading the molecule to some excited state with energy equivalent to the absorbed two photons.

two-photon absorption



$$\alpha = \alpha_0 + \beta I$$



Applications:

optical limiting

fluorescence microscopy

microfabrication

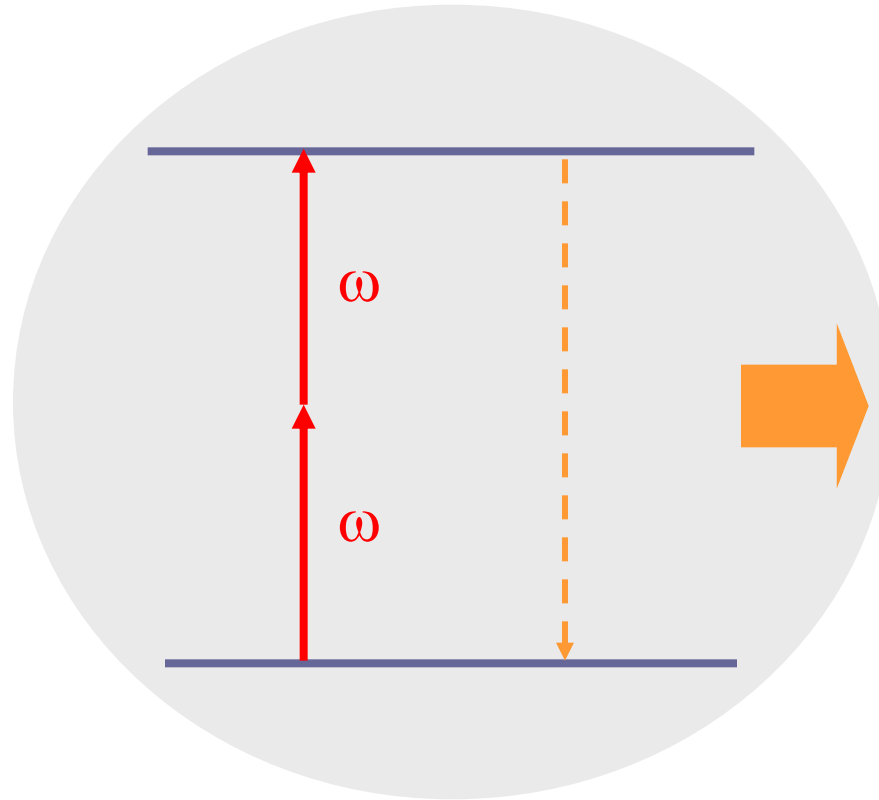
optical limiting



To protect eye and sensors from intense laser pulses

two-photon fluorescence

$$\alpha = \alpha_0 + \beta I$$

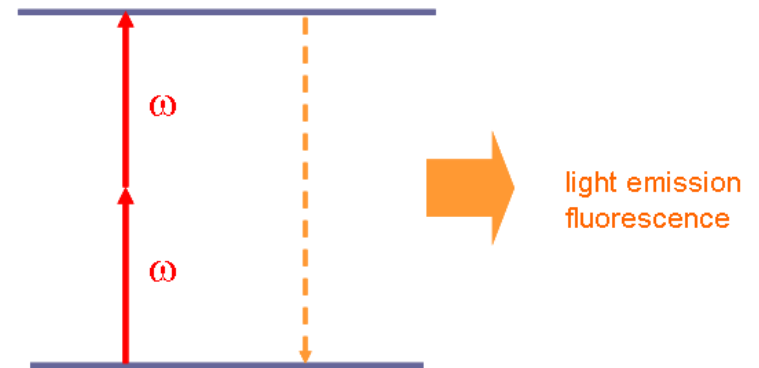
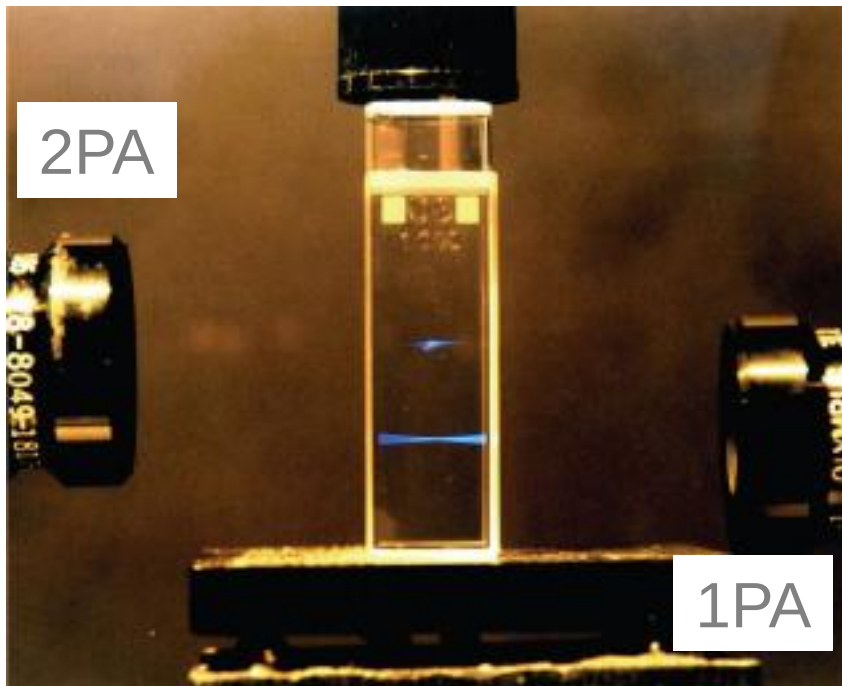


light emission
fluorescence

$$\text{TPA rate constant} \propto \delta I^2$$

localization of the excitation with 2PA

dilute solution of fluorescent dye



$$\begin{aligned} TPA &\propto \delta I^2 \\ I &\sim \frac{1}{z^2} \\ \Rightarrow TPA &\sim \frac{1}{z^4} \end{aligned}$$

spatial confinement of excitation

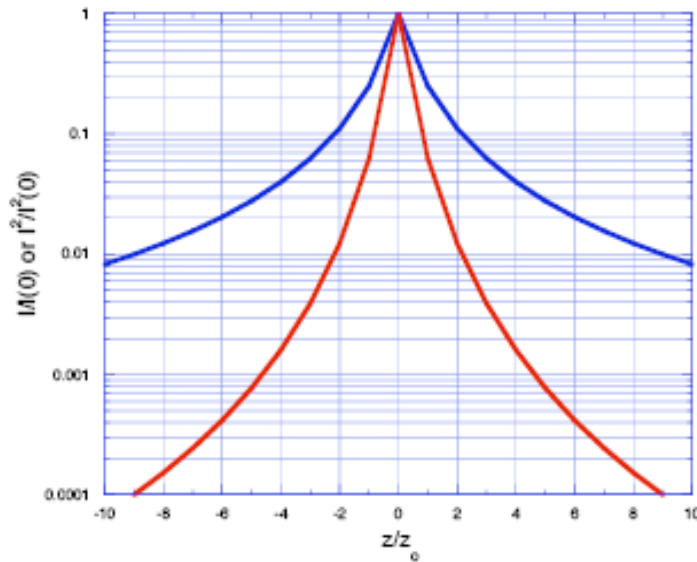
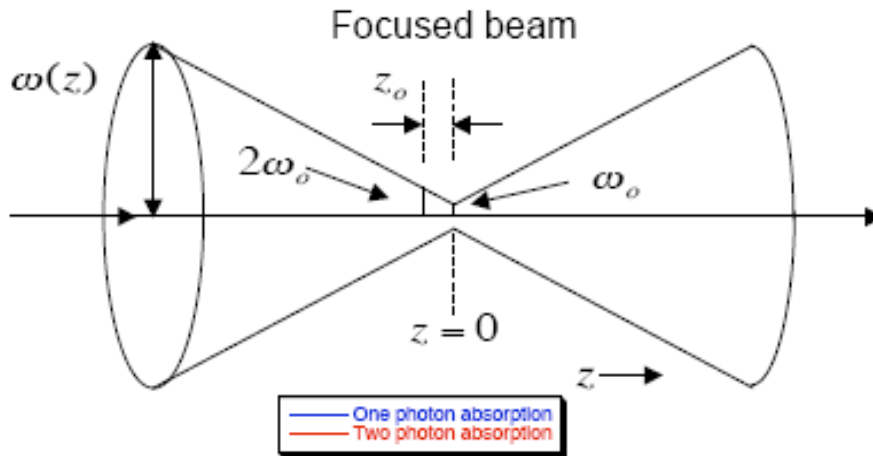
excitation profile along z

Radius, area and intensity of focused beam

$$\omega(z) = \omega_0 \left(1 + \frac{z}{z_0} \right)$$

$$A(z) = \pi \omega^2 = \pi \omega_0^2 \left(1 + \frac{z}{z_0} \right)^2$$

$$I(z) = \frac{E}{At} \propto \frac{1}{A} = \frac{1}{\pi \omega_0^2 \left(1 + \frac{z}{z_0} \right)^2}$$



Normalized excitation rates
(ignoring beam attenuation)

one photon

$$\frac{I(z)}{I(0)} = \left(1 + \frac{z}{z_0} \right)^{-2}$$

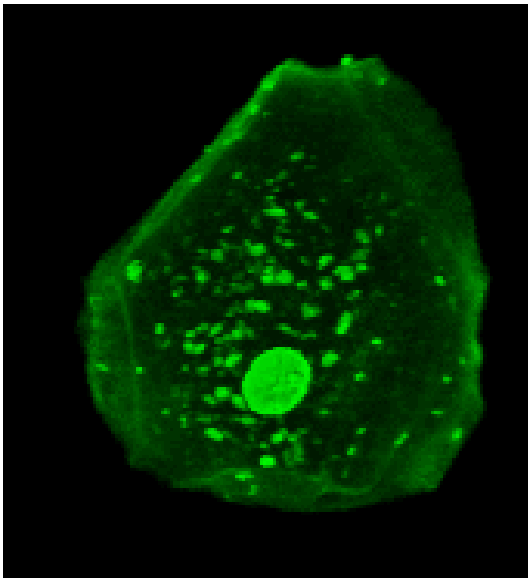
two photon

$$\frac{I(z)}{I(0)} = \left(1 + \frac{z}{z_0} \right)^{-4}$$

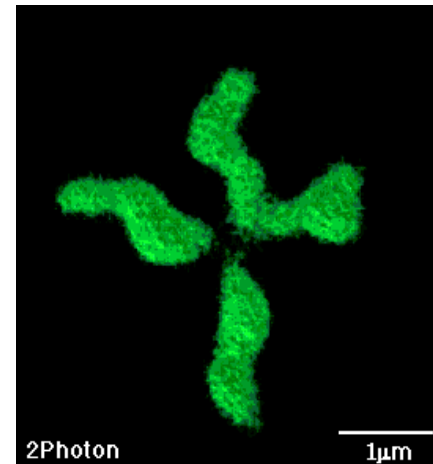
two-photon fluorescence microscopy

➤ Microscopy by two-photon fluorescence

3D image of a cell

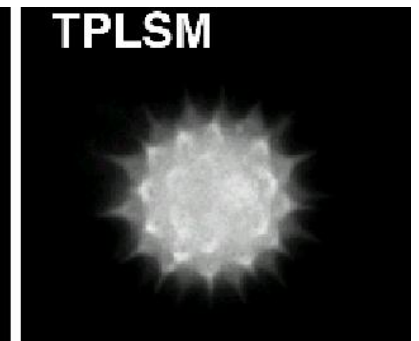
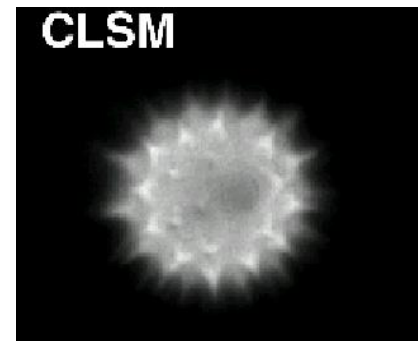
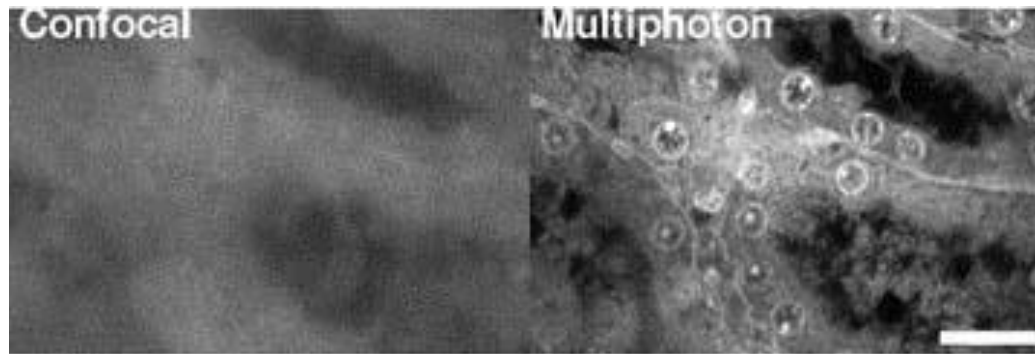


*Laboratory for Optics and
Biosciences
Ecole polytechnique*



Human chromosome

Two-photon fluorescence microscopy

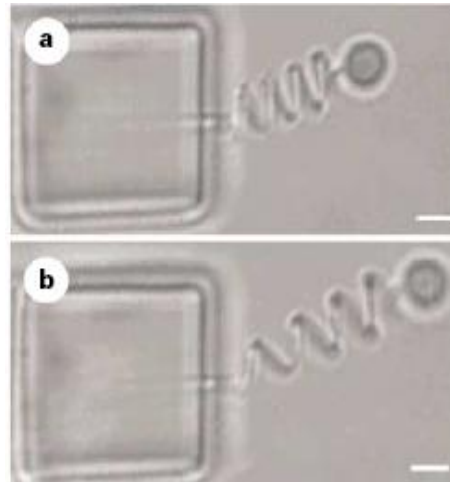


Fluorescent marker $\Rightarrow$ fluorophores

microfabrication

➤ Two-photon polymerization

Nature 412, 697-698 (2001)



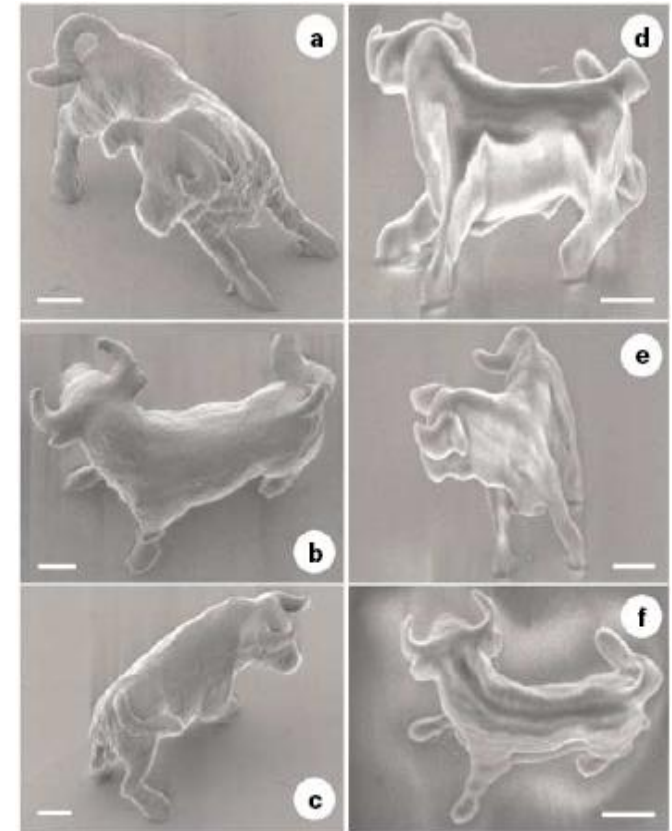
oscillator



Venus statue

Two-photon polymerization

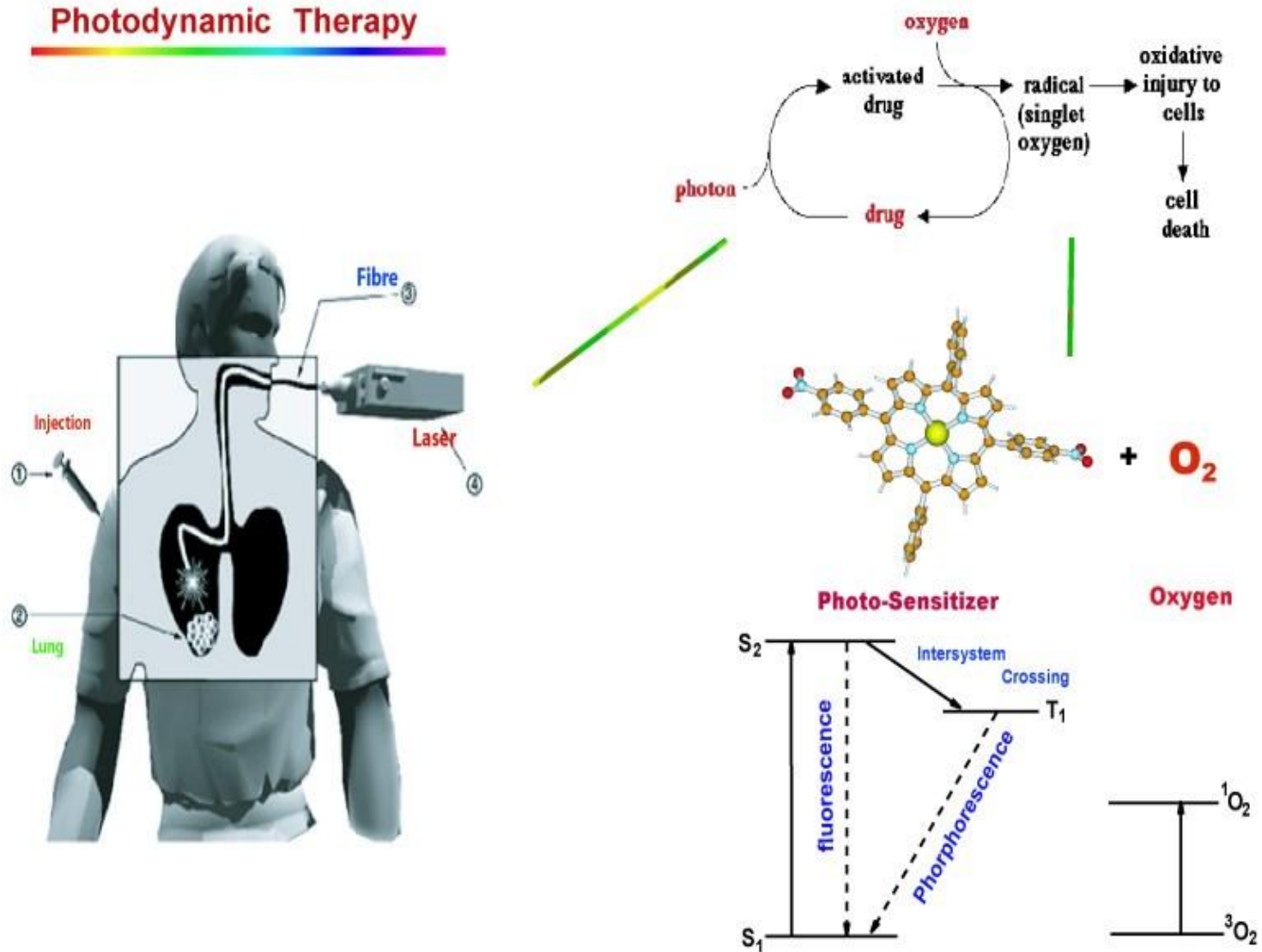
Opt. Exp. 12, 5521-5528 (2004)



Bull

two-photon photodynamic therapy

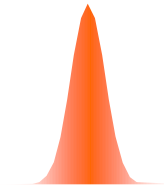
➤ PDT via dois fótons



Real applications in nonlinear optics

Very intense light: femtosecond pulses

Ti:Sapphire lasers



100 fs



50 fs



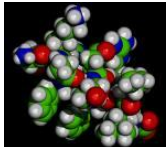
20 fs

$1 \text{ fs} = 10^{-15} \text{ s}$

1 fs



1s



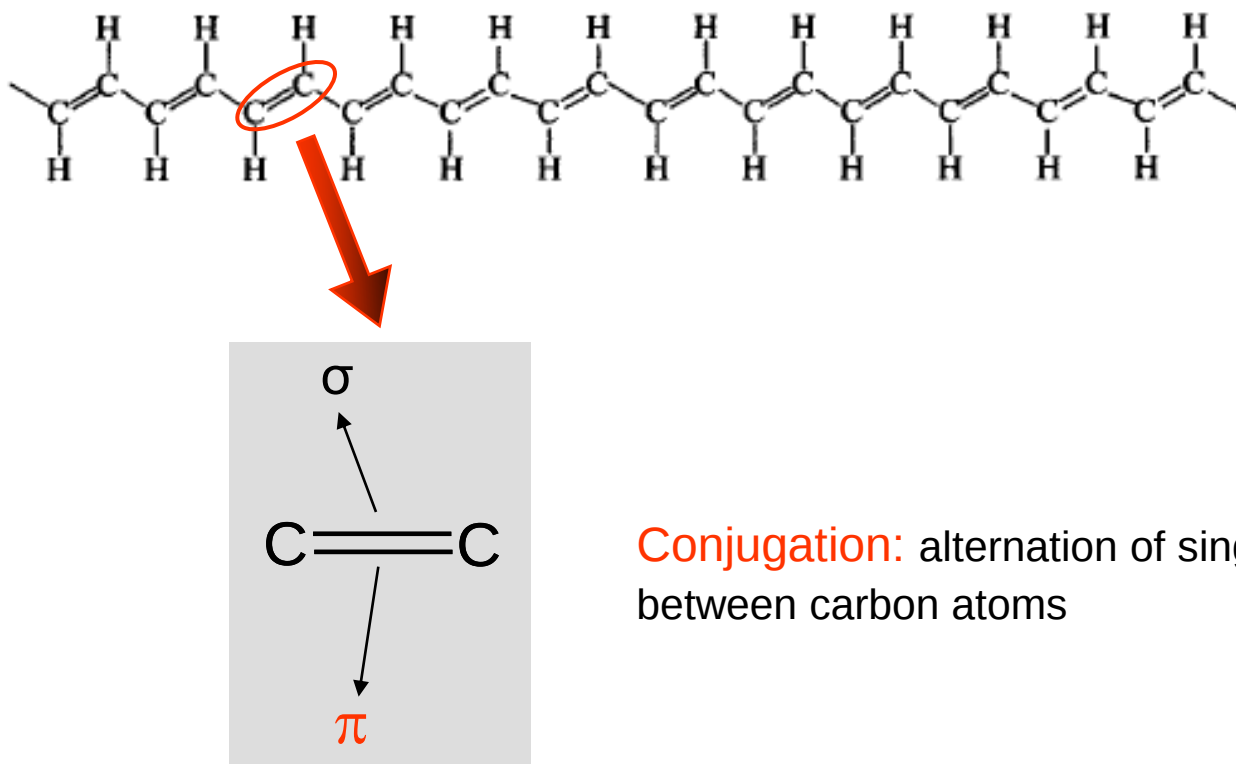
Laser intensities $\sim 100 \text{ GW/cm}^2$
 $1 \times 10^{11} \text{ W/cm}^2$

Laser pointer: 1 mW/cm^2 ($1 \times 10^{-3} \text{ W/cm}^2$)



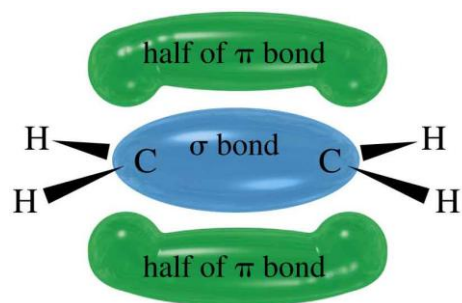
Organic materials

- Flexibility to tune the nonlinear optical response by manipulating the molecular structure
- π -conjugated structures



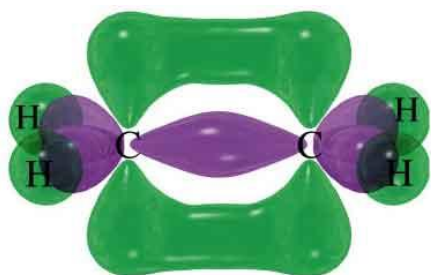
Conjugation: alternation of single and double bonds between carbon atoms

π -conjugation



σ bond: forms a strong chemical bond; localized

π bond: weaker bond; out of the C atoms axis



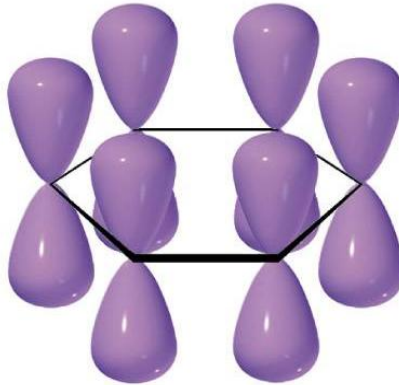
"Free electrons" that are easier to move under an applied electric field

π -conjugation

benzene



p-orbitals



π delocalization
(π -electron cloud)



π bond in conjugated system: delocalized electrons

high optical nonlinearities

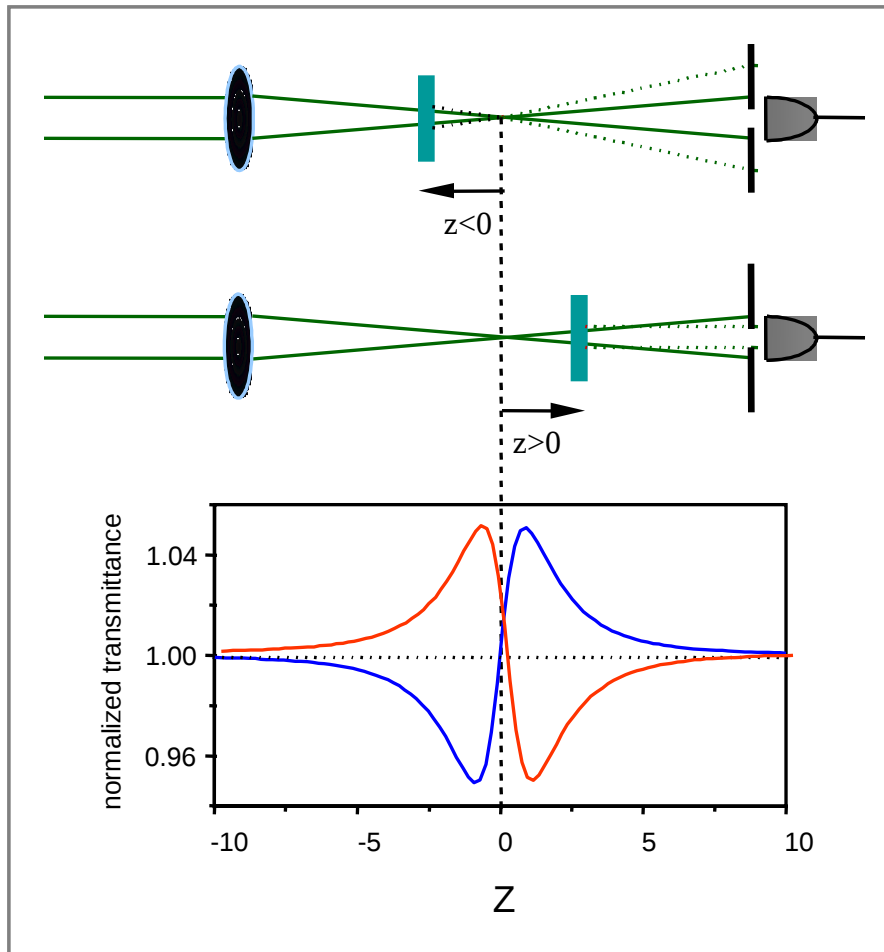


Research

- Understanding the physical principles behind two-photon absorption
- Understanding the relationship between molecular structure and two-photon absorption
- Developing molecules with high optical nonlinearities that can be used for application

Z-scan

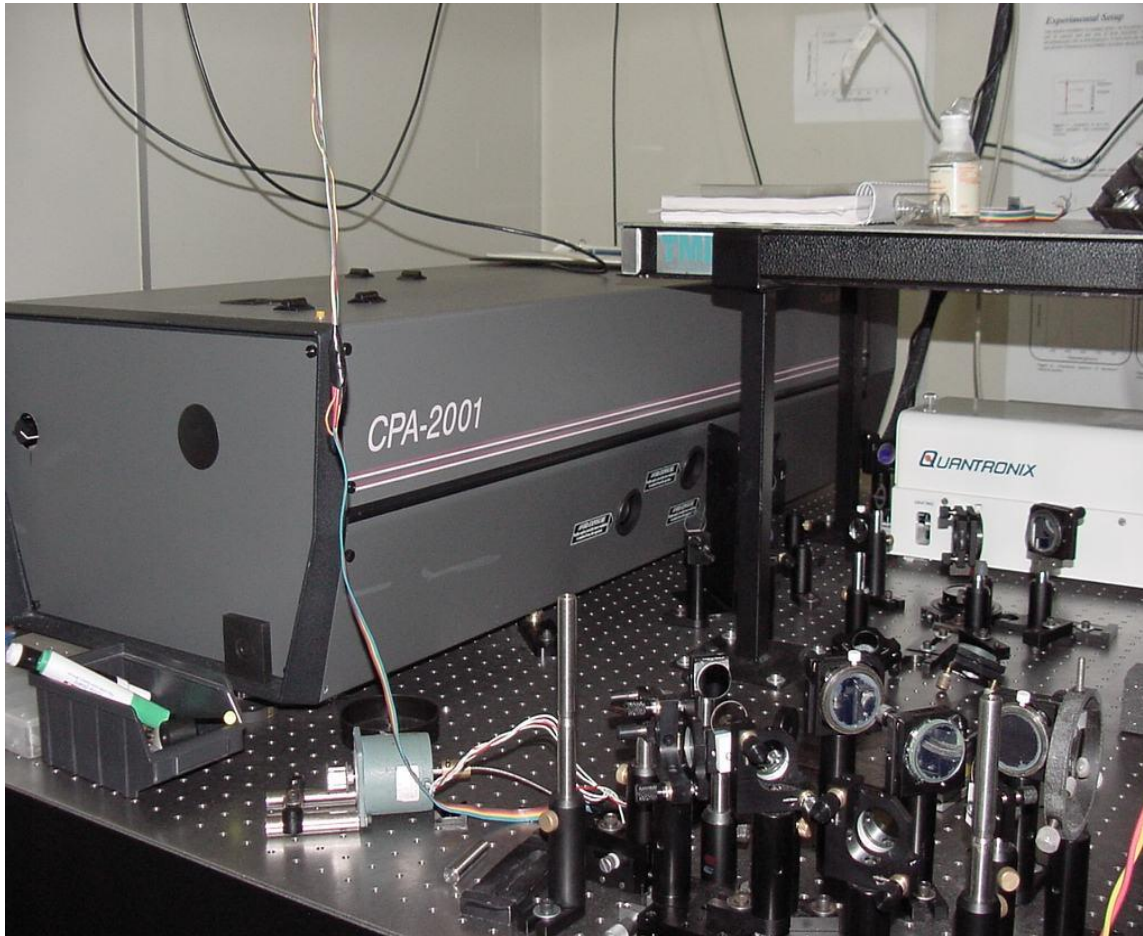
closed aperture Z-scan



$$n = n_0 + n_2 I$$

$$\Delta T \propto n_2 I$$

150 fs laser system



Ti:Sapphire amplifier

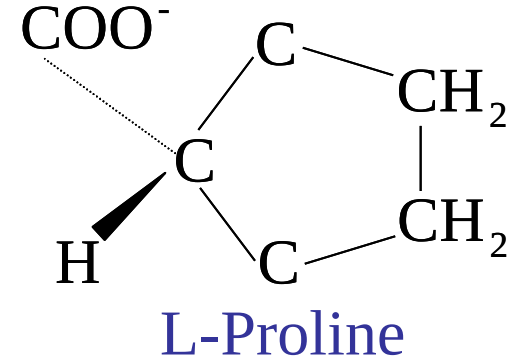
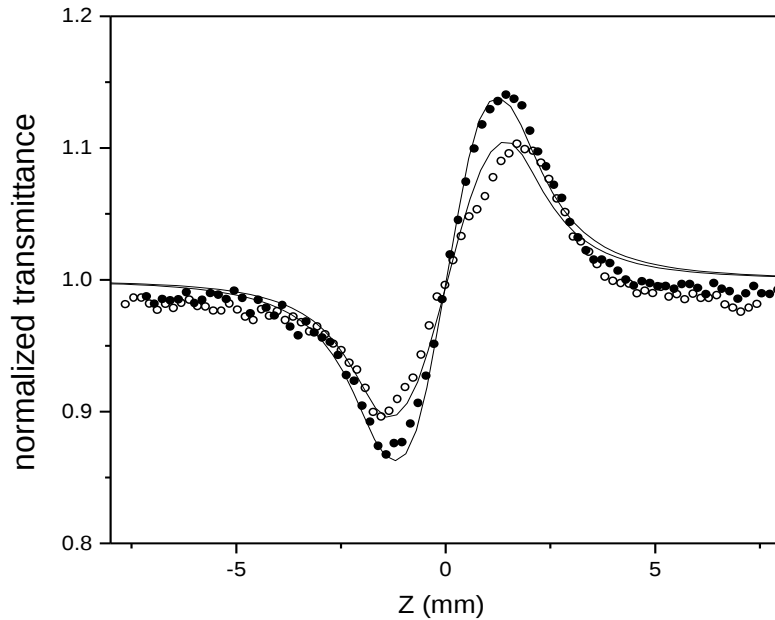
775 nm

150 fs

800 μ J

Nonlinear refraction

Aminoacids



$$\Delta T_{pv} = 0.406 \Delta \Phi_0 = 0.406 \left(\frac{2\pi}{\lambda} n_2 I_0 L \right)$$

150 fs 100 GW/cm²

775 nm

$$n_2 \propto \chi^{(3)}$$

microscopic to macroscopic nonlinearities

Molecular

$$\mu = \mu_0 + \alpha E + \beta E^2 + \gamma E^3 + \dots$$

Material

Electro-optic effect Intensity dependent refraction

$$P = P_0 + \chi^{(1)} E + \chi^{(2)} E^2 + \chi^{(3)} E^3 + \dots$$

For weakly interacting molecules

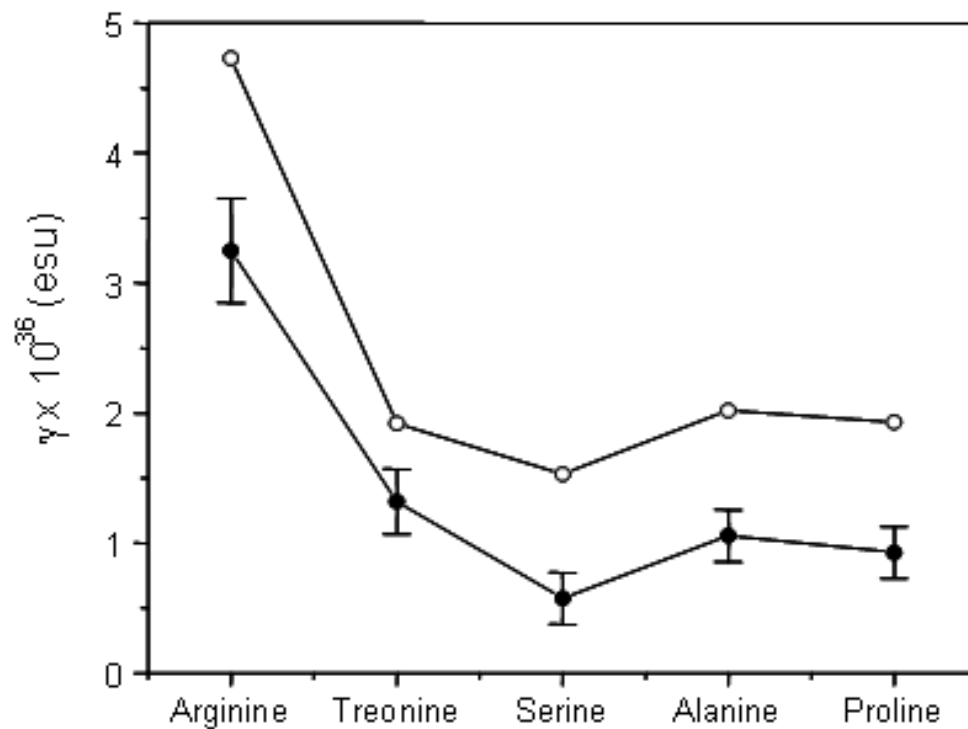
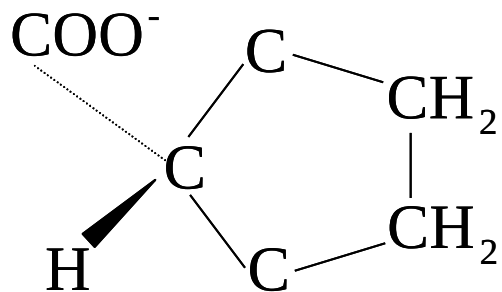
$$\chi^{(3)} = \gamma N L^4$$

Number density

Local field factor
 $(n^2 + 2)/3$

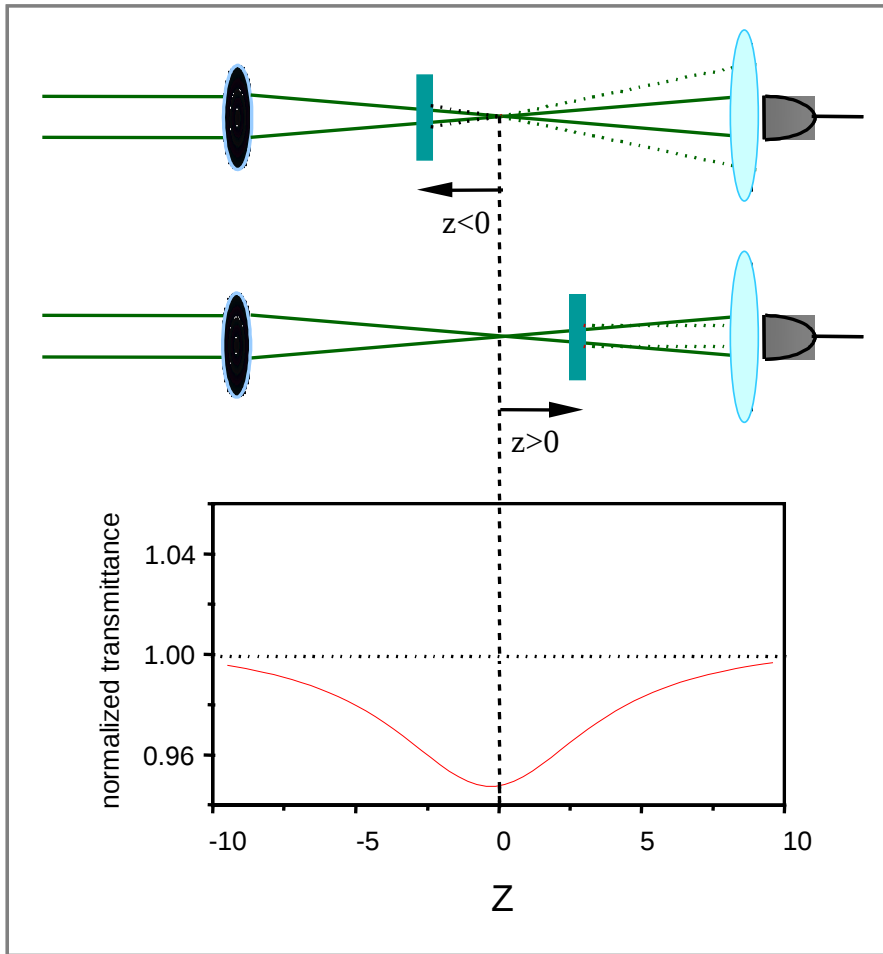
Nonlinear refraction

L-Proline



Z-scan (nonlinear absorption)

open aperture Z-scan



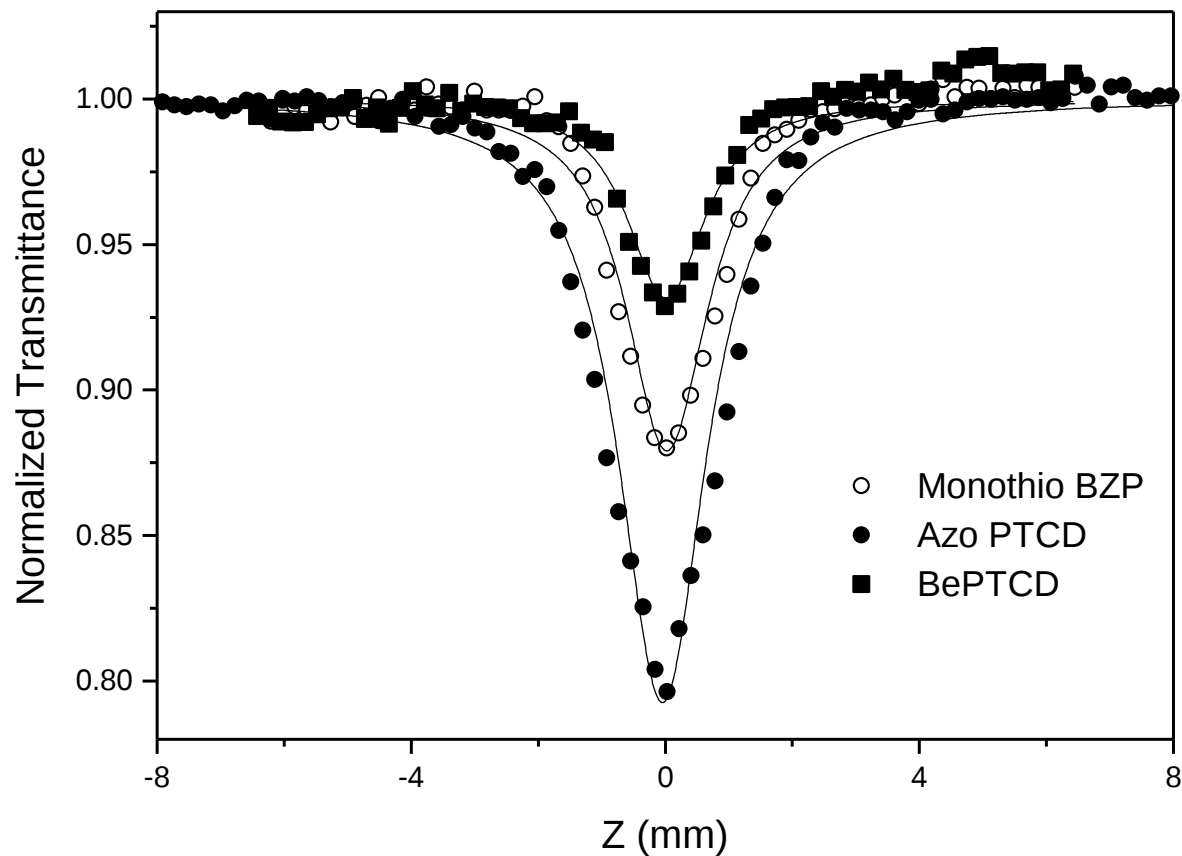
$$\alpha(I) = \alpha_0 + \beta I$$

$$\Delta T \propto \beta I$$

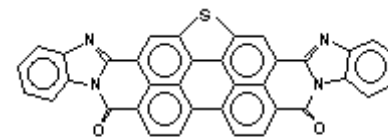
$$T(z) = \sum_{m=0}^{\infty} \frac{[-q_0(z,0)]^m}{(m+1)^{3/2}}$$

$$q_0(z,t) = \beta I_0 L / (1 + z^2 / z_0^2)$$

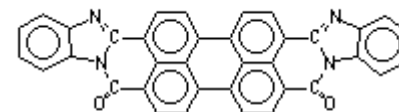
Nonlinear absorption



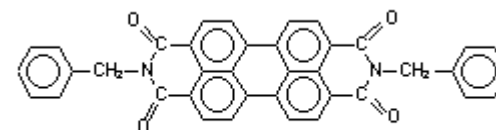
775 nm



Monothio BZP



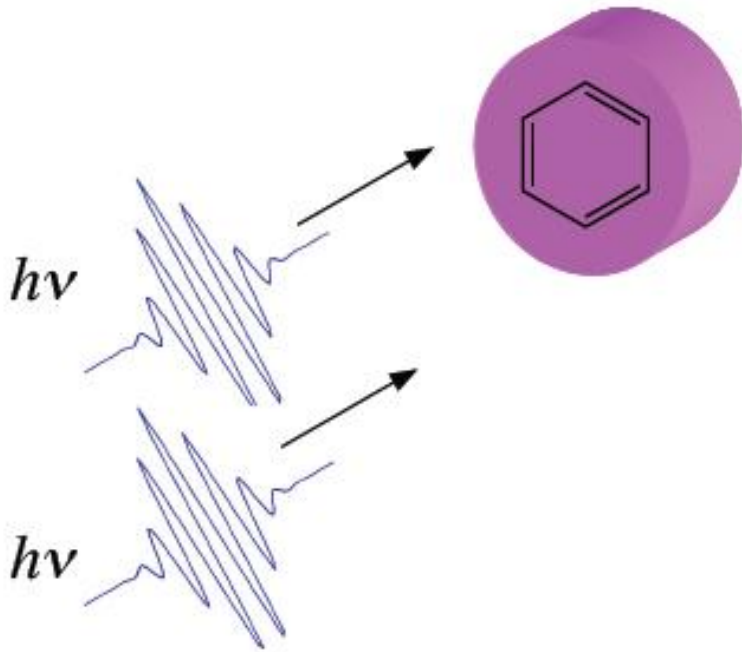
AzoPTCD



BePTCD

two-photon absorption cross-section

A pair of photon incident on a molecule



For two-photon absorption to occur, a pair of photons must be incident within a cross sectional area and within the lifetime of the virtual state, $\tau \sim 10^{-15}$ s

Relation between molecular TPA cross section and TPA absorption coefficient

$$\delta = \frac{h\nu\beta}{N}$$

Representative values for TPA transitions:

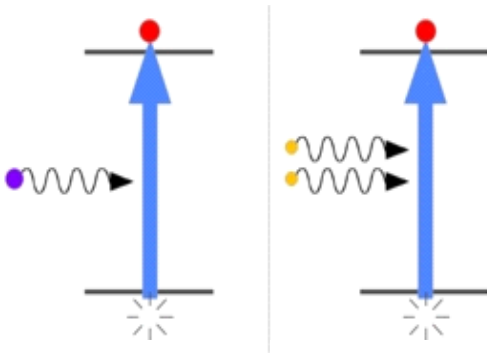
$$\delta = 10 \text{ to } 1000 \times 10^{-50} \text{ cm}^4 \text{ s photon}^{-1}$$

$$\beta = 10 \text{ to } 100 \text{ cm GW}^{-1} \text{ for a neat material}$$

Nonlinear spectrum

nonlinear absorption

$$\alpha = \alpha_0 + \beta I$$



nonlinear refraction

$$n = n_0 + n_2 I$$

intense laser (ultra short pulses)



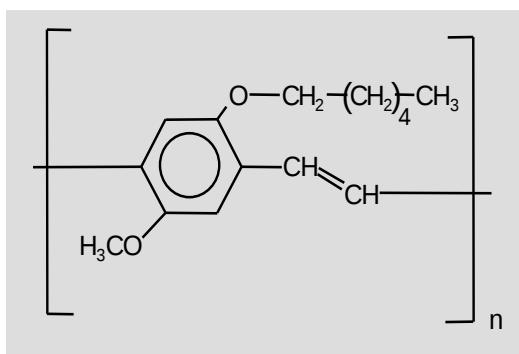
discrete λ s

$$\delta(\lambda) \quad n_2(\lambda)$$

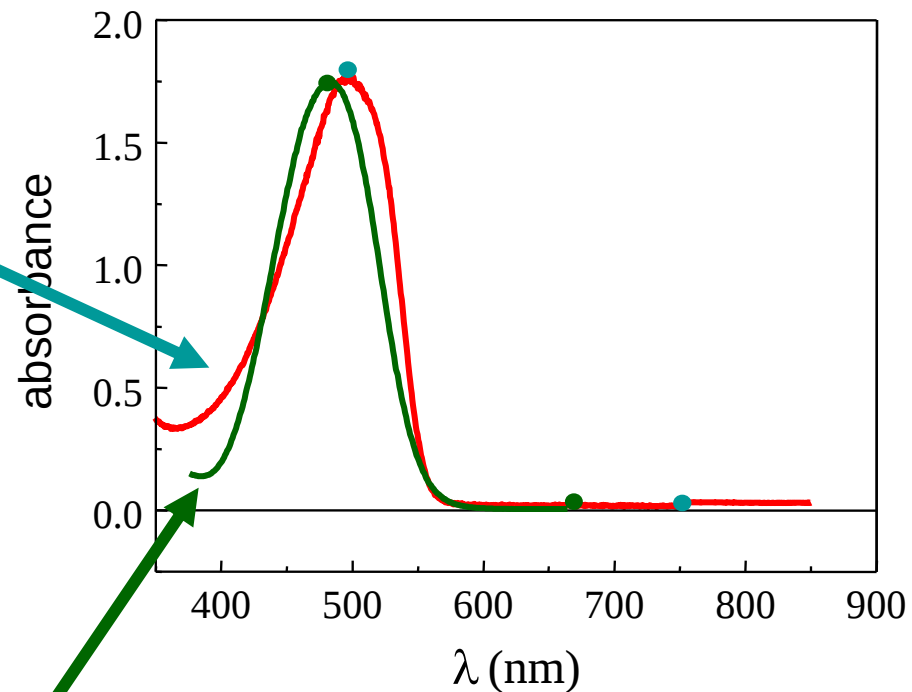
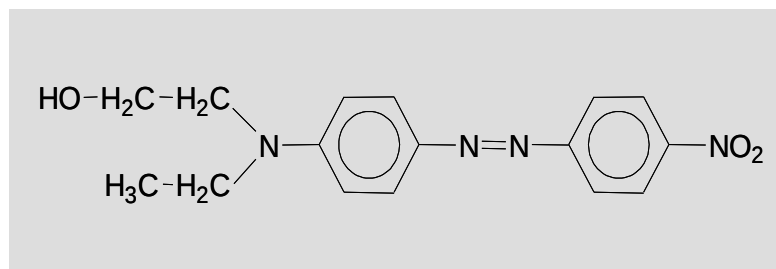
nonlinear spectrum ???

Linear absorption spectrum

MeH-PPV



DR1



Nonlinear absorption spectrum



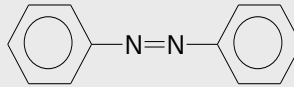
Optical parametric amplifier

460 - 2600 nm

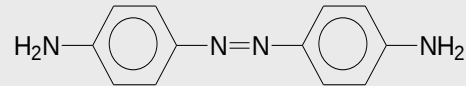
≈ 120 fs

20-60 μ J

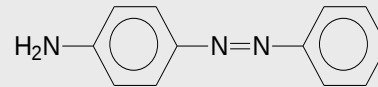
Azoaromatic samples



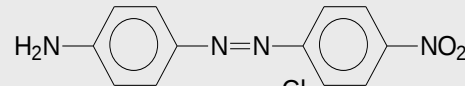
AZO



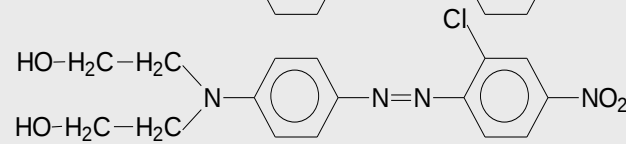
DIAMINO



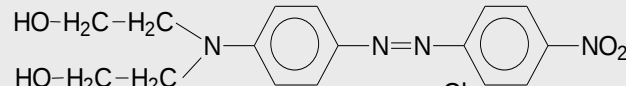
p-AMINO



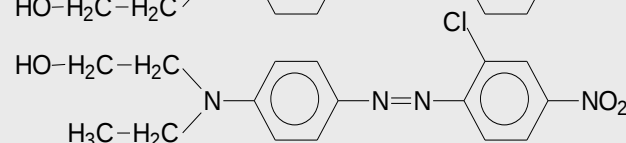
DO3



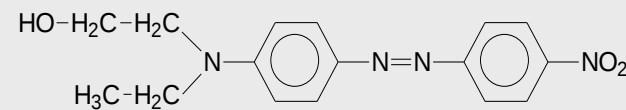
DR19Cl



DR19

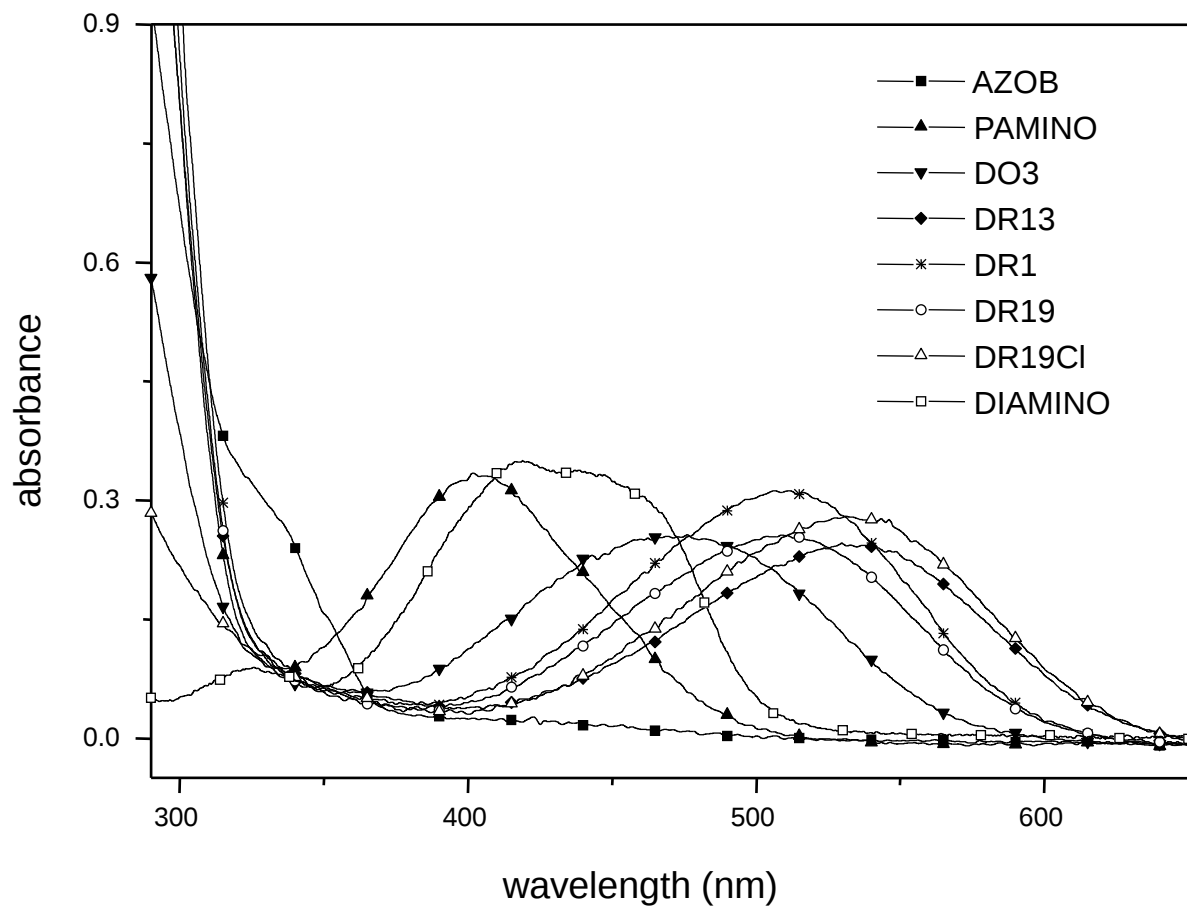


DR13

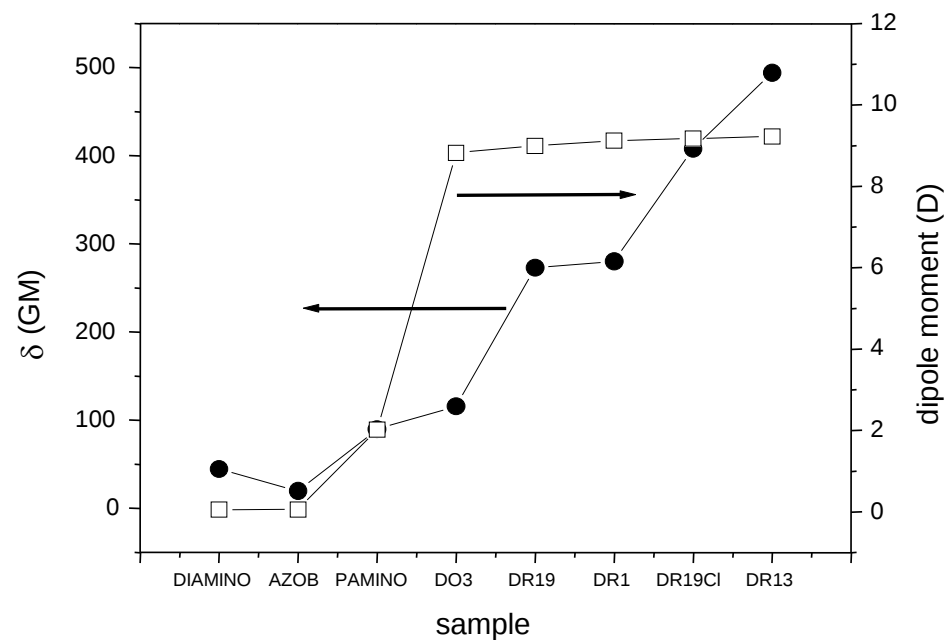
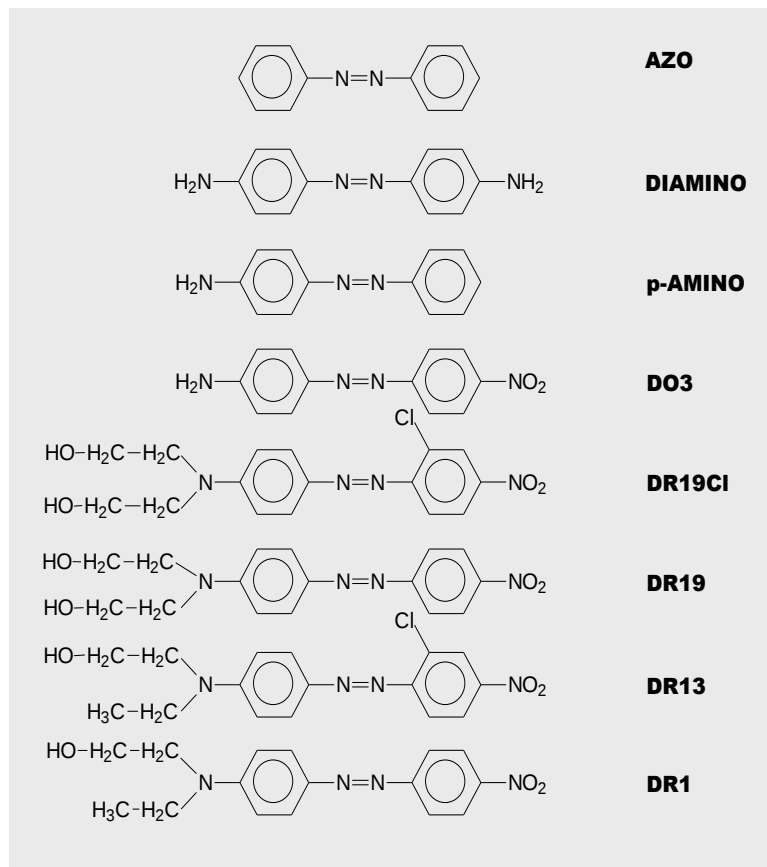


DR1

Linear absorption of azoaromatic compounds

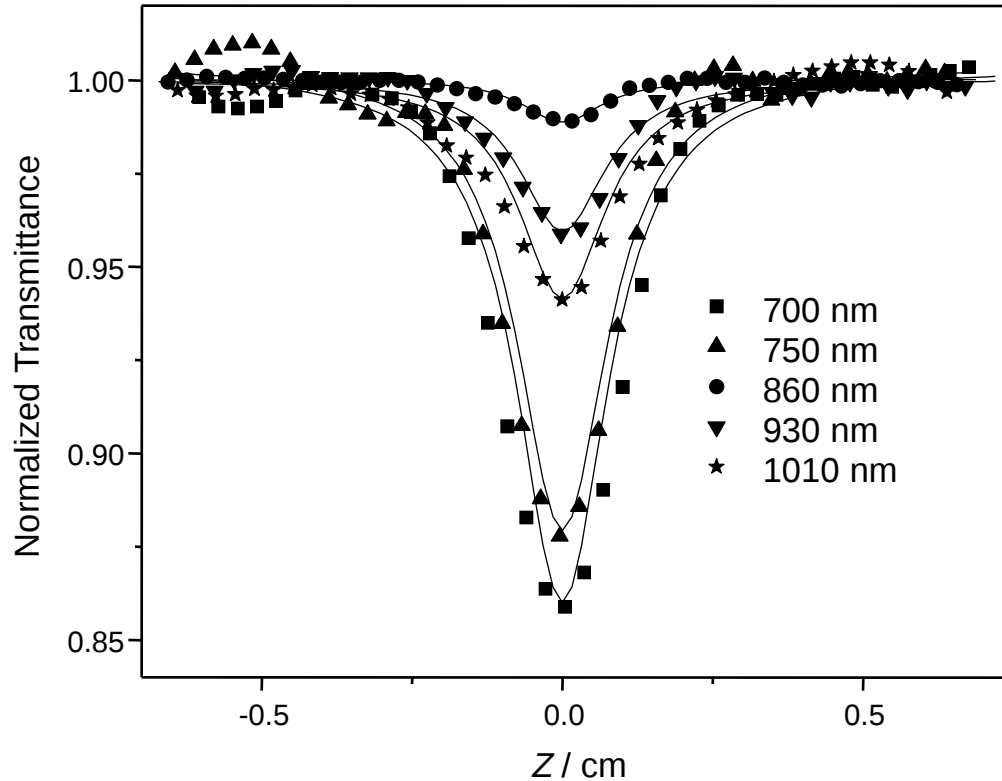


Two-photons absorption

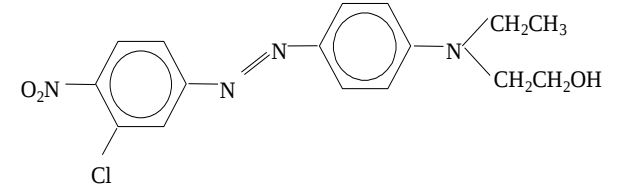


775 nm

Two-photon absorption



DR13

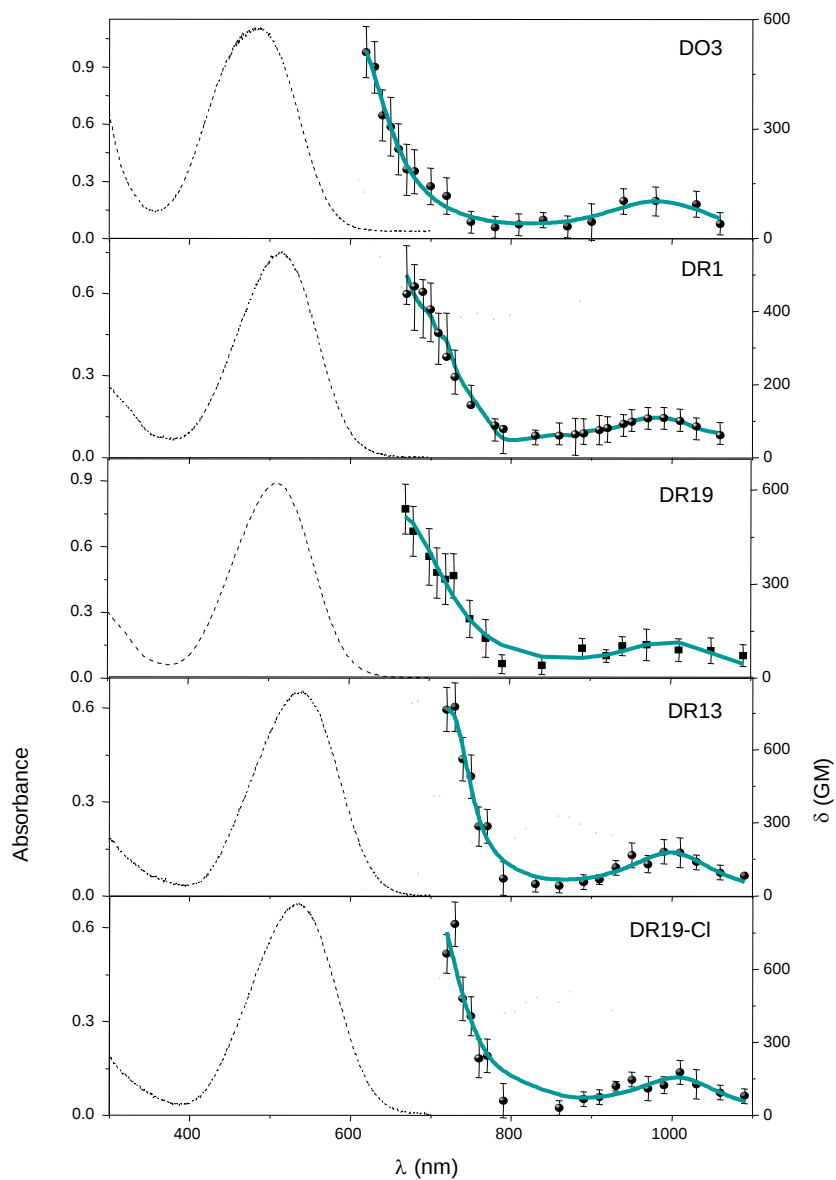


$$T(z) = \sum_{m=0}^{\infty} \frac{[-q_0(z,0)]^m}{(m+1)^{3/2}}$$

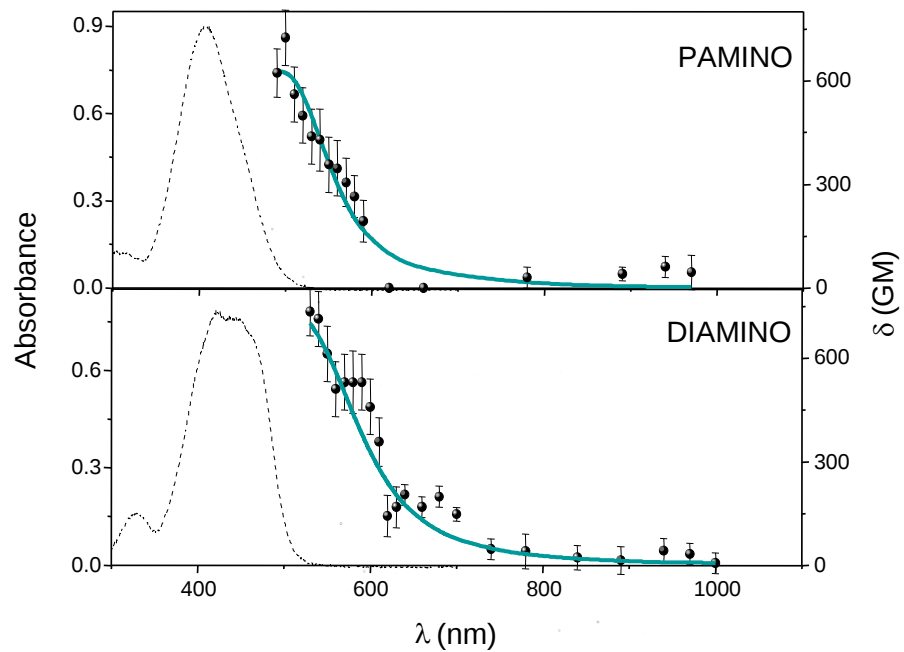
$$\alpha = \alpha_0 + \beta I$$

β : two-photon absorption coefficient

Psuedostilbenos

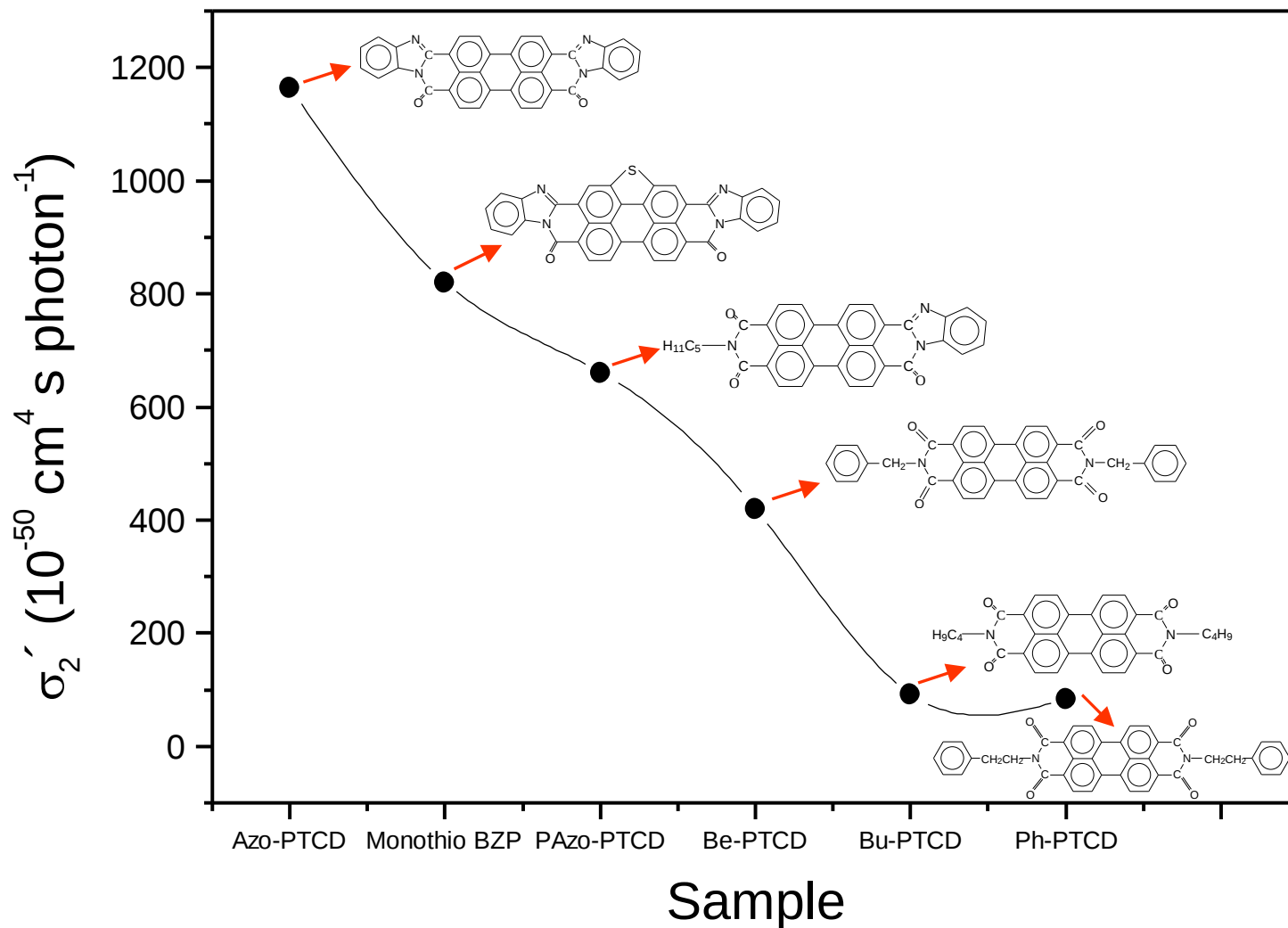


Aminoazobenzenos



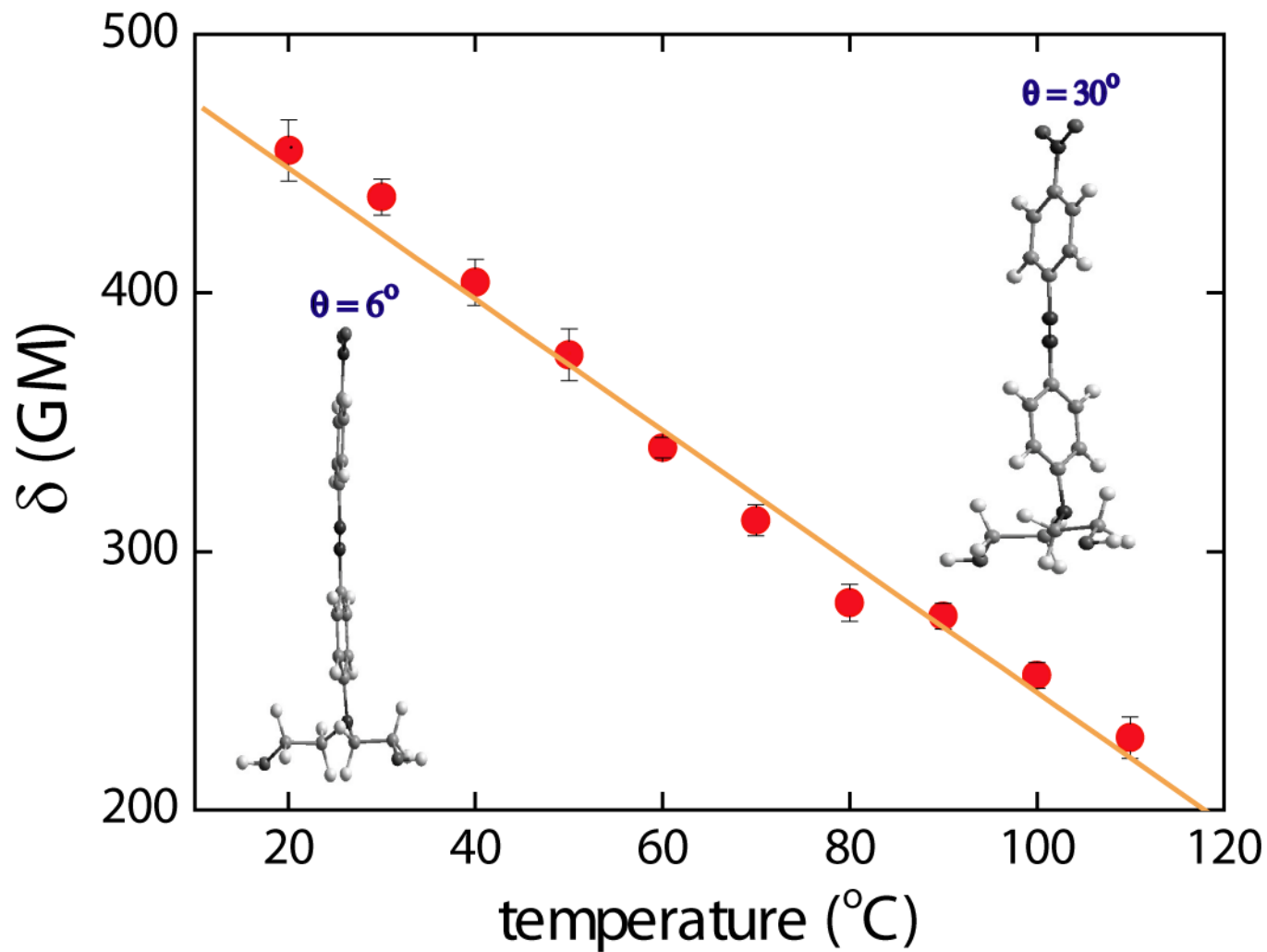
$$\delta(\nu) \propto \frac{\nu^2}{(\nu_{i0} - \nu)^2 + \Gamma_{i0}^2} \left[\frac{A_1}{(\nu_{f10} - 2\nu)^2 + \Gamma_{f10}^2} + \frac{A_2}{(\nu_{f20} - 2\nu)^2 + \Gamma_{f20}^2} \right]$$

Two-photon absorption of perylenes



Correlation between molecular structure and the nonlinear response

Planarity of the π -bridge



Thermally induced torsion in the molecular structure

Molecular design strategy

- Increasing the molecular conjugation
- Adding charged groups to the molecule
- Keep molecular planarity

Increasing the conjugation

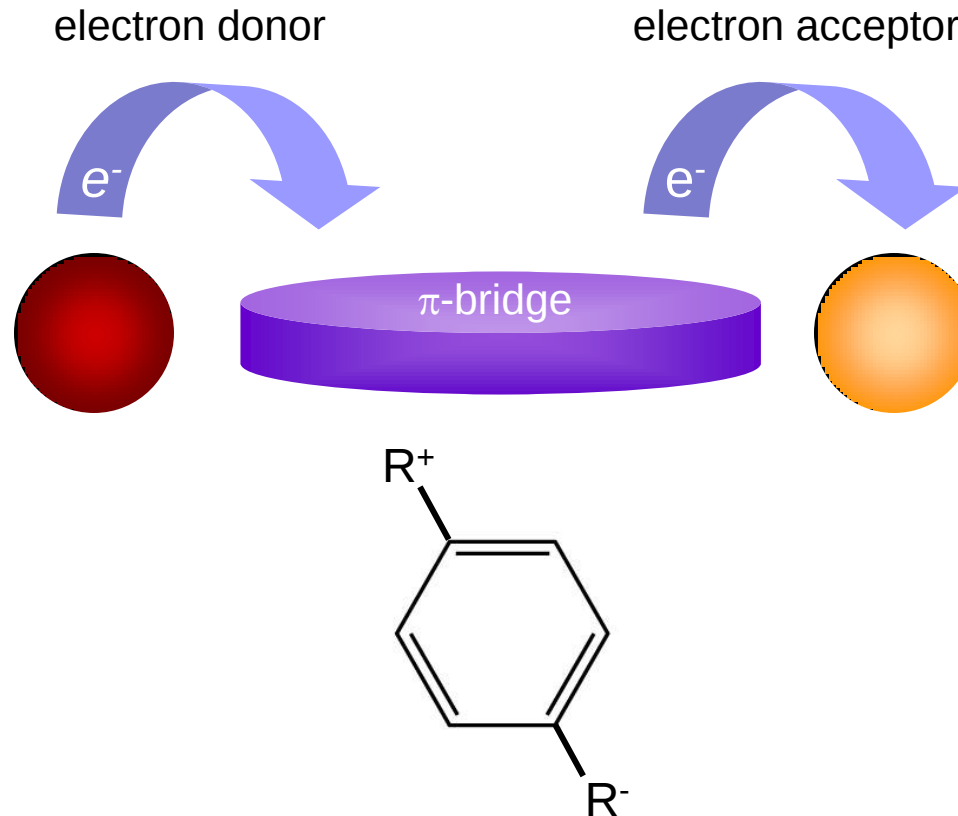


Increase in the optical nonlinearity



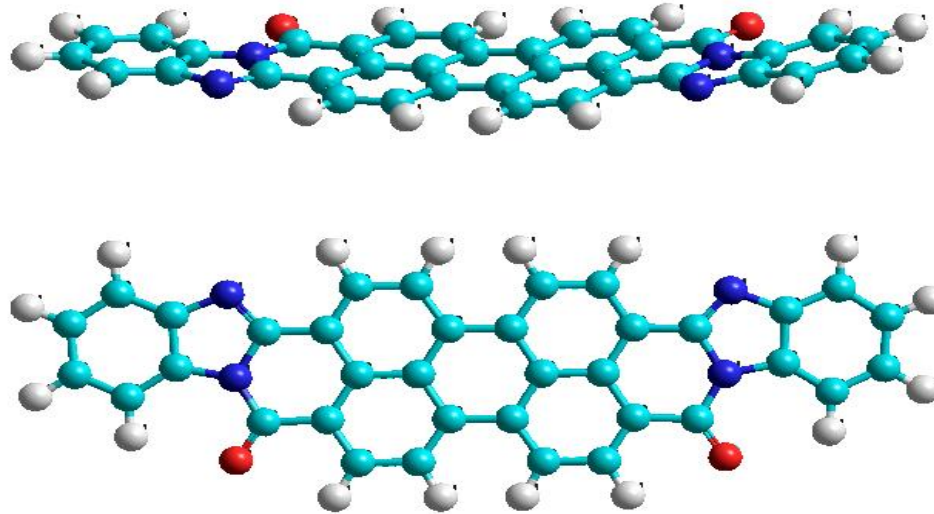
Increasing the π -conjugation

Donor and acceptor groups



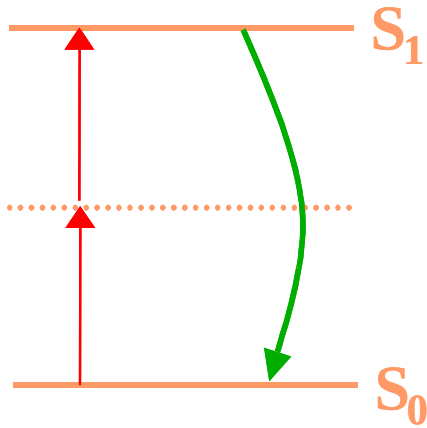
Incorporating electron donor and acceptor groups in a predictable way leads to an enhancement of the optical nonlinearity

Planarity of the π -bridge



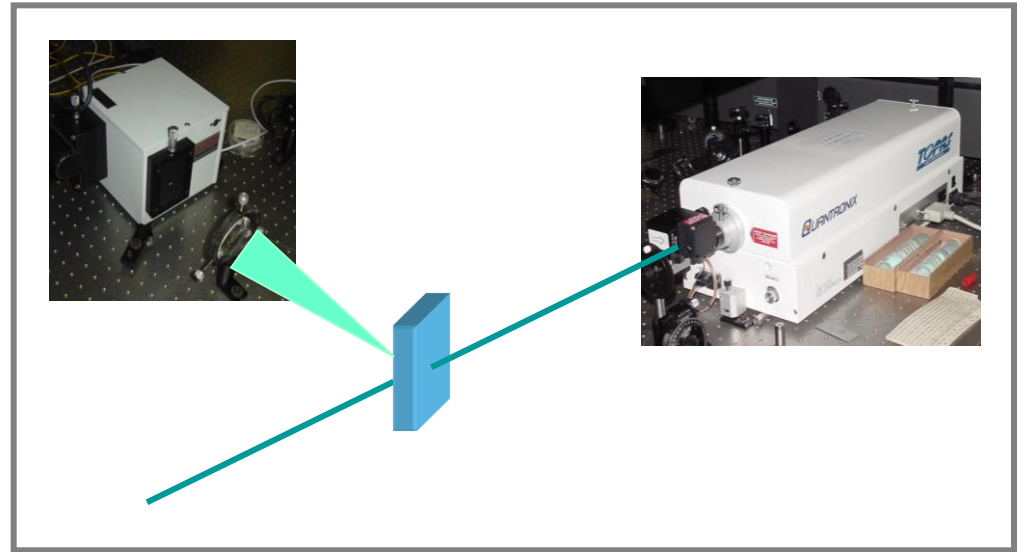
Perylene compounds are very planar molecules, which explains its high optical nonlinearities

Two-photon excited fluorescence

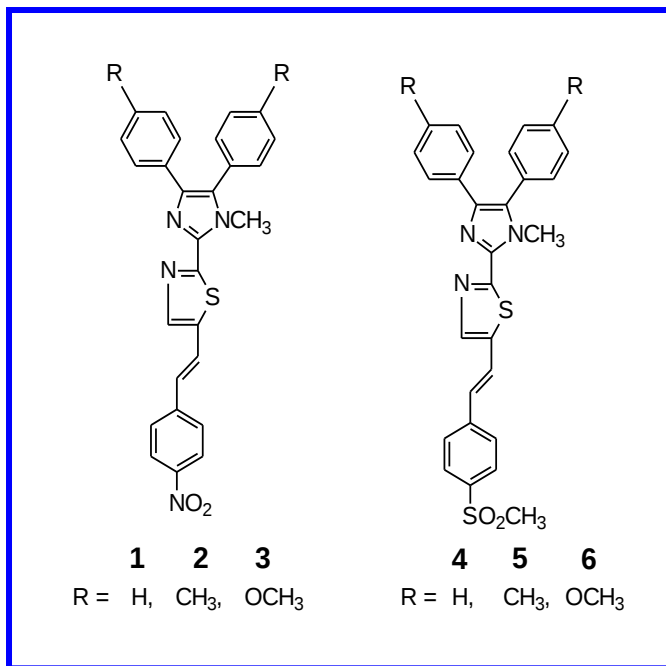
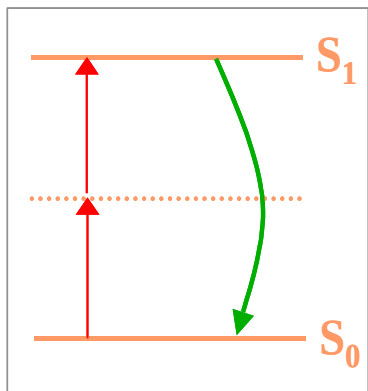


$$\frac{dN_2}{dt} \approx N_1 \delta I^2$$

$$F \approx I^2$$



Multi-photon excited fluorescence



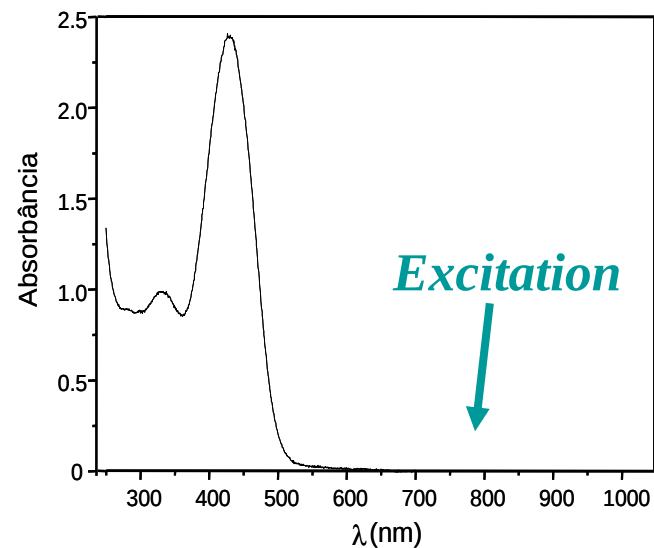
NITRO

SULFONIL

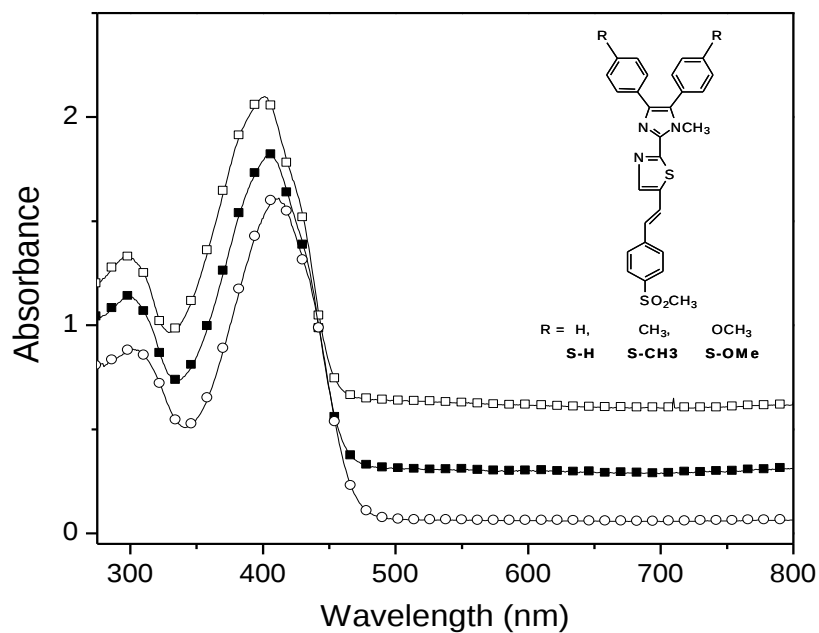
$$\frac{dN_2}{dt} \approx N_1 \delta I^2$$

$$F \approx I^2$$

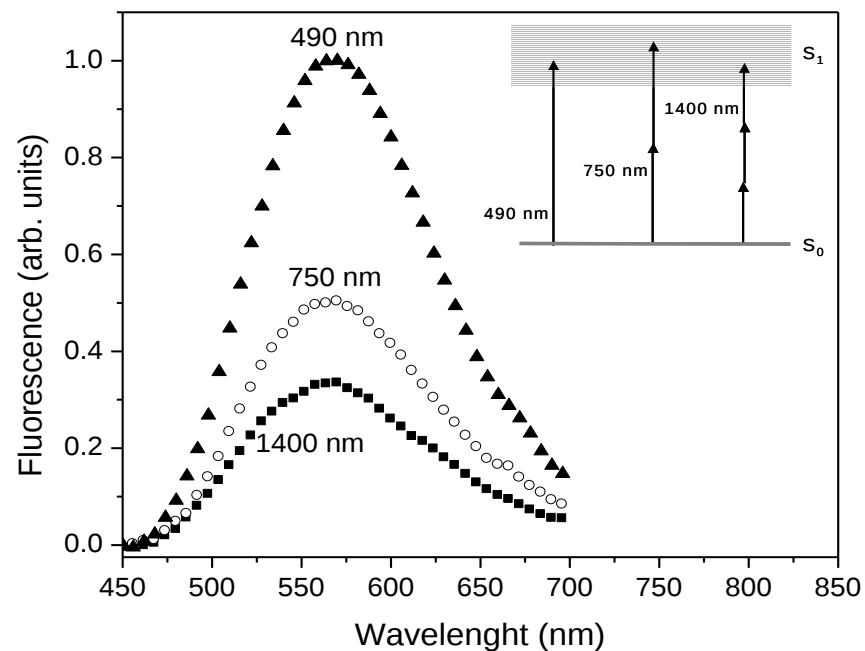
Y-shaped



Fluorescence excited by two and three photons

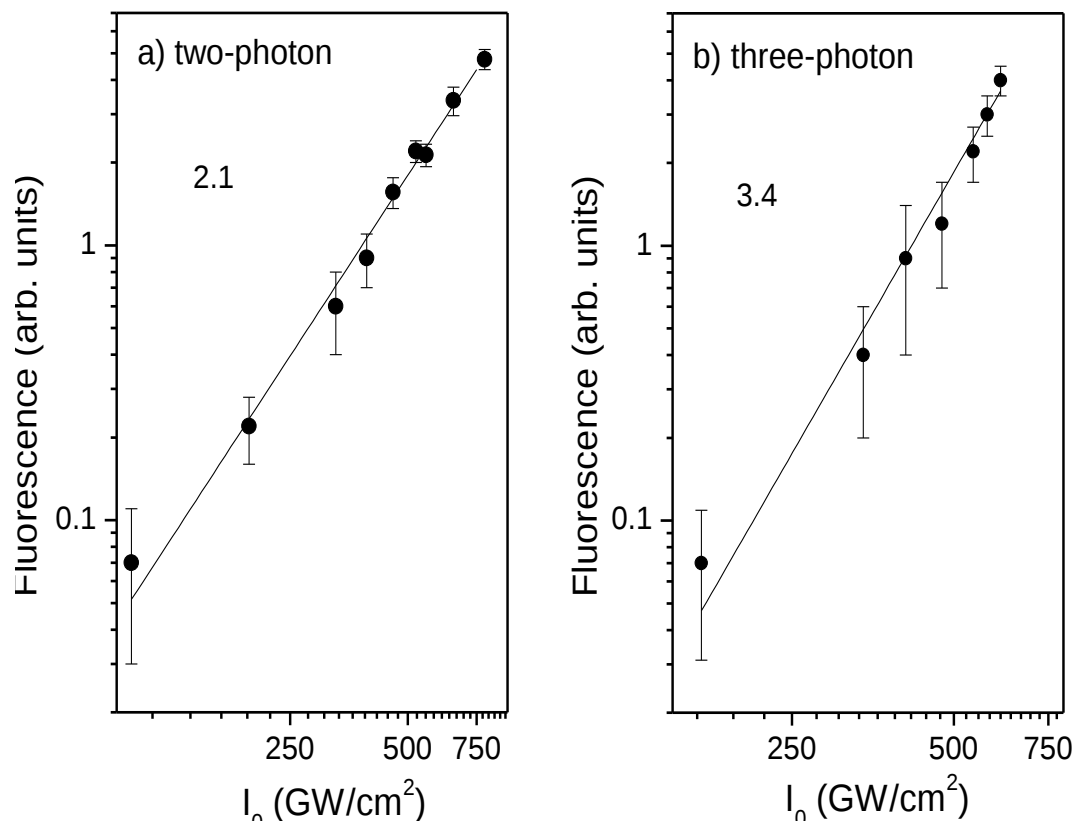


Absorption



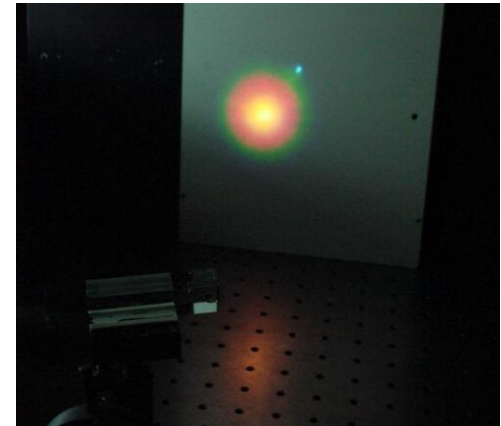
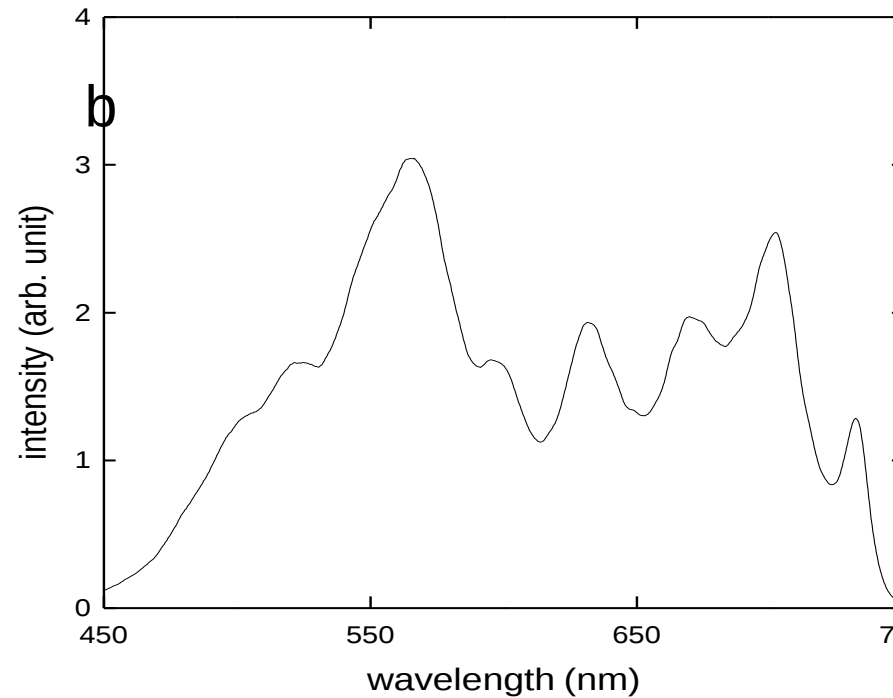
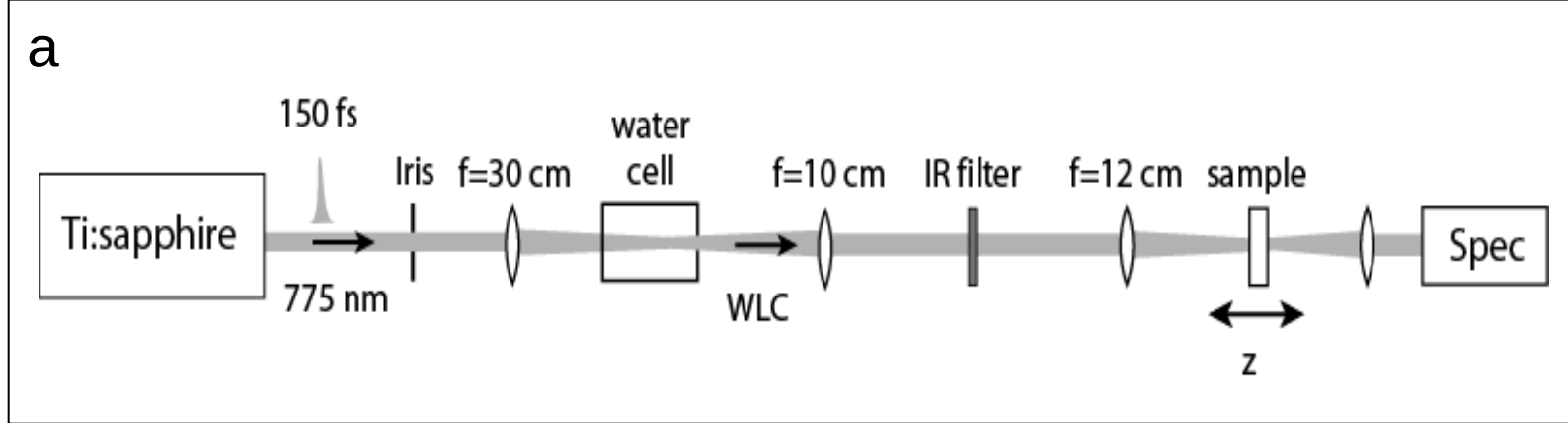
Fluorescence

Two-photon excited fluorescence

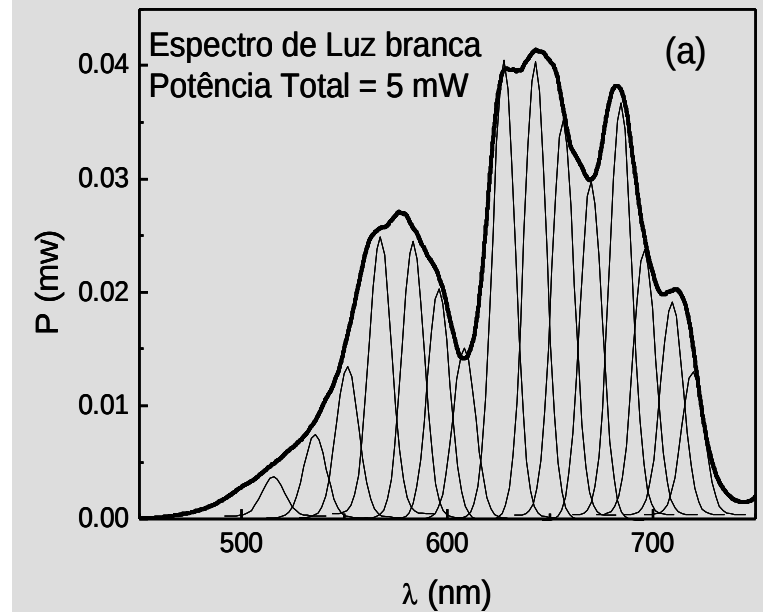
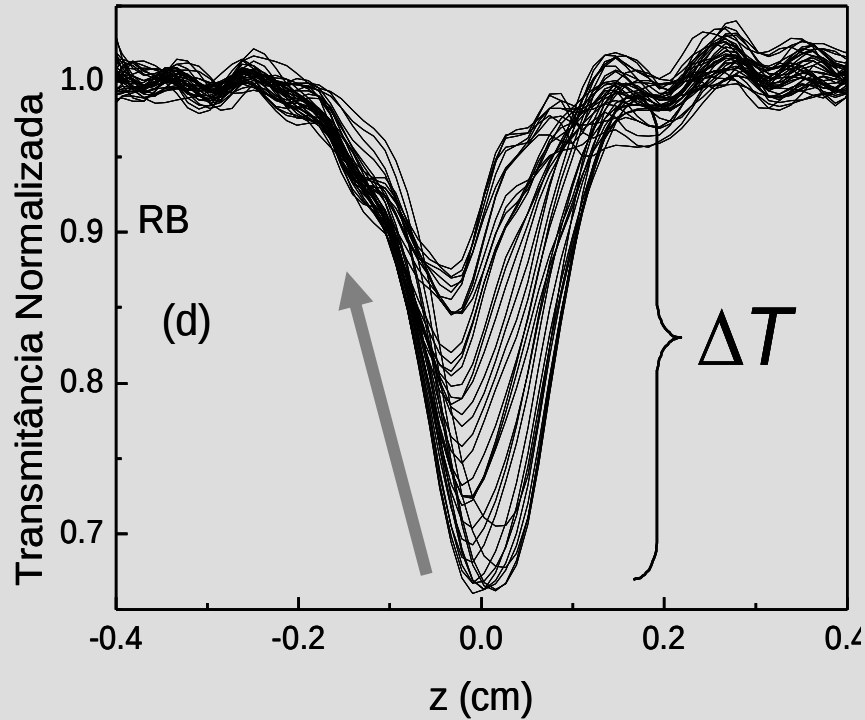


Sample	N (10^{18} molecules/cm ³)	σ_2 (10^{-50} cm ⁴ s) @ 750 nm	σ_3 (10^{-78} cm ⁶ s ²) @ 1400 nm
S-H	3.2	500	4.5
S-CH ₃	2.6	1450	5.6
S-Ome	2.4	1550	7.3

White light continuum Z-scan



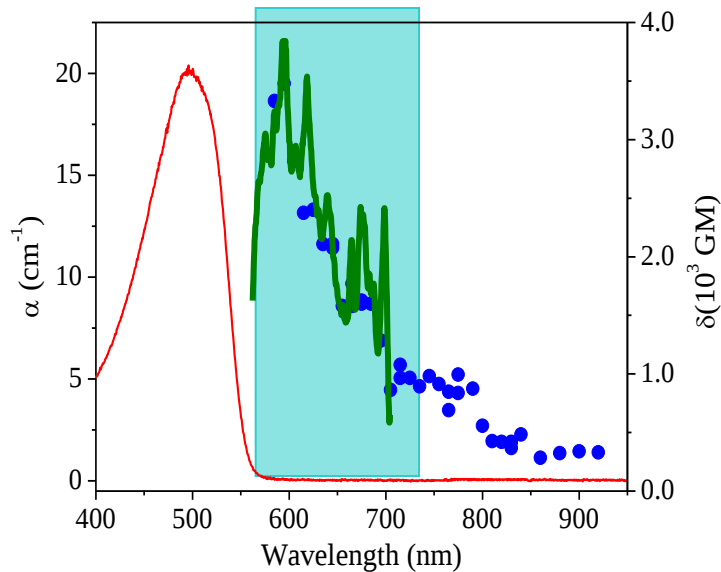
White light continuum Z-scan



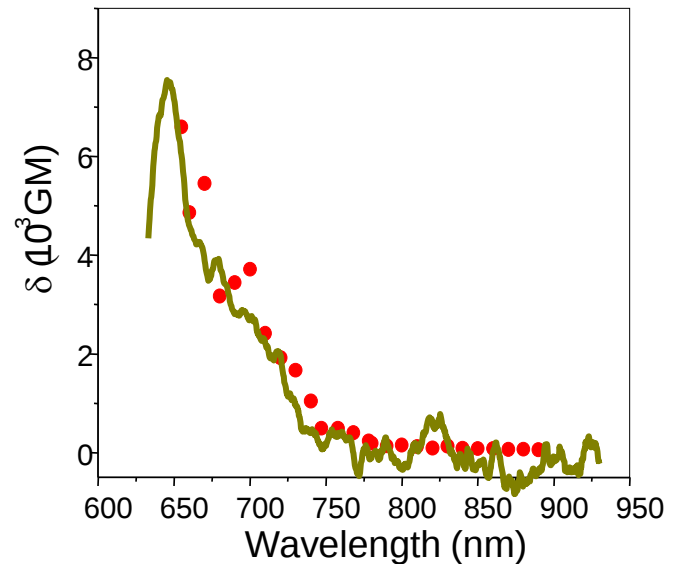
White light continuum Z-scan

Non resonant effects

MeH-PPV



Perylenes derivatives

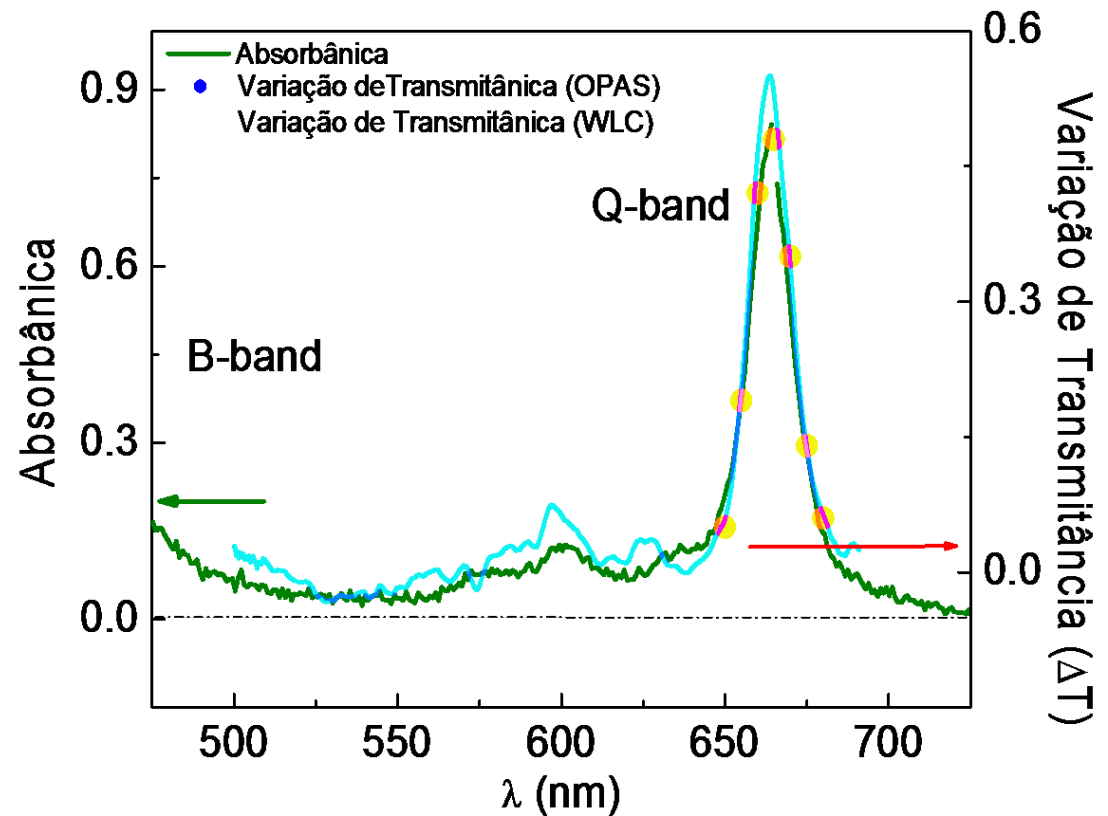
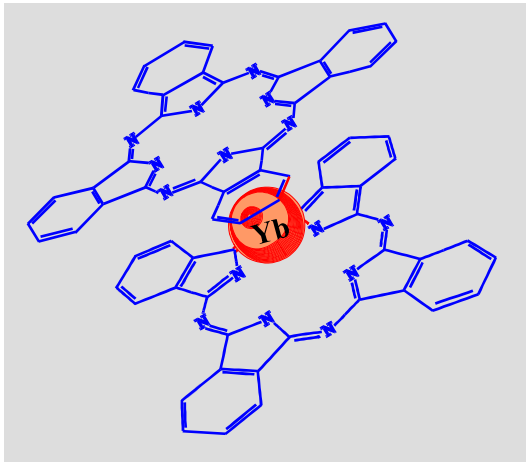


each measurement takes only a few minutes

White light continuum Z-scan

Resonant effects

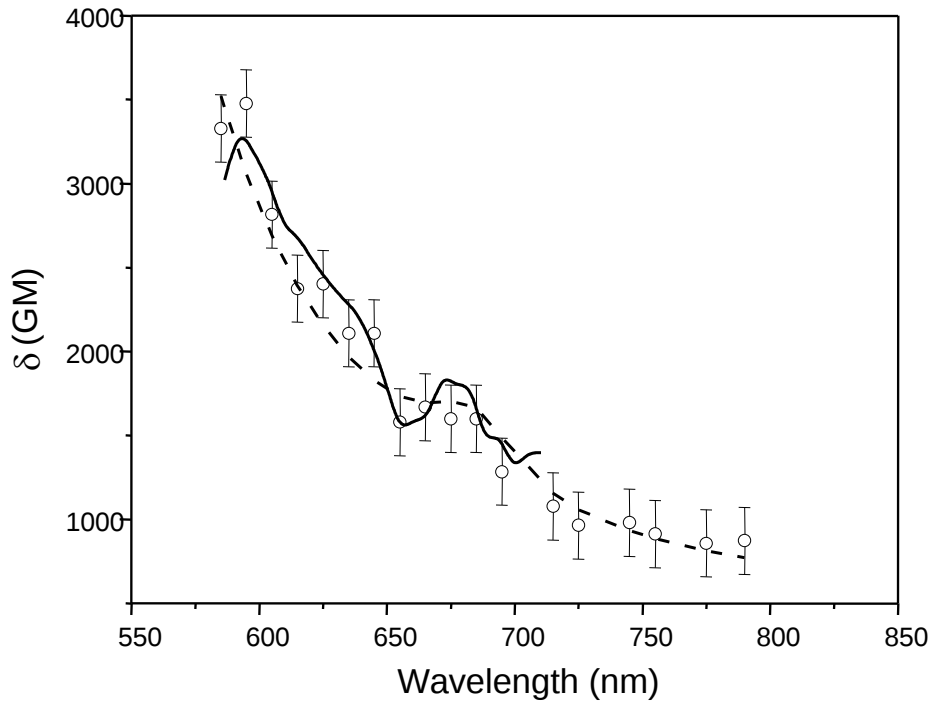
Bisftalocianina de Yb



White light continuum Z-scan

Two-photon absorption

MEH-PPV

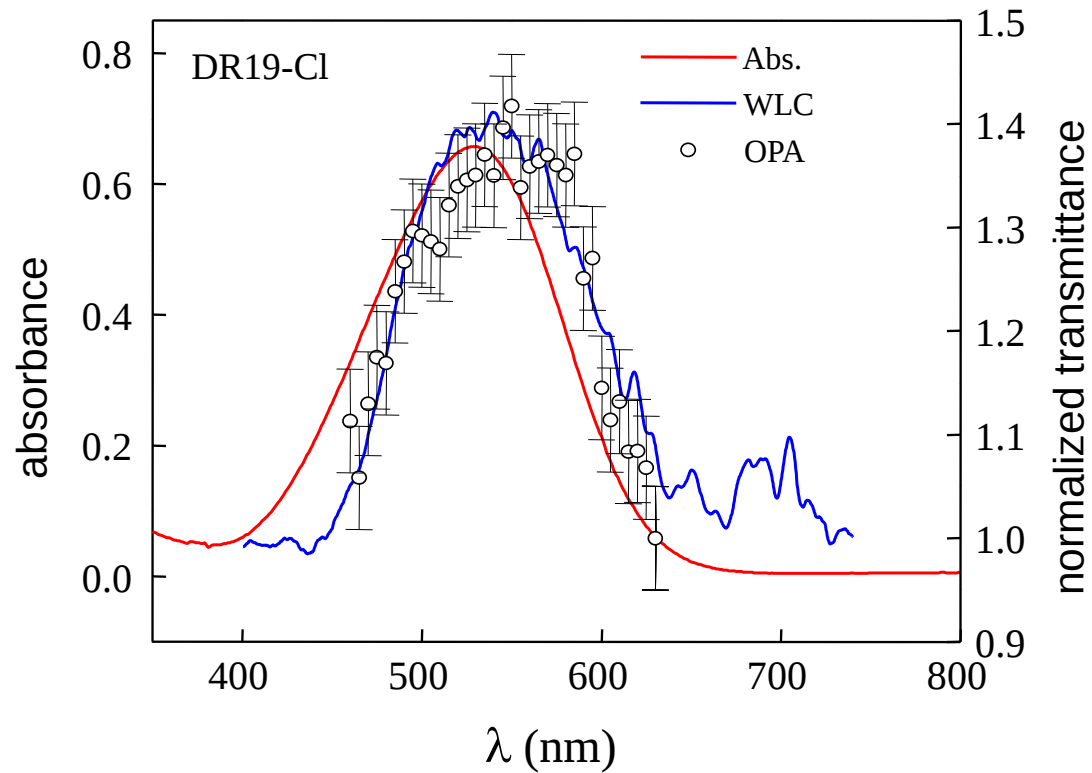


Circles : discrete Z-scan measurements

Dashed line: theoretical model

Solid line: Degenerate two-photon absorption cross-section spectra obtained from WLC Z-scan

White light continuum Z-scan



Resonant nonlinearities

saturated absorption

Thank you !

<http://www.fotonica.ifsc.usp.br>

<http://www.photonics.ifsc.usp.br>

